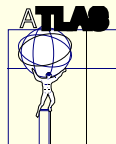


Beyond Standard Model with ATLAS at LHC

Ian Hinchliffe LBNL

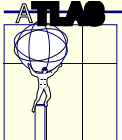
January 13, 2004



Outline

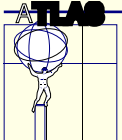
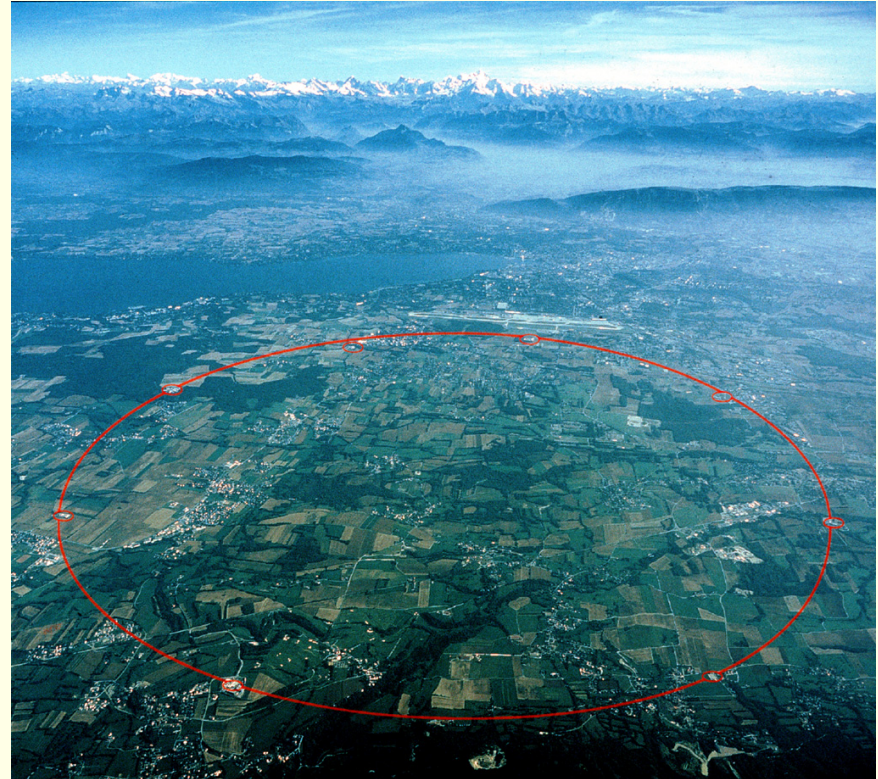
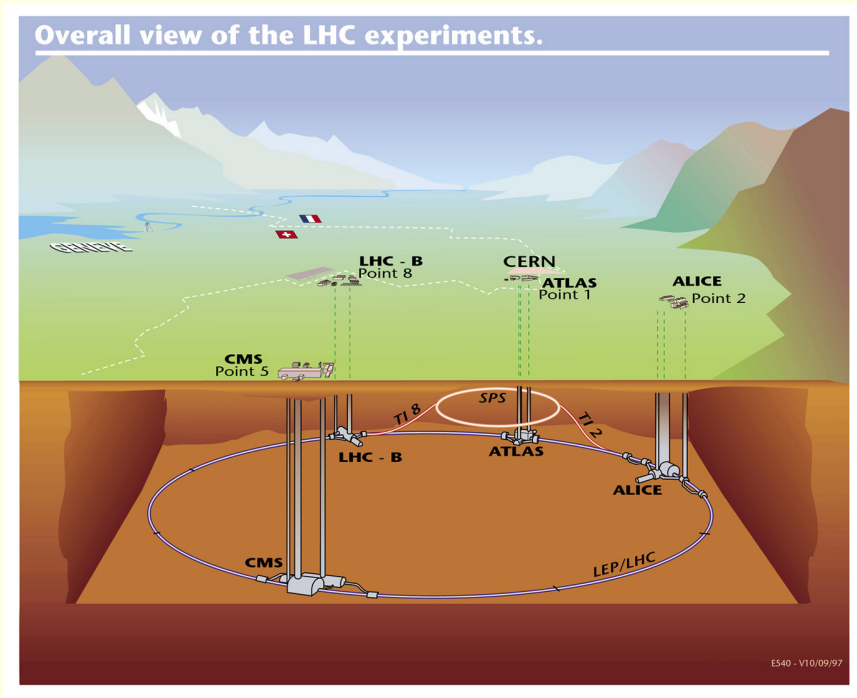
- LHC status
- ATLAS status
- Physics topics
 - SUSY
 - Extra Dimensions
 - Little Higgs Models

Many physics studies concentrate on “ultimate” goals.
I will give an indication of what might happen quickly



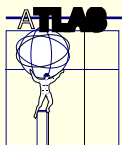
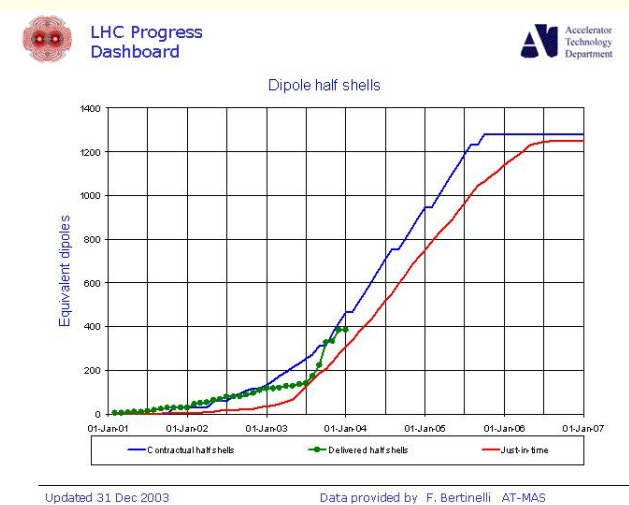
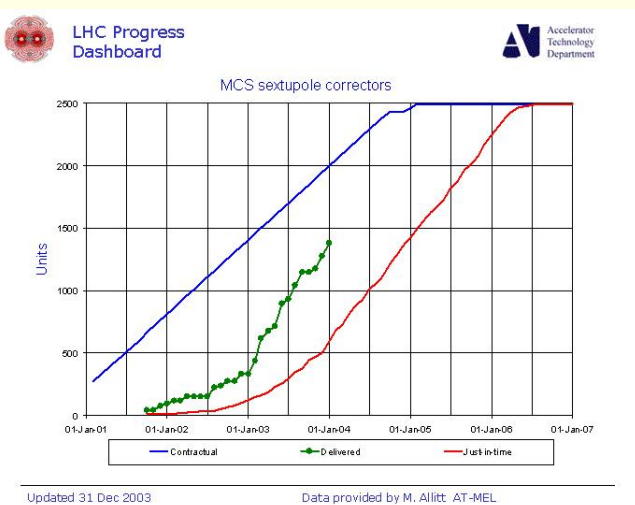
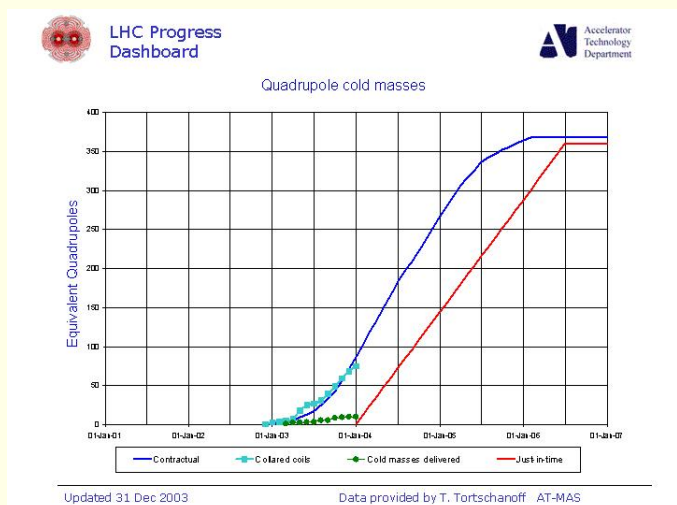
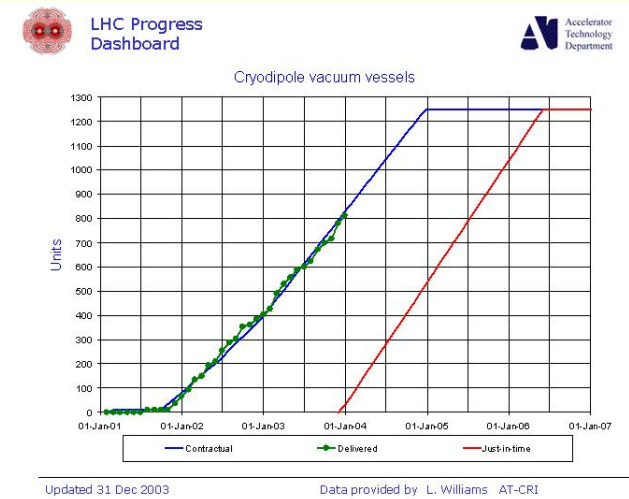
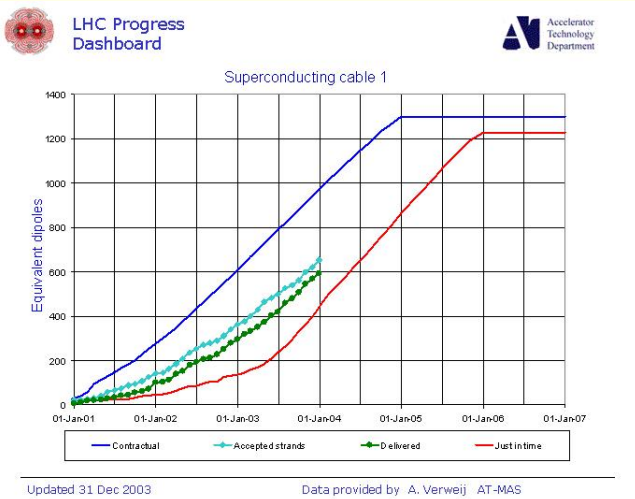
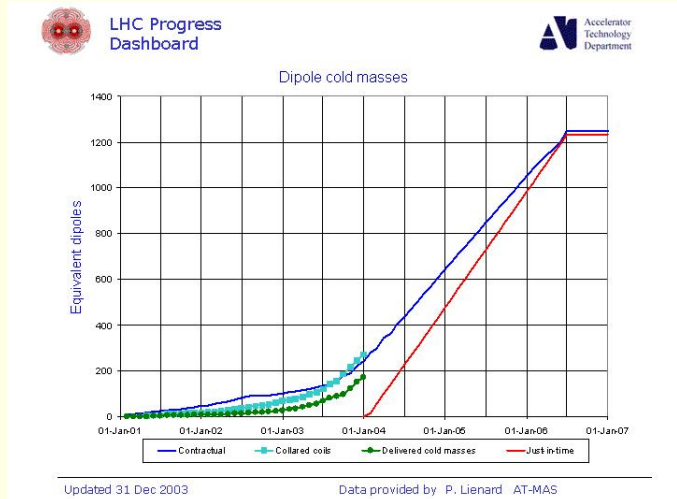
LHC Status and Schedule

“Overall, the project’s cost is stable and its schedule unchanged, foreseeing first beam in April 2007 with first collisions following in June.” L. Maiani Dec 19 2003.



Status is updated monthly at

<http://lhc-new-homepage.web.cern.ch/lhc-new-homepage/DashBoard/index.as>

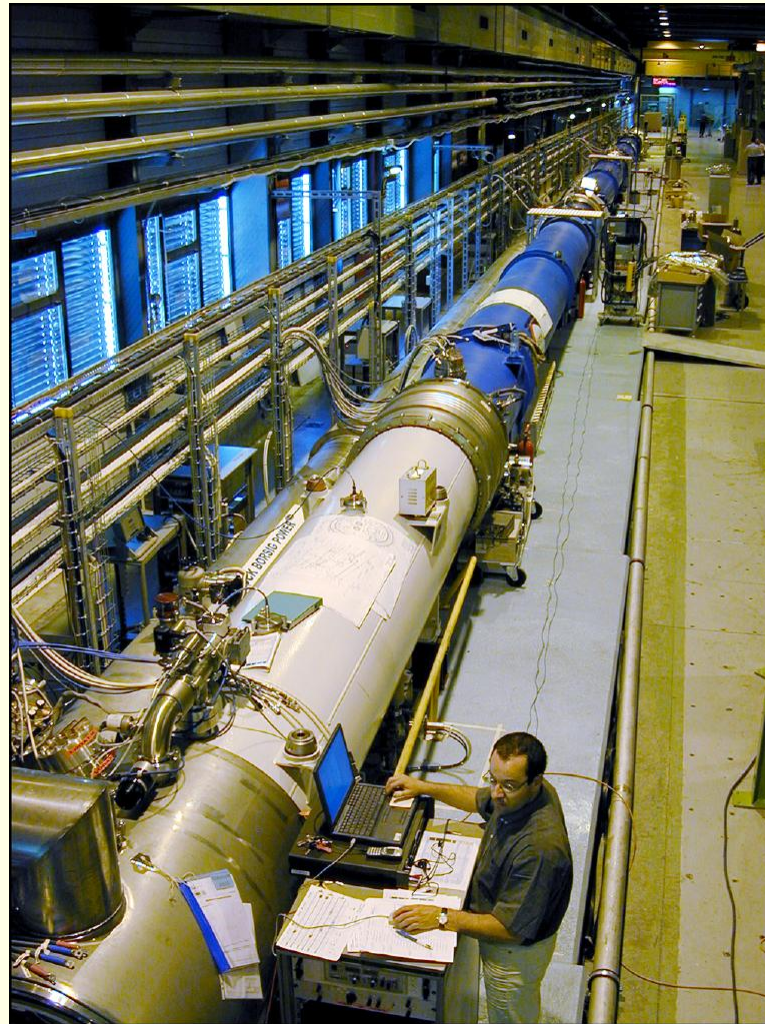




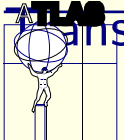
First magnet installed into transfer line December 2003



Transfer magnets awaiting installation

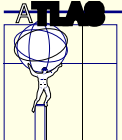


String Test 2001

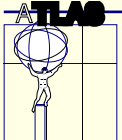


LHC operation

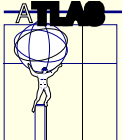
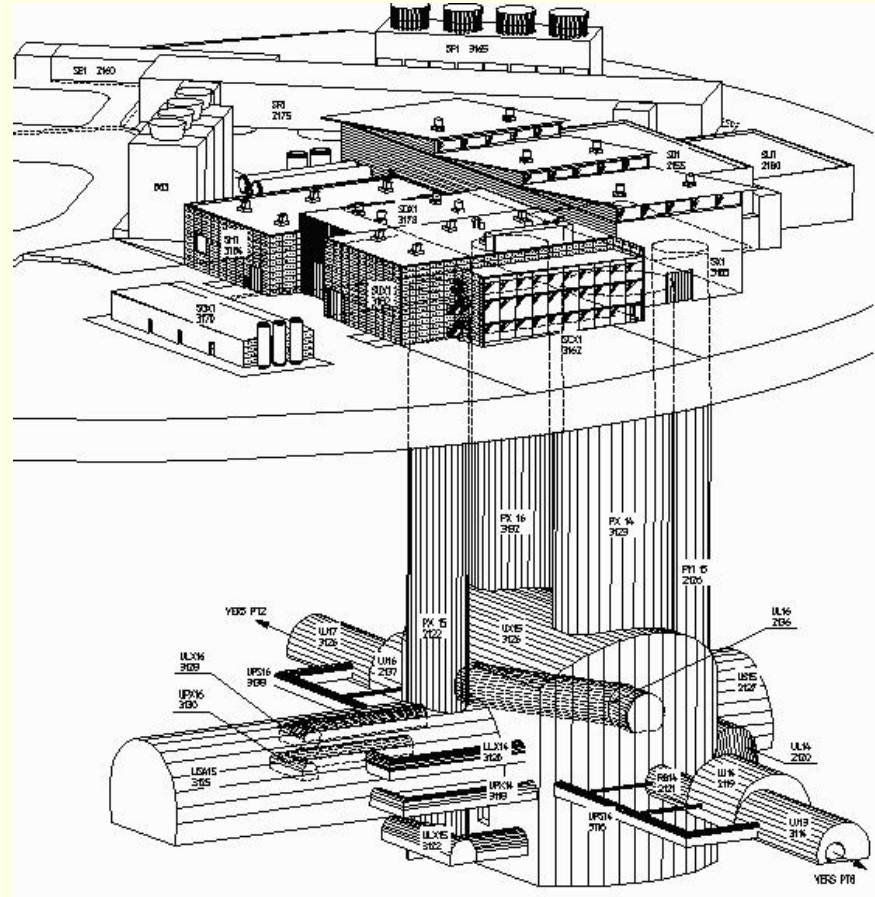
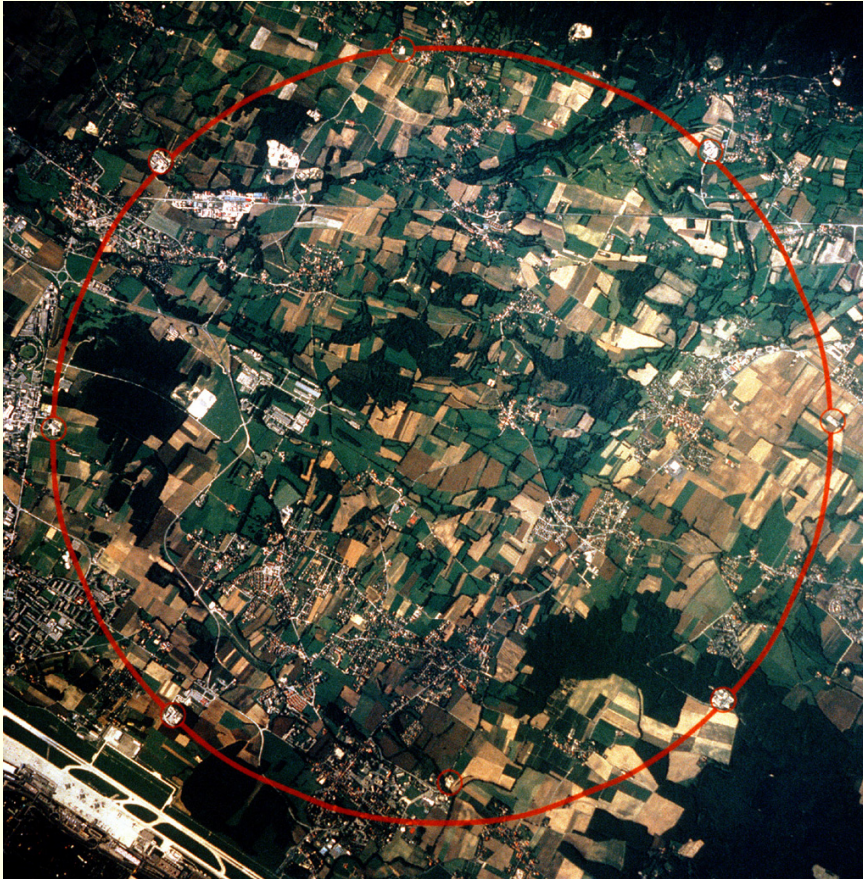
- Single Beam operation – April 2007
- Collisions – June 2007
- Operation in “low luminosity mode” for 3 years $2 \times 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$
- 1 month per year of heavy ion running.
- Full luminosity in ~ 2010 $10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$, multiple interactions per crossing cause some degradation in performance *e.g.* b-tagging.
- Some ATLAS elements have been staged and will not be available at turn-on. Middle layer of pixels, some muon chambers, little impact at low luminosity.
- Trigger/DAQ staging means less rate – impacts *b*–physics: Could be restored with extra funding.



Atlas-Buildings and location



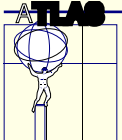
Surface building – across street from CERN main gate



Above



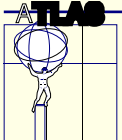
Below



Last weeks photos

LHC Beam is at *A* and *C*

In the center is the support structure for the detector



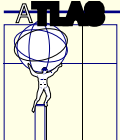
Overview of ATLAS

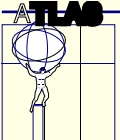
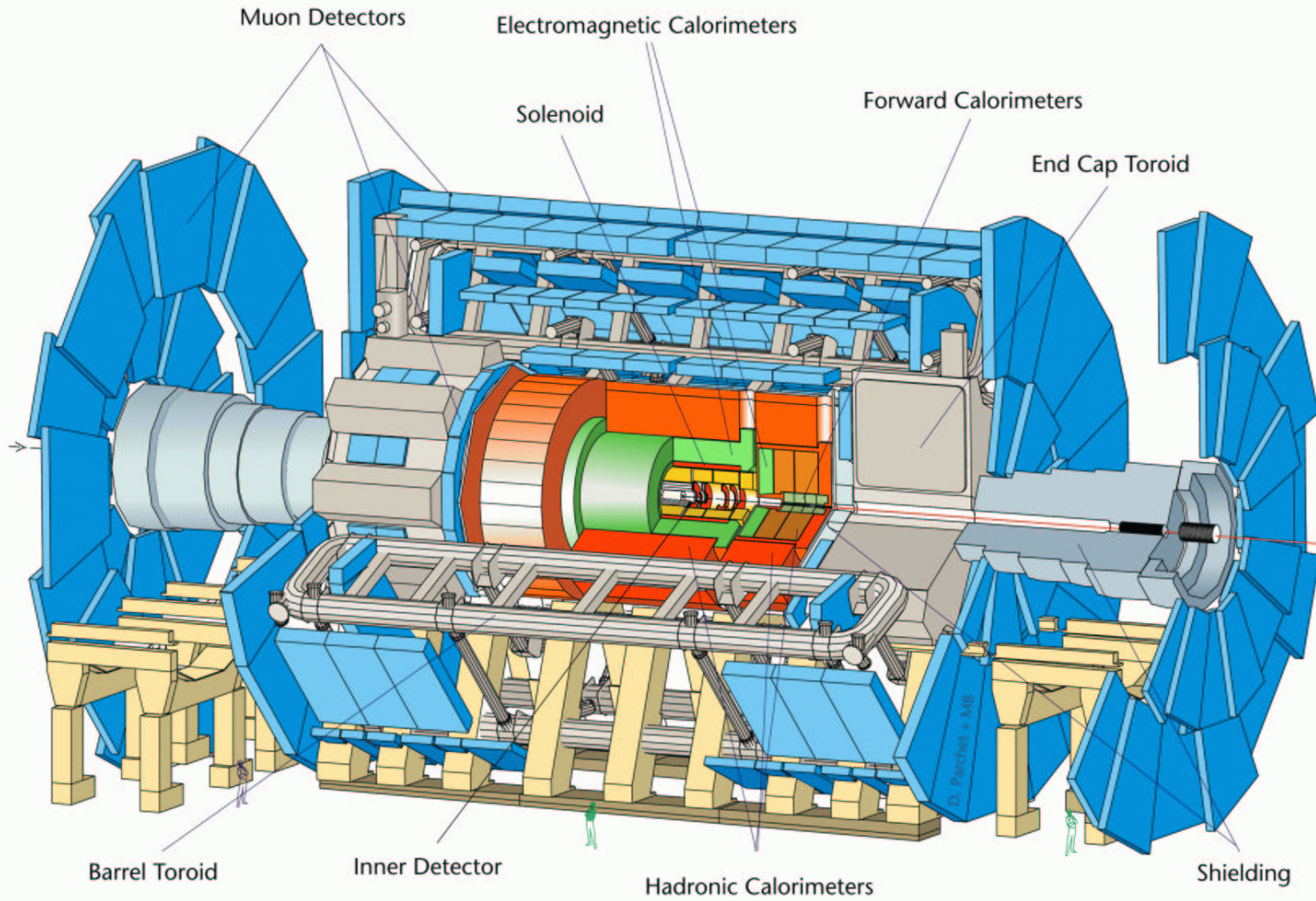
ATLAS and CMS are aimed at “new physics”

“Full acceptance” for physics objects, *i.e.* leptons and jets, missing E_T

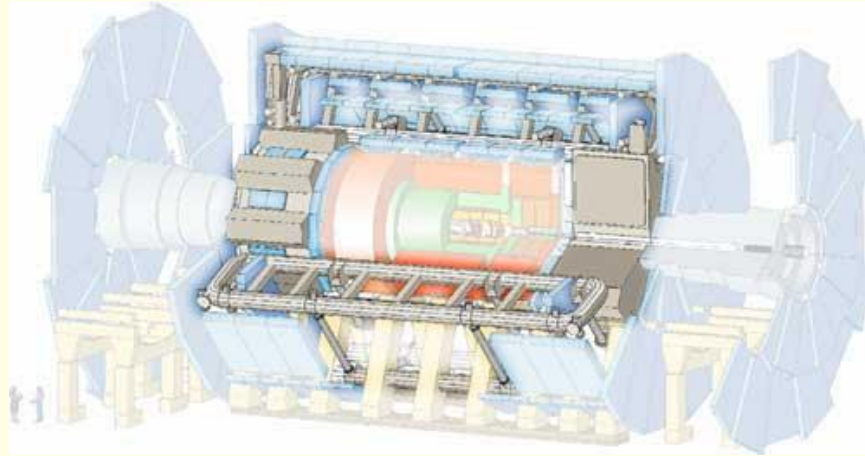
Many detector choices driven by specific physics goals (*e.g.* LiAr Calorimeter) Equal response for e and μ

Physics performance is expected to be similar to CMS, technology choices are quite different





Magnet system



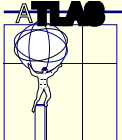
Solenoid – Central tracking



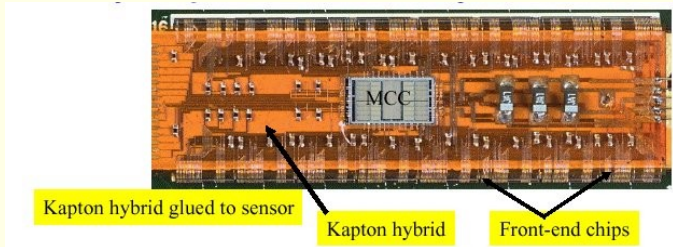
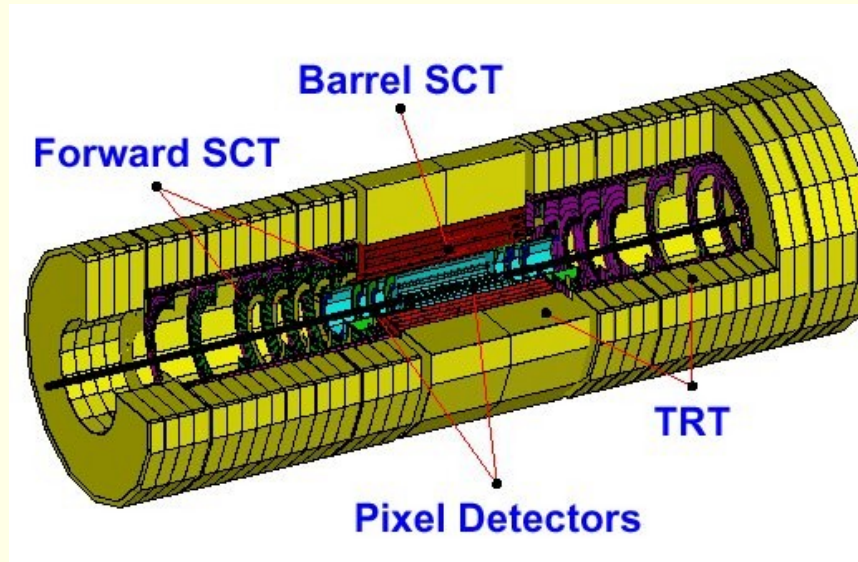
Muon endcap



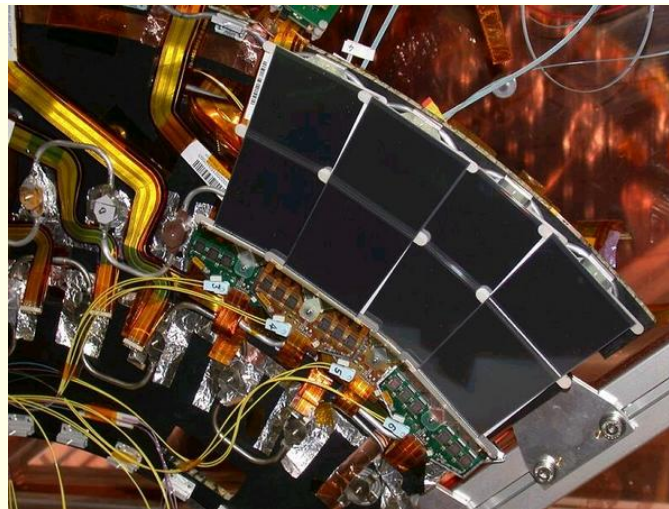
Central toroid under assembly



Inner Detector



Pixel Hybrid

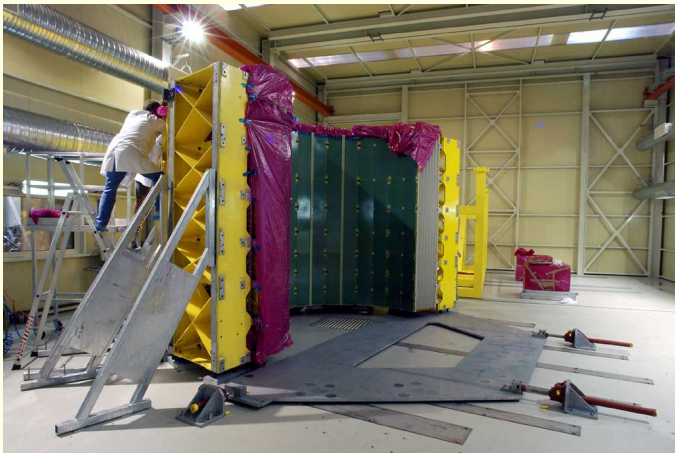
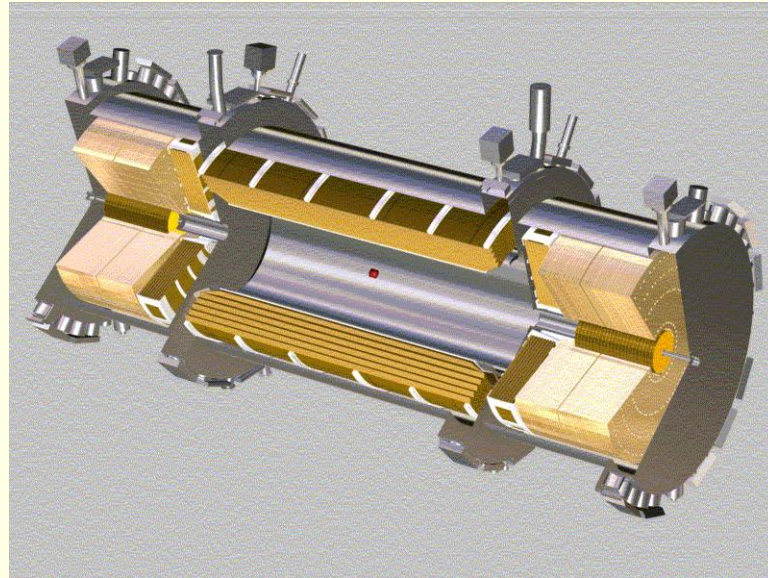


Forward Si Strip Module

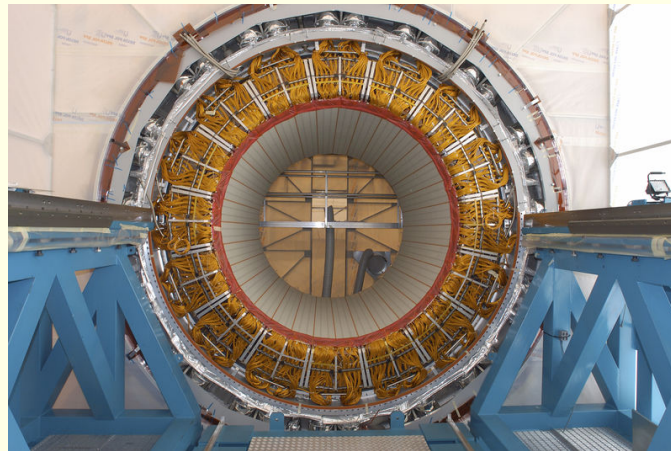


Forward TRT wheel

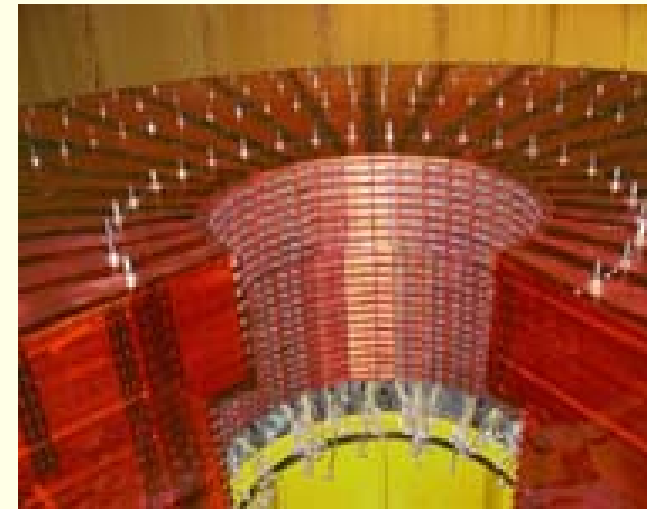
LiAr (EM) Calorimeter



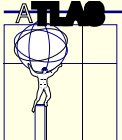
Barrel EM



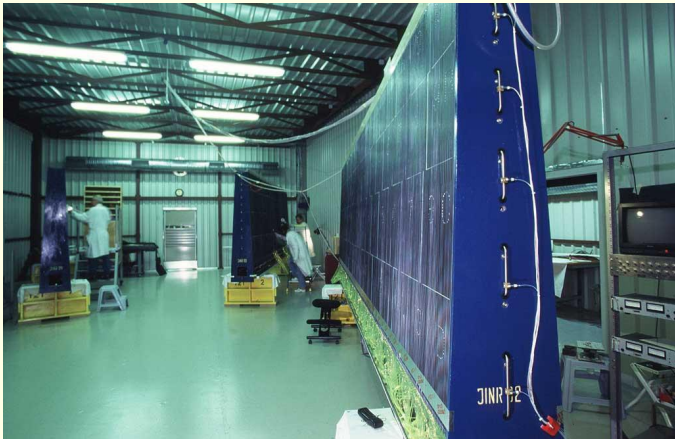
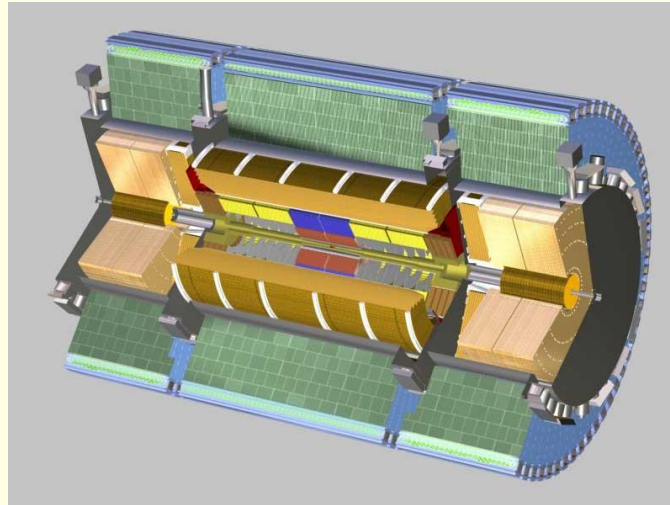
Barrel Cryostat



hadronic end cap



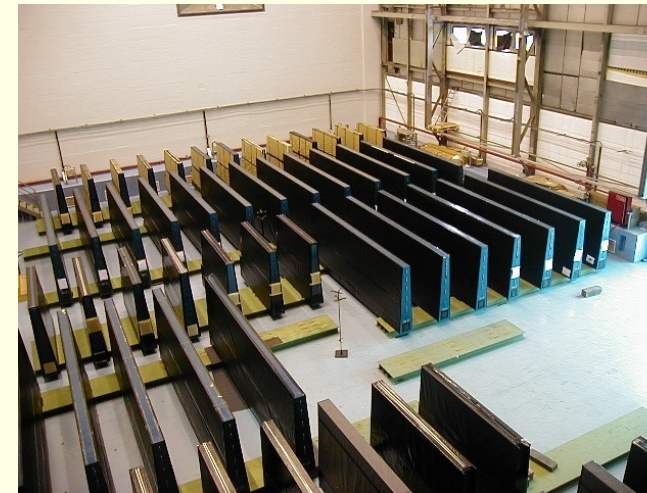
Tile (Hadronic) Calorimeter



Single element

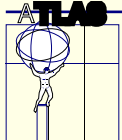
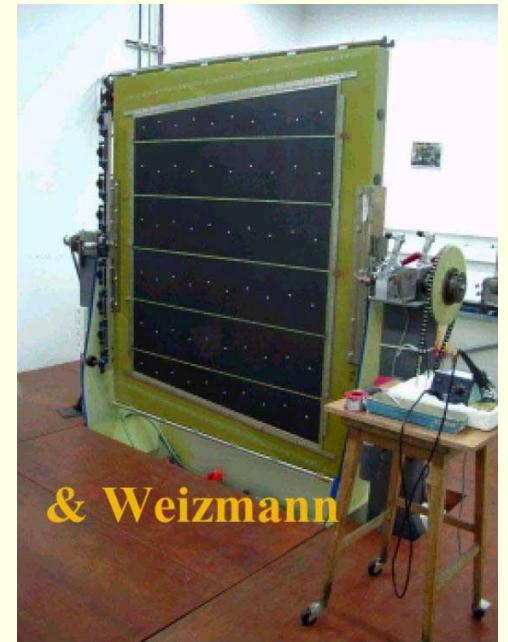
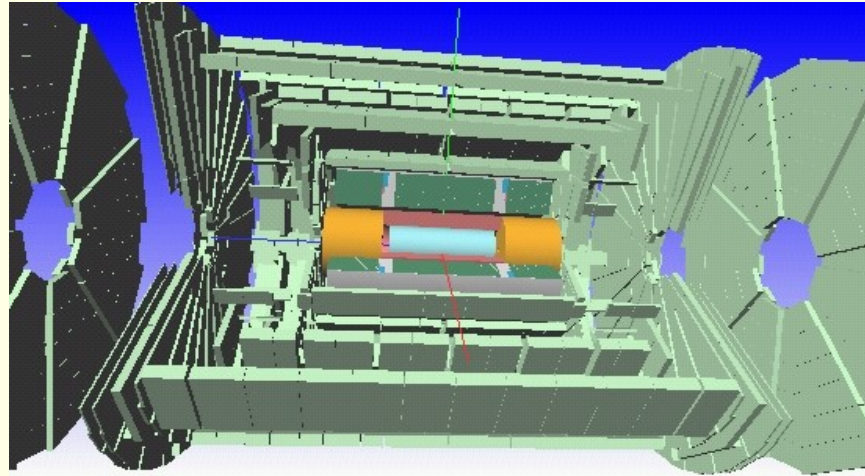


Barrel



Sections in storage

Muons



Characteristic New physics signatures at LHC

Not all present in all models

Heavy objects decay into Standard Model particles with high energy

\cancel{E}_T from ν or other new particles

High Multiplicity of large p_t jets

Many isolated leptons – from W , Z or directly produced

Copious b production – “democratic decays?”

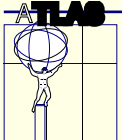
Large Higgs production – this may be a standard model particle

Isolated Photons

Quasi-stable charged particles – like a heavy muon.

N.B. Production of heavy objects implies subset these signals

Important for triggering considerations



Backgrounds – Measuring and Calculating

At present, we rely on MC for signal and background estimates

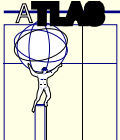
There are uncertainties in rates from PDF's, higher order QCD

Most of these do no matter at the moment, **They will matter once data appears**

My concern: underlying and min-bias events

Affects process that need forward jet tagging *e.g.* *WW – scattering* or central jet veto

Will be measured once data exists and MC will be tuned to agree...



Little Higgs Models

All data consistent with SM ($g - 2$???)

New particles of mass $\lesssim 10\text{TeV}$ are constrained EW fits, FCNC limits *etc*

Calculate with a cut off $\Lambda = 10\text{TeV}$

$$\text{top loop } \delta m_h^2 = \frac{3}{8\pi^2} \lambda_t^2 \Lambda^2 \sim (2\text{TeV})^2$$

$$\text{W/Z loops } \delta m_h^2 \sim \alpha_w \Lambda^2 \sim -(750\text{GeV})^2$$

$$\text{Higgs loop } \delta m_h^2 \sim \frac{\lambda}{16\pi^2} \Lambda^2 \sim -(1.25m_h)^2$$

$$m_h^2 \sim (100\text{GeV})^2$$

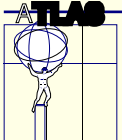
Fine tuning of Higgs mass seems to require something else $\sim 1\text{TeV}$

Most dangerous terms are top loop, Higgs loop, W/Z loops

Solve these and problem is $\gtrsim 10\text{TeV}$ where we know nothing

SUSY solves it up to $\sim M_{Planck}$ by removing all quadratic divergences.

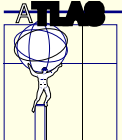
Can arrange *ad-hoc* cancellations by adding a few particles but need a symmetry



Little Higgs models (2)

- Models try to arrange new particles to cancel these effects
- Do this by extending the symmetries of the Standard Model so that the cancellations are forced by the new symmetries – SUSY is best example
- Need a theory with a broken global symmetry to get a massless Goldstone boson.
- Must break the symmetry “in a small way” so that this Goldstone Boson can have interactions and a VEV and play the role of the Higgs.
- Will solve the hierarchy problem; cancellations will appear as needed.
- The models are not simple (they may be “elegant”) and not complete.

Arkani-Hamed, Georgi, Burdman, Schmalz,

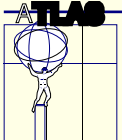


LHC signals

What is the minimal stuff??

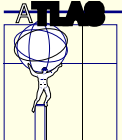
- Something to cancel the top loop.
In the example this is T decays via $T \rightarrow Zt$, $T \rightarrow Wb$, $T \rightarrow ht$ with BR in the proportion 1 : 2 : 1
Ratio is test of model
- Something to deal with the W loop
In the example this is the gauge bosons of the other $SU(2) \times U(1)$.
Once the masses are specified their couplings have one free parameter (θ)
- Something to deal with the H loop
In the example here this is the Higgs triplet ϕ which is produced via WW fusion
- Very small effects $< 5\%$ in $h \rightarrow gg$ and $h \rightarrow \gamma\gamma$

Masses and decays are model dependent. Higgs sector is most model dependent



Expected range of masses

- Fine tuning means that $f = \frac{\Lambda}{4\pi} < 1TeV(\frac{m_H}{200GeV})^2$
- $m_T < 2TeV(\frac{m_H}{200GeV})^2$
- $M_{WH} < 6TeV(\frac{m_H}{200GeV})^2$
- $m_\phi < 10TeV$



New Quark

Properties determined by two parameters λ_1/λ_2 and mass.

Two production mechanisms $qb \rightarrow q'T$ and $gg \rightarrow T\bar{T}$: Former depends on $t - T$ mixing and therefore on λ_1/λ_2

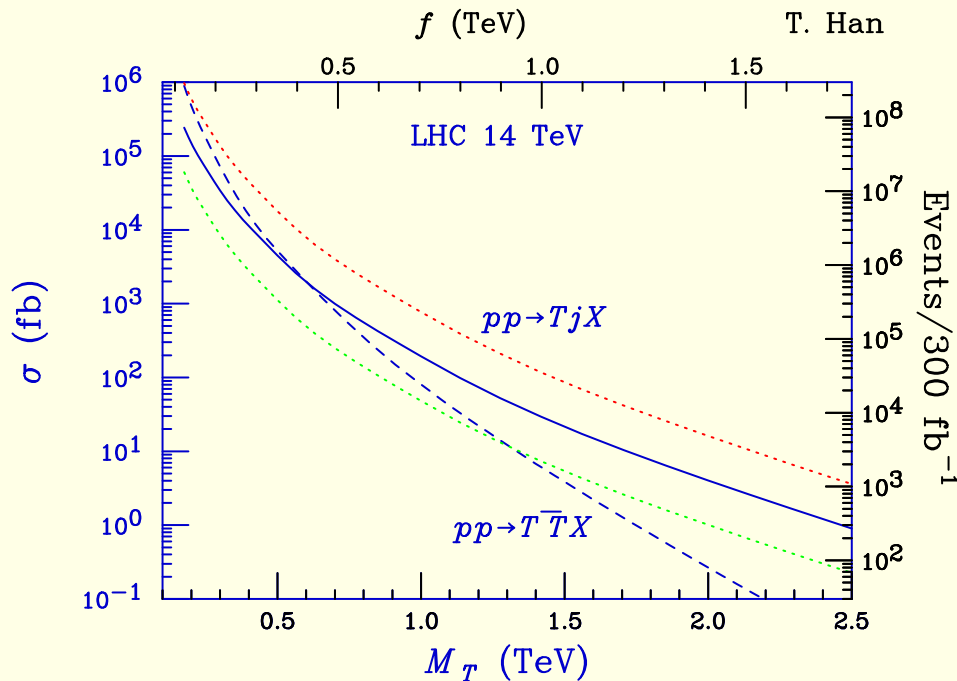


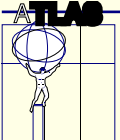
Figure from Han

Single production dominates at large masses

Three single production curves are for $\lambda_1/\lambda_2 = 2, 1, 0.5$

Width is small

Single Production is used in the following: note recoil jet.



$$T \rightarrow Zt$$

Reconstruct from $Z \rightarrow l^+l^-$ and $t \rightarrow bl\nu$

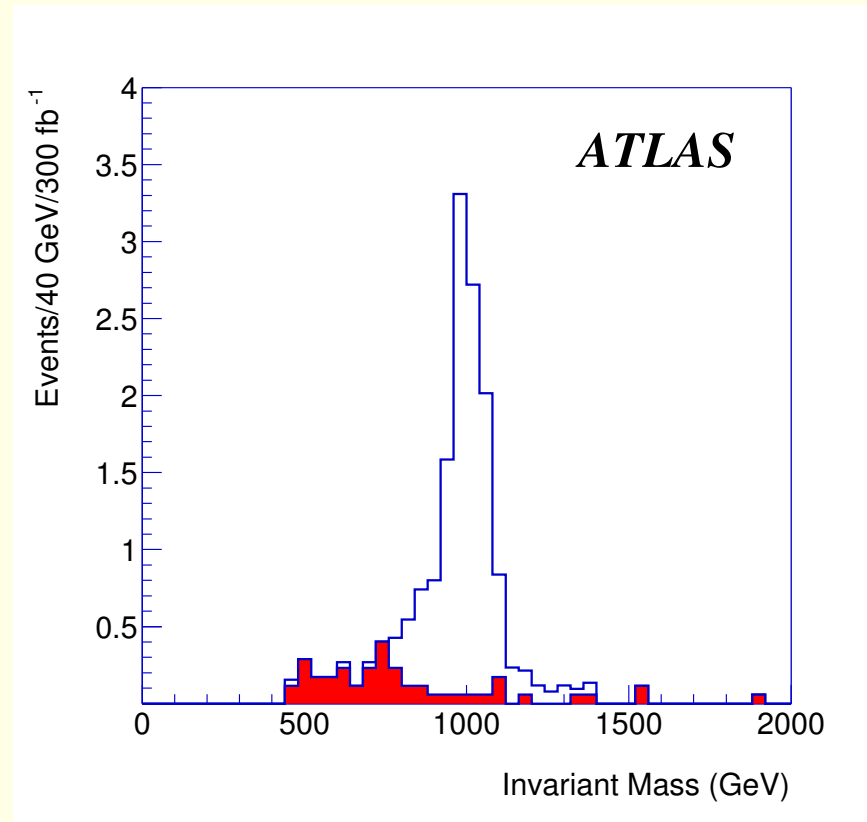
Three isolated leptons (either e or μ) with $p_T > 40$ GeV and $|\eta| < 2.5$ one of which has $p_T > 100$ GeV

No other leptons with $p_T > 15$ GeV

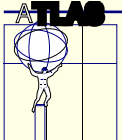
One pair of leptons within 10 GeV of Z mass.

$\cancel{E}_T > 100$ GeV

At least one tagged b -jet with $p_T > 30$ GeV



Background is dominated by tbZ



$$T \rightarrow Wb$$

Reconstruct from $T \rightarrow b\ell\nu$

One isolated lepton (either e or μ) with $p_T > 100$ GeV and $|\eta| < 2.5$

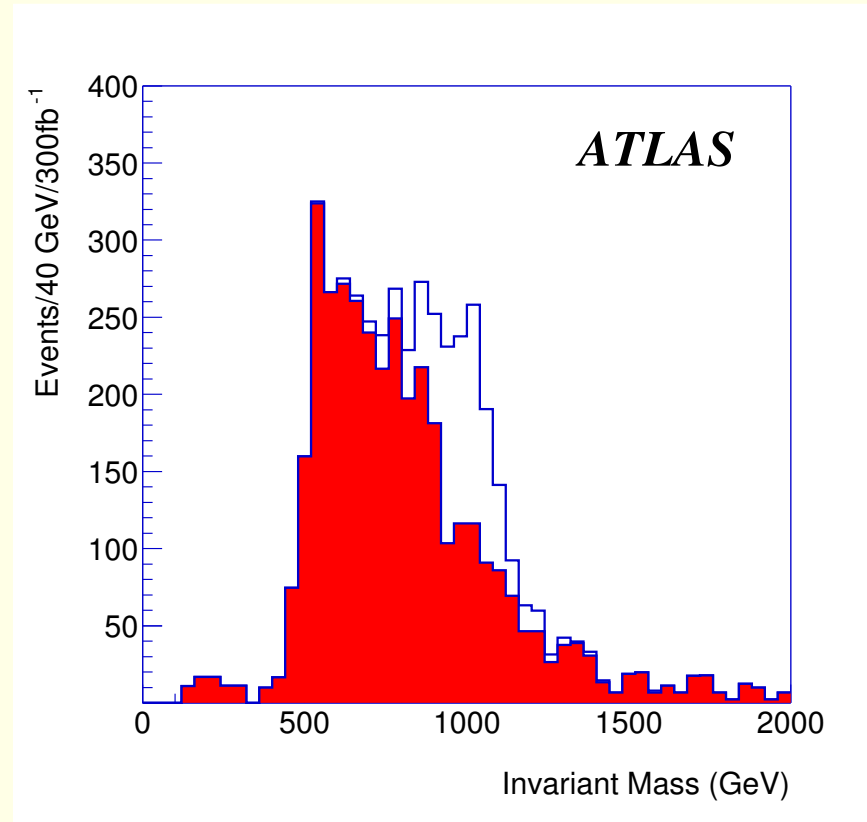
No other leptons with $p_T > 15$ GeV

No more than 2 jets with $p_T > 50$ GeV and

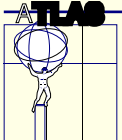
$M(j1, j2) > 200$ GeV

$\cancel{E}_T > 100$ GeV

at least one tagged b -jet with $p_T > 200$ GeV



Background is dominated by $t\bar{t}$



$$T \rightarrow ht$$

Reconstruct from $h \rightarrow b\bar{b}$ and $t \rightarrow bl\nu$

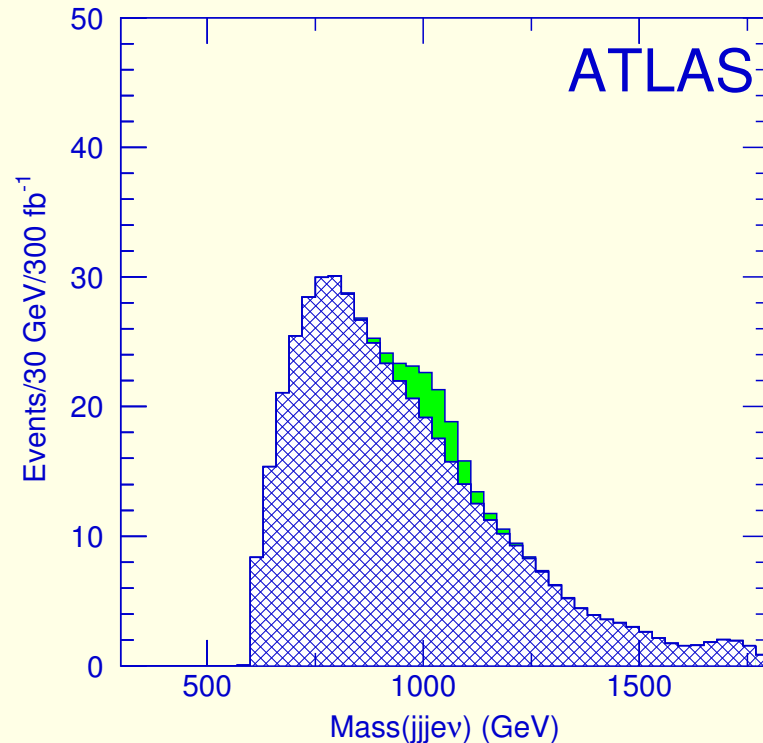
One isolated e or μ with $p_T > 100$ GeV
and $|\eta| < 2.5$.

Three jets with $p_T > 130$ GeV.

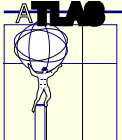
Four jets with $p_T > 15$ GeV.

At least one jet tagged as b -jet

Mass of dijet system within 20 GeV of
Higgs mass (assumed to be 120 GeV)



Background dominated by $t\bar{t}$

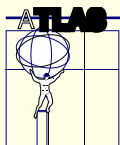


New Bosons

Expect two neutral and two charged: Z_H, A_H, W_H^\pm

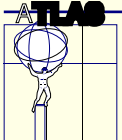
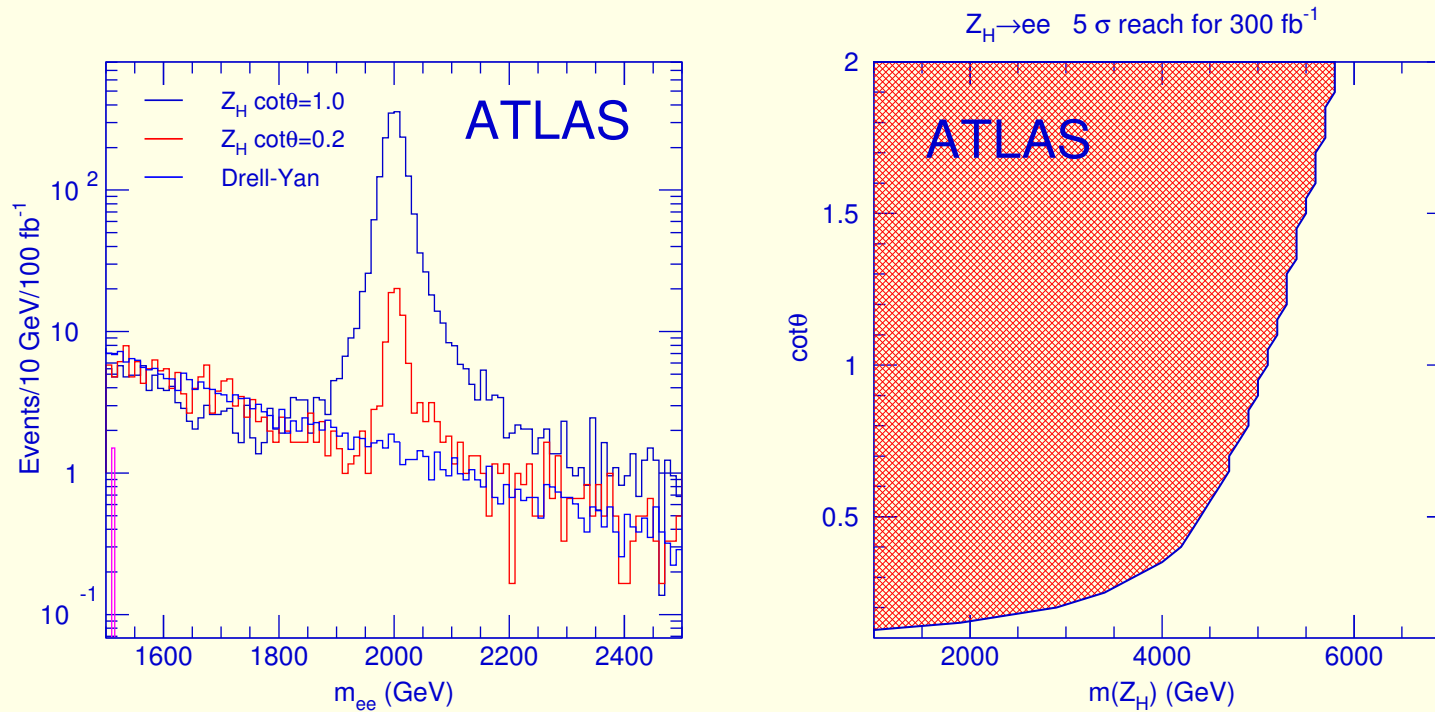
Model has two additional couplings corresponding to the extra $SU(2) \times U(1)$,

Bosons will be discovered via leptonic decays **But critical test is cascades such as**
 $Z_H \rightarrow Zh$



New Bosons – Leptonic decays

Clear signal over Drell-Yan background. Plot shows 2 TeV mass for Z_H



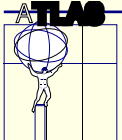
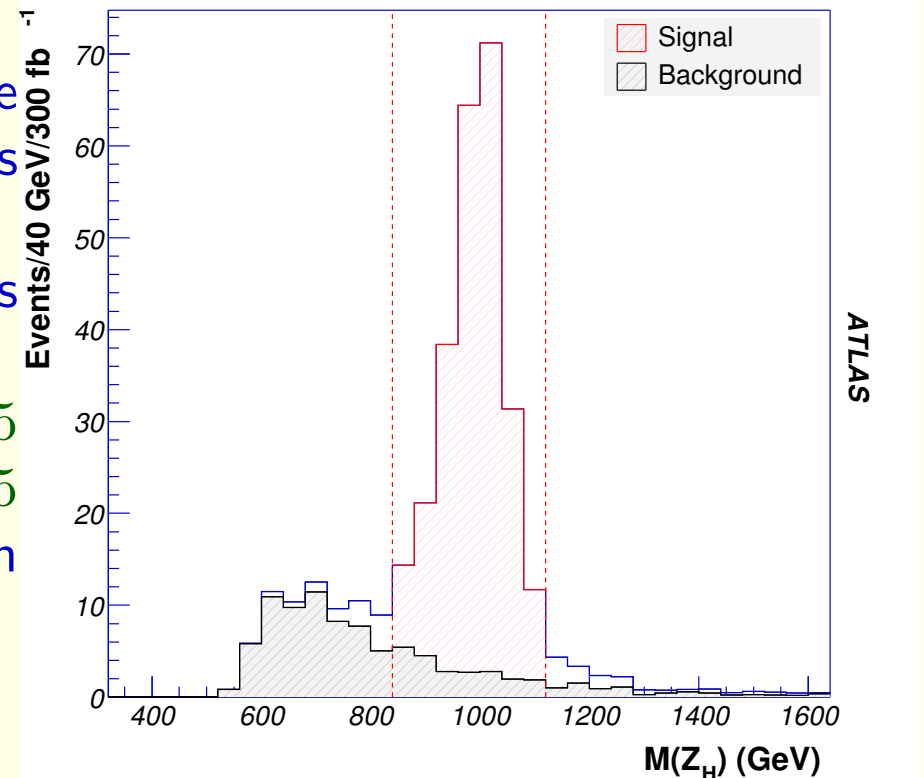
New Bosons – Cascade decay $Z_H \rightarrow Zh \rightarrow \ell^+ \ell^- b\bar{b}$

Two leptons of opposite charge and same flavor with $p_T > 6(5)$ GeV for muons (electrons) and $|\eta| < 2.5$

The lepton pair should have a mass between 76 and 116 GeV

Two reconstructed b – jets with $p_T > 25$ and $|\eta| < 2.5$, which are within $\Delta R < 1.5$

The b – jet pair should have a mass between 60 and 180 GeV



$$Z_H \rightarrow Zh, h \rightarrow \gamma\gamma$$

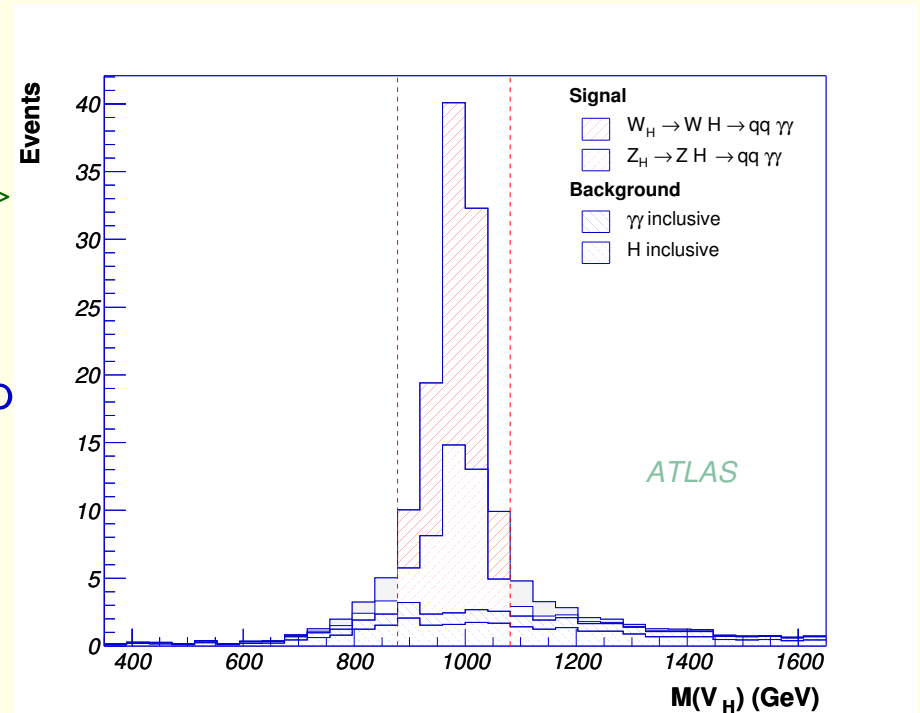
Must use all hadronic mode of Z : Cannot distinguish W_H from Z_H

Two isolated photons one having $p_T(1) > 25$ GeV, $p_T(2) > 40$ GeV.

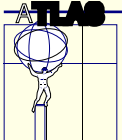
$$M(\gamma\gamma) = m_h \pm 2\sigma$$

The jet pair with invariant mass closest to M_W is selected.

Pair has a combined $p_T > 200$ GeV



Can also extract signal via Jacobian peak in the P_T dist of Higgs



Extra Higgs

ϕ^{++} produced by WW fusion: So must use the forward tagging jets

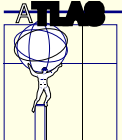
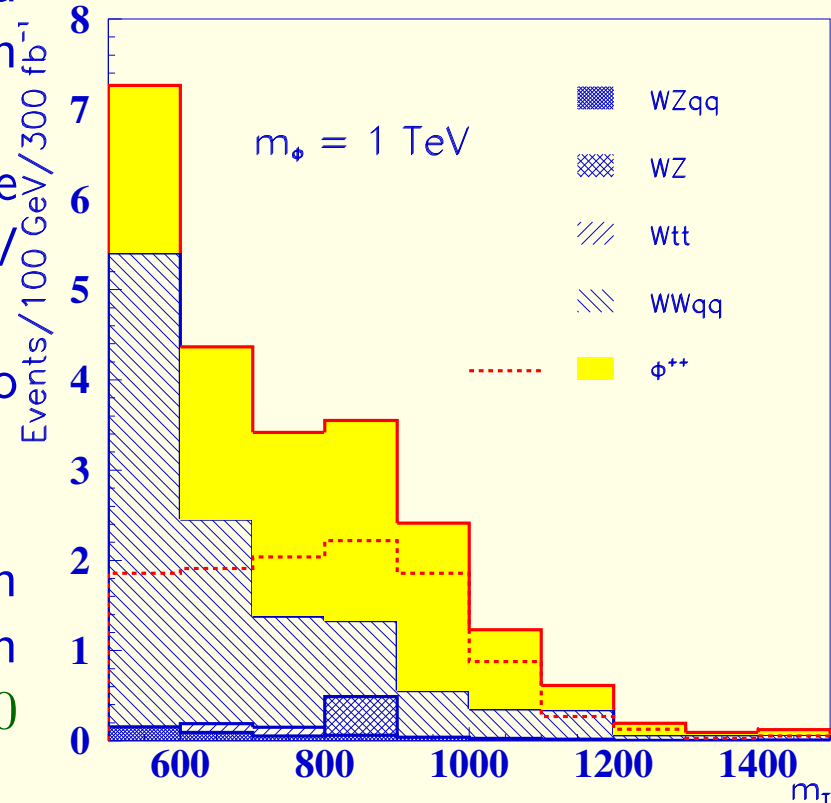
Two reconstructed positively charged isolated leptons (electrons or muons) with $|\eta| < 2.5$

One of the leptons was required to have $p_T > 150$ GeV and the other $p_T > 20$ GeV
 $|p_{T1} - p_{T2}| > 200$ GeV

the difference in pseudorapidity of the two leptons $|\eta_1 - \eta_2| < 2$.

$\cancel{E}_T > 50$ GeV

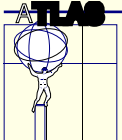
Two jets each with $p_T > 15$ GeV, with rapidities of opposite sign, separated in rapidity $|\eta_1 - \eta_2| > 5$; one jet has $E > 200$ GeV and the other $E > 100$ GeV



Summary of sensitivity

- T Observable in both $h(120)t$ (up to mass of 1.2 TeV) and Zt (up to mass 1.0 TeV): Wb is observable up to 1.3 TeV for $\lambda_1/\lambda_2 = 1$
- Z_H observable in e^+e^- to mass of 4.5 TeV for $\cot\theta = 0.5$
 $Z_H \rightarrow Zh(120) \rightarrow Zb\bar{b}$ observable for mass up to 2 TeV
 $Z_H \rightarrow Zh(120) \rightarrow Z\gamma\gamma$ observable for masses up to 1.1 TeV
- ϕ^{++} may be observable in W^+W^+ at 1.5 TeV
- More work needed for $m_h \gtrsim 150$ GeV

LHC finds it or motivation disappears



Hadron Production of Sparticles

LHC is likely to be above threshold for many sparticles

A consistent model must be used for simulation. Most popular is SUGRA

Unification all scalar masses (m_0) at GUT scale

Unification all gaugino masses ($m_{1/2}$) at GUT scale

Universal A and B

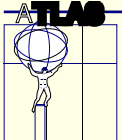
$|\mu|$ and B are traded off for M_Z and $\tan\beta = v_1/v_2$

So **five** parameters $\tan\beta = v_1/v_2$ $\text{sign}(\mu)$ A , $m_{1/2}$ and m_0 gives full mass spectrum and decays

Glino mass strongly correlates with $m_{1/2}$, slepton mass with m_0 .

Studies have also been done for Gauge, or Anomaly mediated models.

Enough cases have now been studied that given a complete set of masses and decay rates, we can usually estimate what can be done at LHC.



SUSY in hadron colliders

Inclusive signatures provide evidence up to 2.5 TeV for squarks and gluinos.

Everything is produced at once; squarks and gluinos have largest rates.

Production of Sparticles with only E-W couplings (e.g sleptons, Higgs) may be dominated by decays not direct production.

Must use a consistent model for simulation: cannot discuss one sparticle in isolation.

Makes studies somewhat complicated and general conclusions difficult to draw.

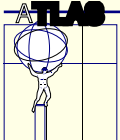
Studies shown here are not optimized

Large event rates are used to cut hard to get rid of standard model background.

Dominant backgrounds are combinatorial from SUSY events themselves.

Studies shown here are not optimized; large event rates are exploited to cut hard to get rid of standard model background.

Full program difficult to estimate, depends on masses and branching ratios



Inclusive analysis at LHC

These studies tend to be conservative

Reach is shown for various inclusive signals

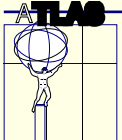
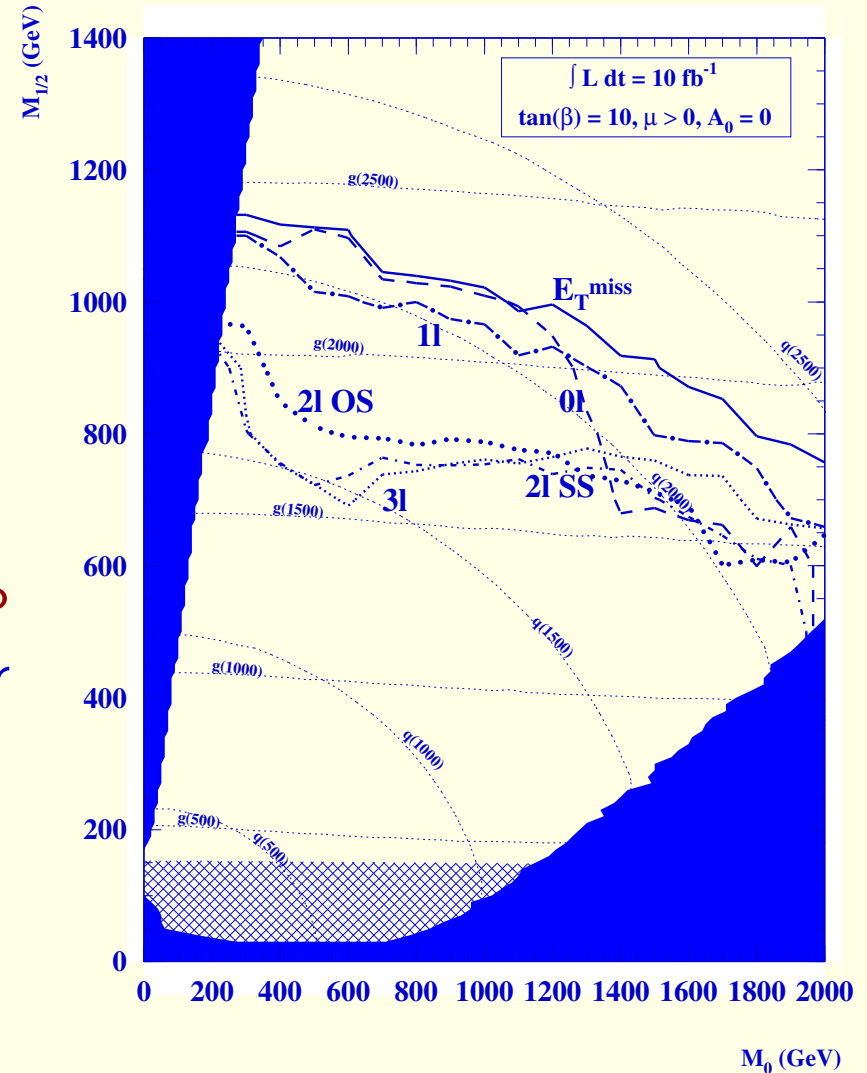
Jets plus missing E_T

Multileptons of same and opposite sign

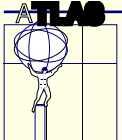
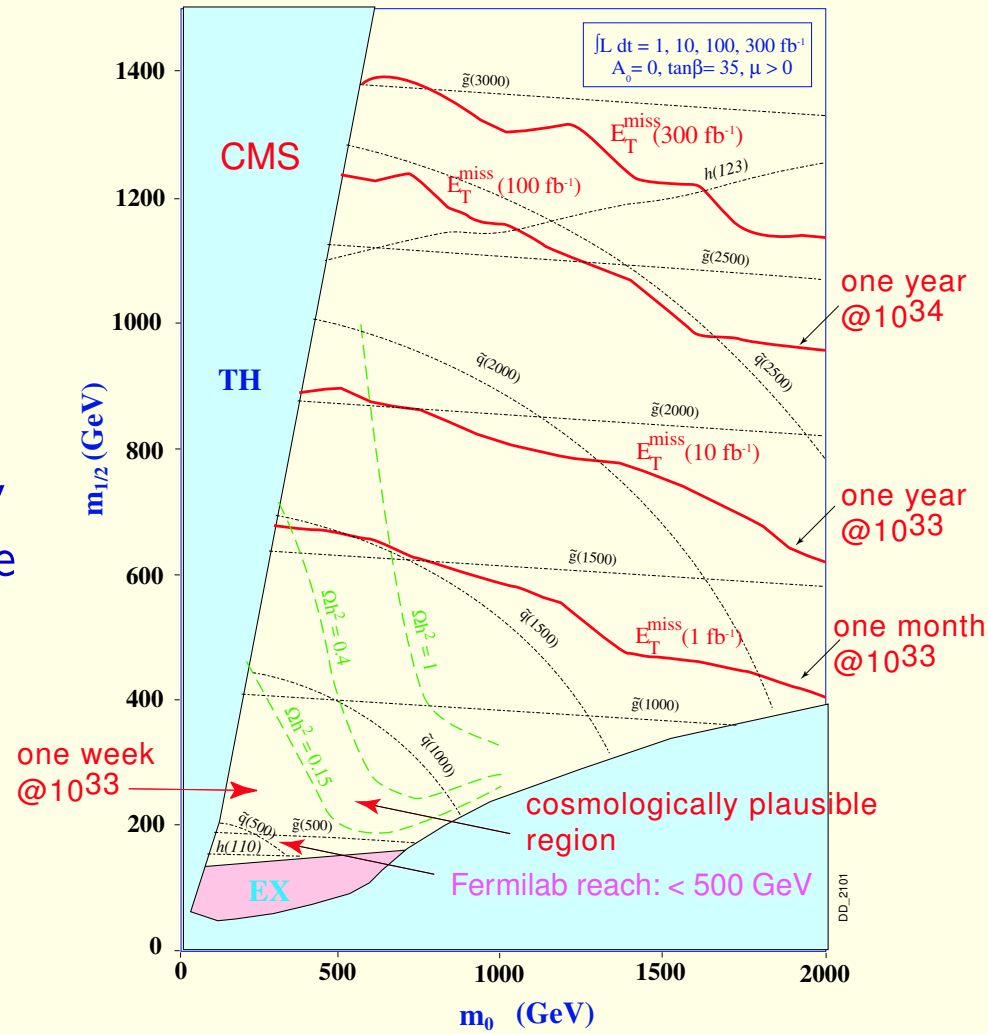
Shown for SUGRA

Shaded regions excluded by theory or LEP

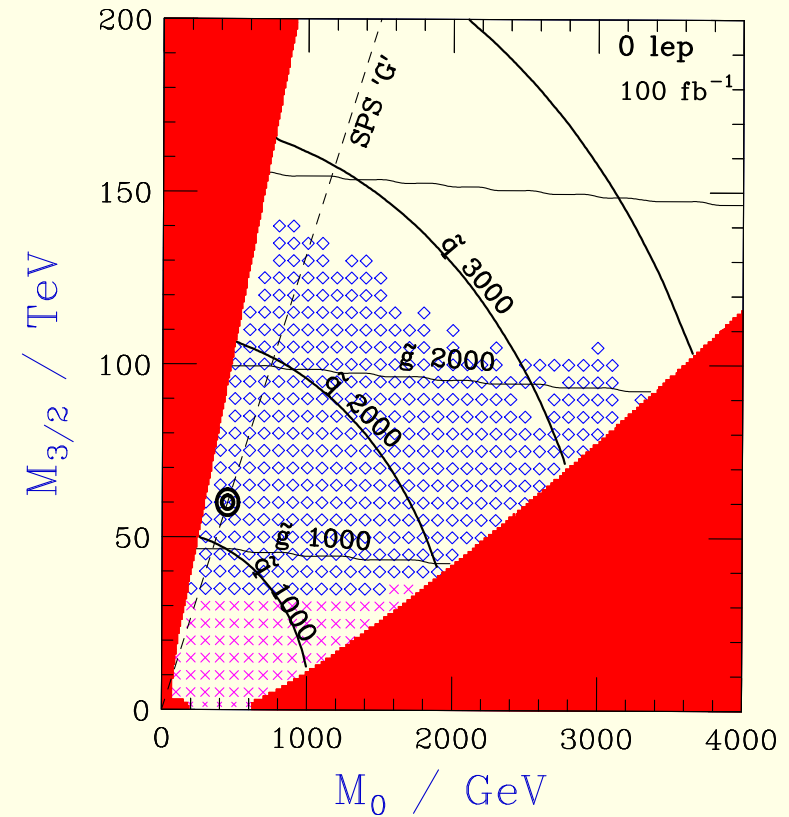
Extends to gluino masses of over 2 TeV for 10fb^{-1}



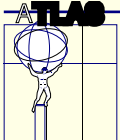
Plot shows evolution of reach with luminosity
 Notice that a few 0.1fb^{-1} covers most of the region favored by fine tuning arguments



Reach is similar in other models
 Example of anomaly mediated model
 Shaded pink region is excluded by LEP



In general reach depends mainly on $M_{\tilde{g}}$ and $M_{\tilde{q}}$ provided $M_{\tilde{\chi}_1^0} \ll M_{\tilde{g}}, M_{\tilde{q}}$
 rather model independent

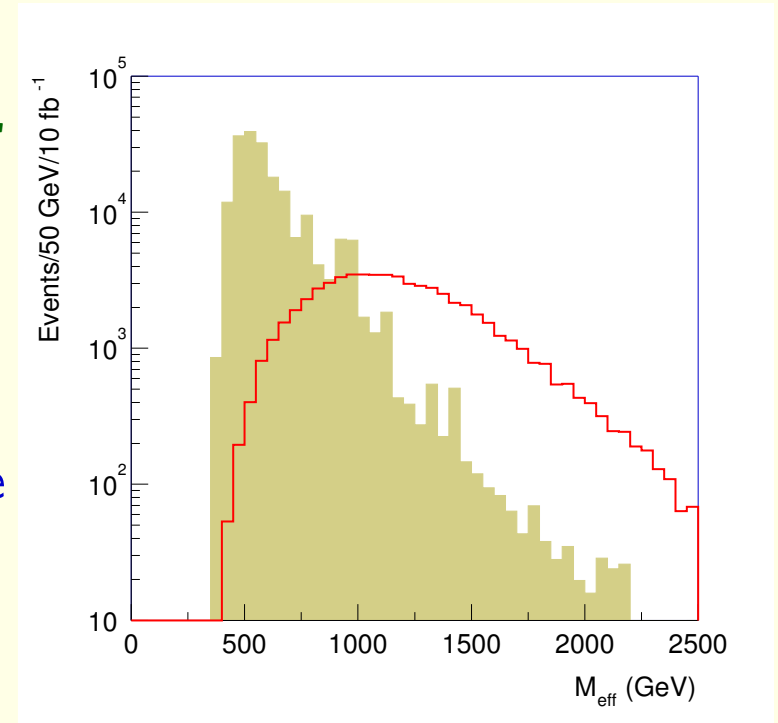


Estimating the scale

Select events with at least 4 jets and Missing E_T
A simple variable

$$M_{\text{eff}} = P_{t,1} + P_{t,2} + P_{t,3} + P_{t,4} + \cancel{E}_T$$

At high M_{eff} non-SM signal rises above background
note scale

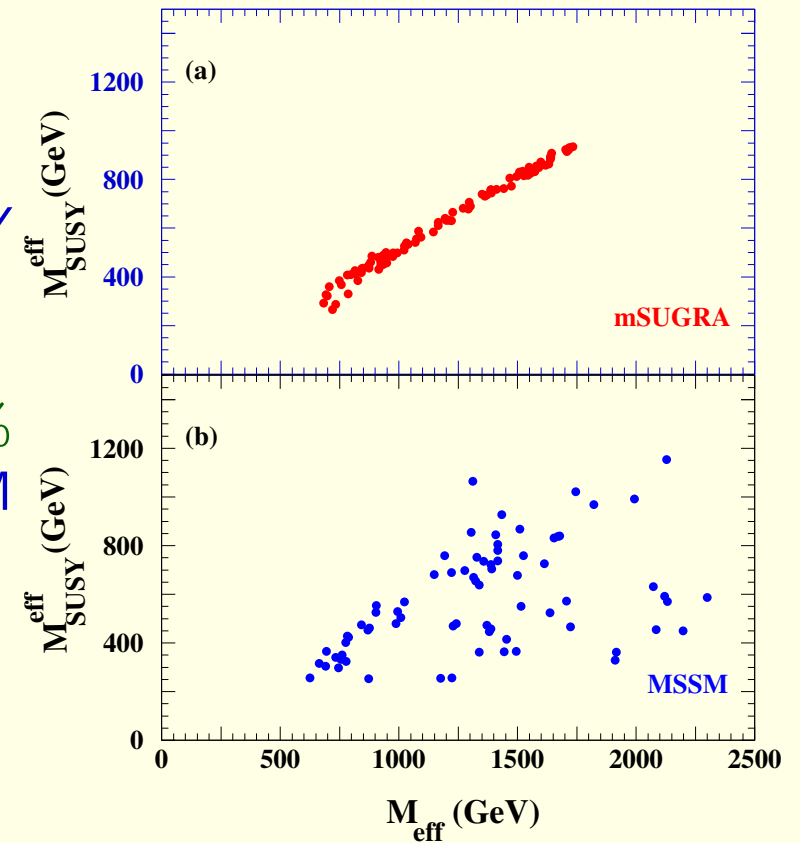


Peak in M_{eff} distribution correlates with SUSY mass scale

$$M_{\text{SUSY}} = \min(M_{\tilde{u}}, M_{\tilde{g}})$$

Will determine gluino/squark masses to $\sim 15\%$ in SUGRA, much poorer in a more general MSSM

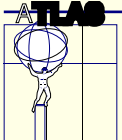
15 parameters were varied



Note that rate information is difficult to use as BR are not known

Must reconstruct decays to get more information

Examples follow



Identifying typical decays

Assume $M_{\tilde{g}} > M_{\tilde{q}}$ (similar results in reverse case)

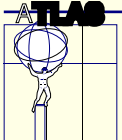
Then typically

$$B(\tilde{q}_L \rightarrow \tilde{\chi}_2^0 q) \sim 1/3, \quad B(\tilde{q}_L \rightarrow \tilde{\chi}_1^\pm q') \sim 2/3, \quad B(\tilde{q}_R \rightarrow \tilde{\chi}_1^0 q) \sim 1.$$

If channels are open, two body decays such as $\tilde{\chi}_2^0 \rightarrow \tilde{\ell}^+ \ell^-$, $\tilde{\chi}_2^0 \rightarrow Z \tilde{\chi}_1^0$, $\tilde{\chi}_2^0 \rightarrow h \tilde{\chi}_1^0$ usually dominate

Otherwise $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \ell^+ \ell^-$ via virtual slepton

So a good idea to look for leptons



Leptonic final states

Isolated leptons indicate presence of t , W , Z , weak gauginos or sleptons

Straightforward case

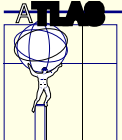
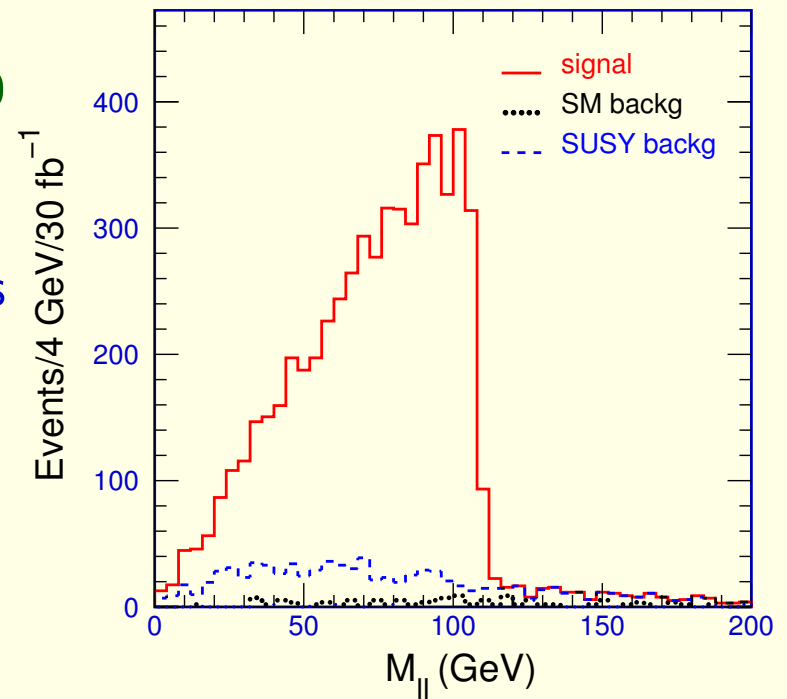
Decay chain is $\tilde{\chi}_2 \rightarrow \tilde{\ell}^+ \ell^- \rightarrow \tilde{\chi}_1 \ell^+ \ell^-$

- 2 isolated opposite sign leptons; $p_t > 10$ GeV
- ≥ 4 jets; one has $p_t > 100$ GeV, rest $p_t > 50$ GeV
- $\cancel{E}_T > \max(100, 0.2M_{eff})$

Mass of opposite sign same flavor leptons is constrained by decay

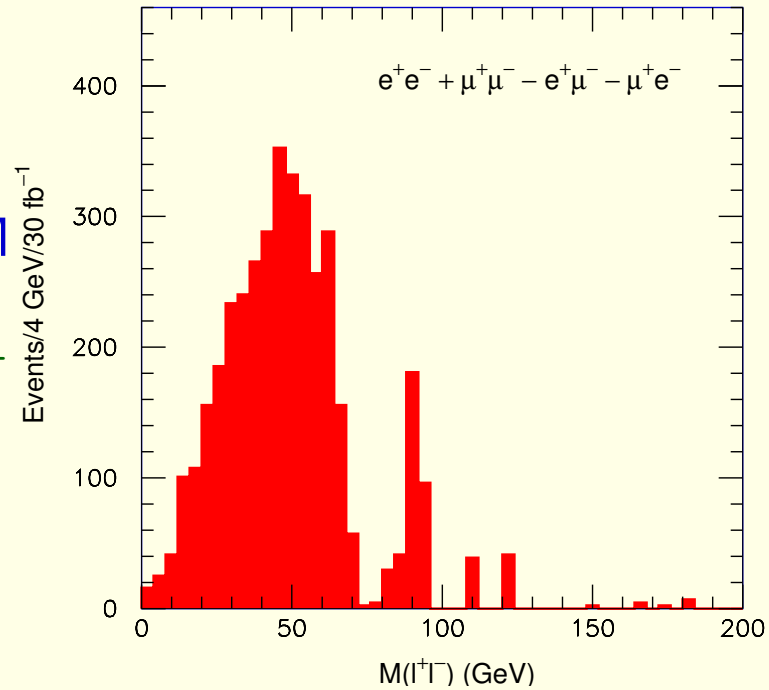
$$M_{\ell\ell} = \sqrt{(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{\ell}}^2)(M_{\tilde{\ell}}^2 - M_{\tilde{\chi}_1^0}^2)}/M_{\tilde{\ell}}.$$

Standard Model background is dominated by $t\bar{t}$
Other SUSY events (mainly $\tilde{\chi}_1^\pm$ decays also contribute)

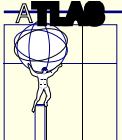


Flavor subtraction remove the SM background and cleans up signal

This example has both $\tilde{\chi}_2^0 \rightarrow \tilde{\ell}^+ \ell^-$ and $\tilde{\chi}_2^0 \rightarrow Z \tilde{\chi}_1^0$,



Must add jets to this to try to get full decay chains

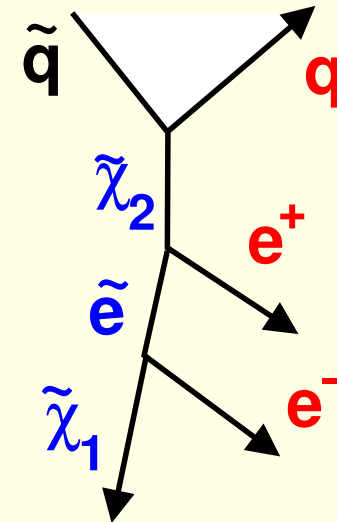


Squark masses

Attempt to find $\tilde{q}_L \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\ell}\ell \rightarrow q\ell\ell\tilde{\chi}_1^0$

Identify and measure decay chain

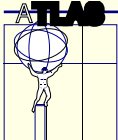
- 2 isolated opposite sign leptons; $p_t > 10$ GeV
- ≥ 4 jets; one has $p_t > 100$ GeV, rest $p_t > 50$ GeV
- $\cancel{E}_T > \max(100, 0.2M_{eff})$

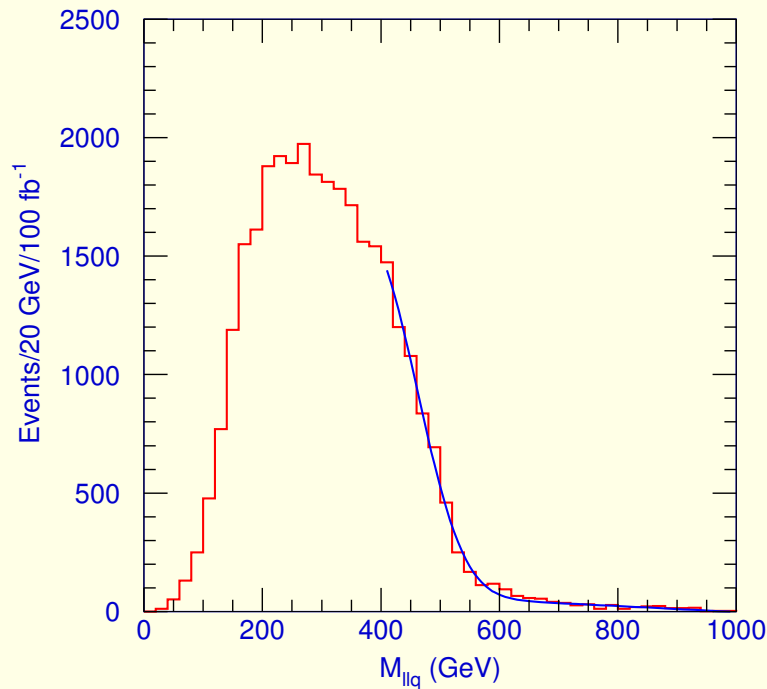


Mass of $q\ell\ell$ system has max at

$$M_{\ell\ell q}^{\max} = \left[\frac{(M_{\tilde{q}_L}^2 - M_{\tilde{\chi}_2^0}^2)(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{\chi}_1^0}^2)}{M_{\tilde{\chi}_2^0}^2} \right]^{1/2} = 552.4 \text{ GeV}$$

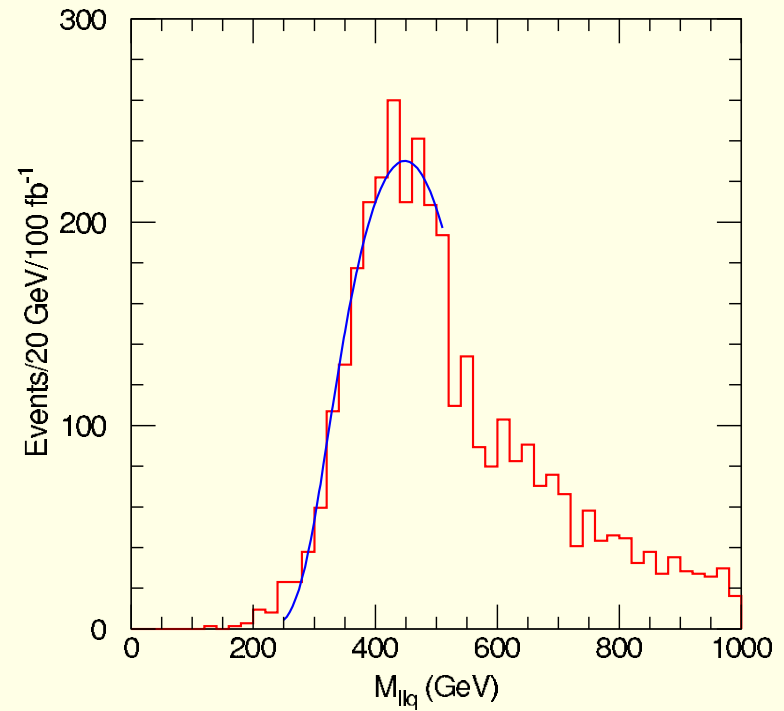
and min at 271 GeV (in the example shown)



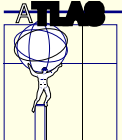


smallest mass of possible *lljet* combinations

Kinematic structure clearly seen
Can also exploit *ljjet* mass



largest mass of possible *lljet* combinations

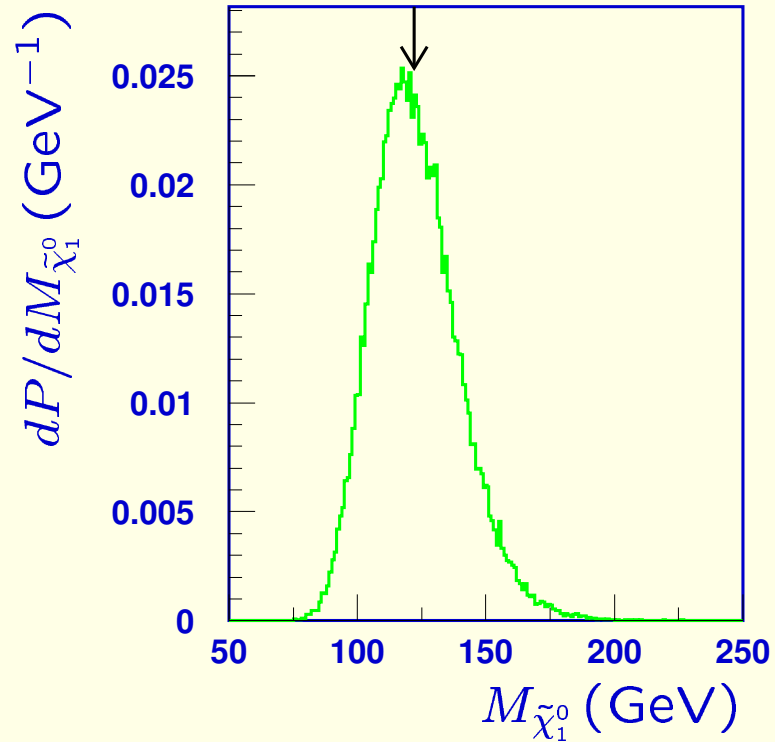
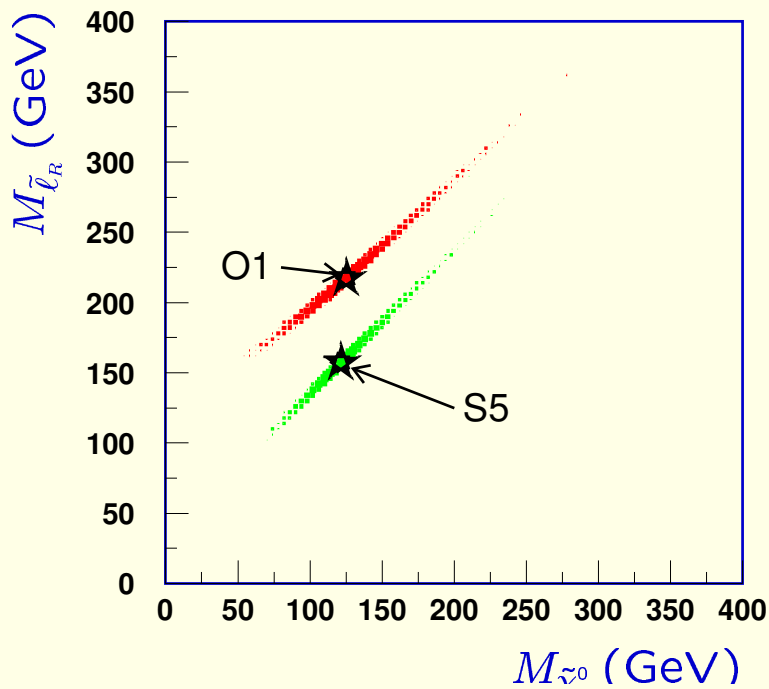


Can now solve for the masses. Note that no model is needed

Very naive analysis has 4 constraints from $lq, llq_{upper}, llq_{lower}, ll$ masses

4 Unknowns, $m_{\tilde{q}_L}, m_{\tilde{e}_R}, m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_1^0}$

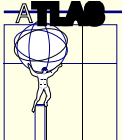
Errors are 3%, 9%, 6% and 12% respectively



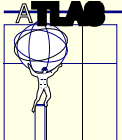
correlations $m_{\tilde{e}_R}$ vs. $m_{\tilde{\chi}_1^0}$

LSP mass

Mass of unobserved LSP is determined



Errors are strongly correlated and a precise independent determination of one mass reduces the errors on the rest.



What about \tilde{q}_R ?

$\tilde{q}_r \tilde{q}_r \rightarrow qq\tilde{\chi}_1^0\tilde{\chi}_1^0$ produces clean events

$$m_{T2}^2(\chi) \equiv \min_{q_T^{(1)} + q_T^{(2)} = \cancel{E}_T} \left[\max \left\{ m_T^2(p_T^{j(1)}, q_T^{(1)}; \chi), m_T^2(p_T^{j(2)}, q_T^{(2)}; \chi) \right\} \right]$$

Event selection

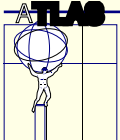
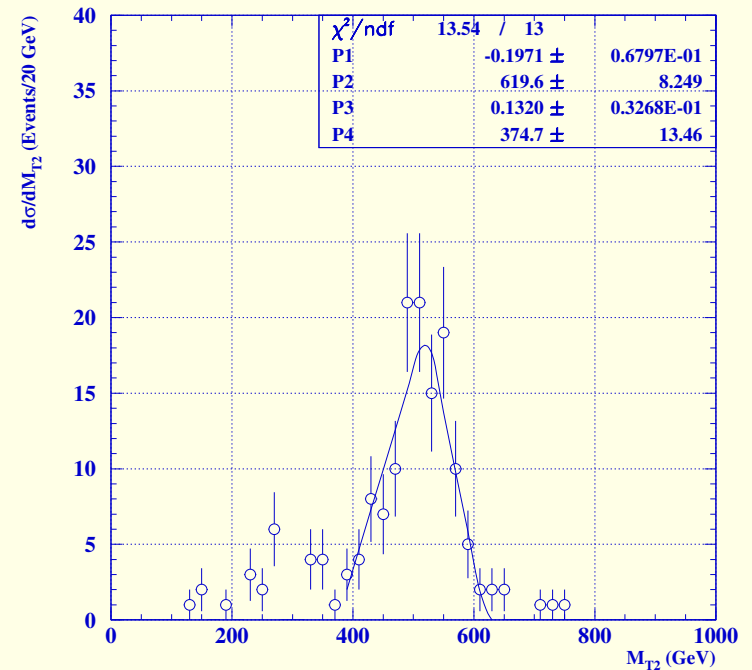
Two jets with $P_T > 150$ GeV

$\cancel{E}_T > 200$ GeV

No other jets with $P_T > 40$ GeV

Clear structure

Determines a combination of M_{q_r} and $M_{\tilde{\chi}_1^0}$



Decays to Higgs

If $\chi_2^0 \rightarrow \chi_1^0 h$ exists then this final state followed by $h \rightarrow b\bar{b}$ results in discovery of Higgs at LHC.

In these cases $\sim 20\%$ of SUSY events contain $h \rightarrow b\bar{b}$

Event selection

$\cancel{E}_T > 300$ GeV

≥ 2 jets with $p_T > 100$ GeV and ≥ 1 with

$|\eta| < 2$

No isolated leptons (suppresses $t\bar{t}$)

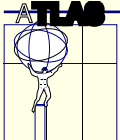
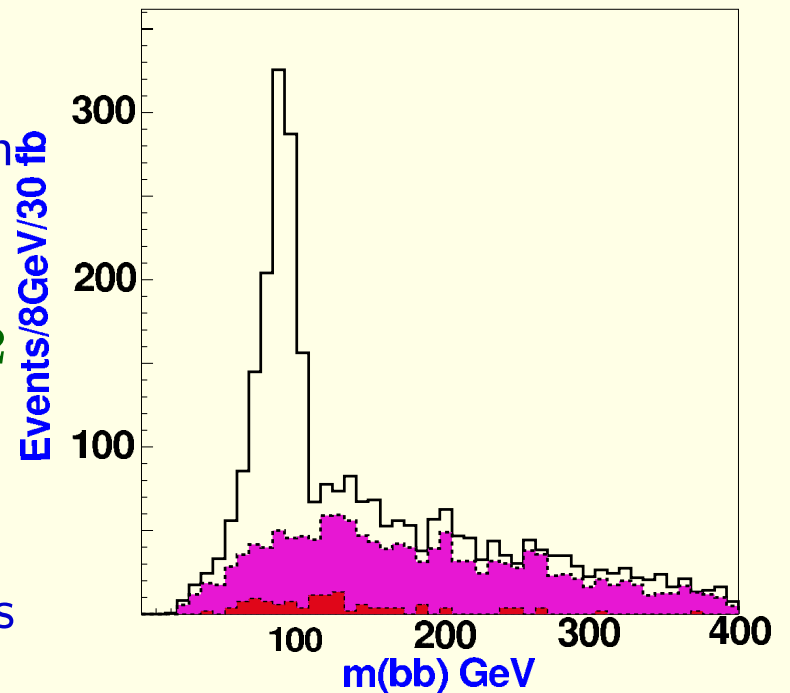
Only 2 b-jets with $p_{T,b} > 55$ GeV and $|\eta| < 2$

$\Delta R_{b\bar{b}} < 1.0$ (suppresses $t\bar{t}$)

Clear peak in $b\bar{b}$ mass

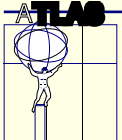
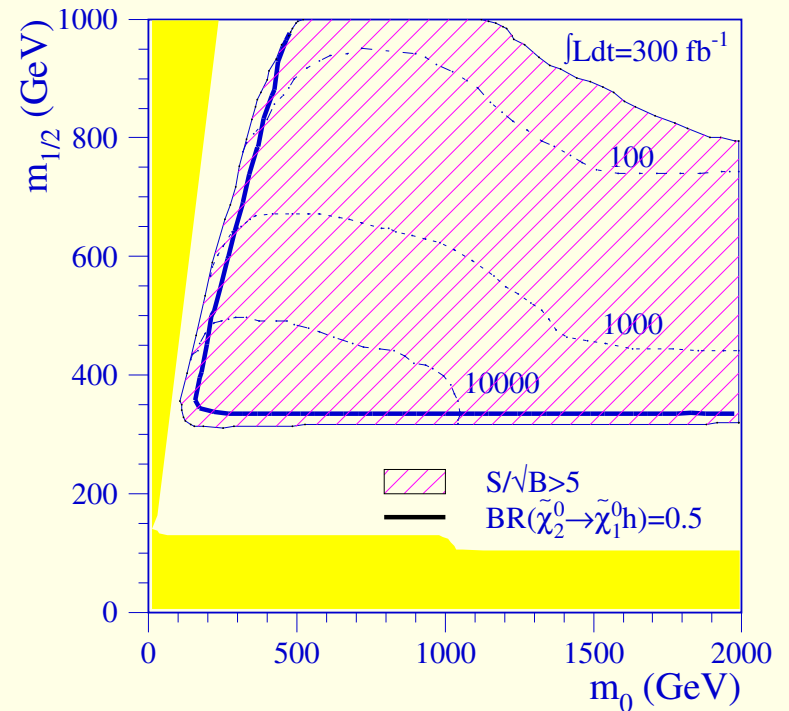
Very small standard model background (pale)

Dominant background is other SUSY decays (dark)



Generally applicable

This method works over a large region of parameter space in the SUGRA Model
Hatched region has $S/\sqrt{B} > 5$
Contours show number of reconstructed Higgs
Channel is closed at low $m_{1/2}$

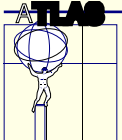
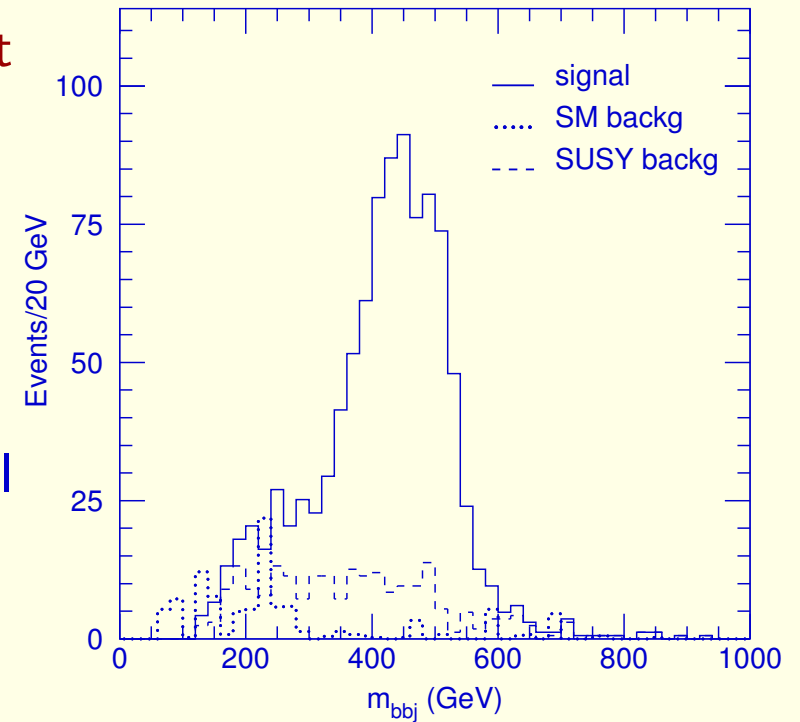


Combine with a jet to attempt to get
 $\tilde{q} \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\chi}_1^0$

Take $b\bar{b}$ around the peak and combine with all jets

Plot the combination with the smallest mass

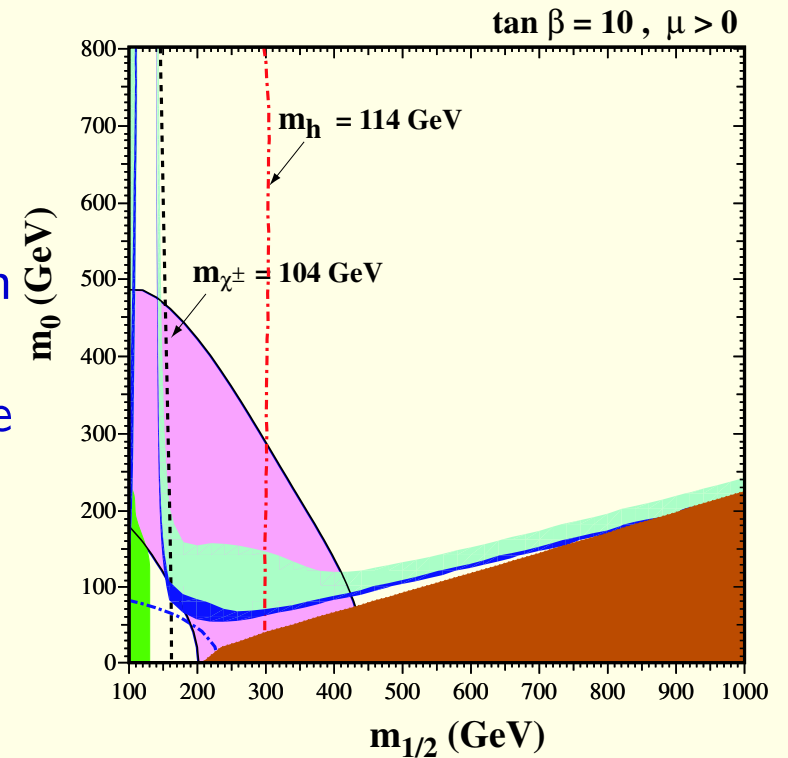
Again we see upper kinematic limit



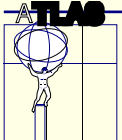
Preferred regions?

It would be nice to know where to look

If we really believe in minimal SUGRA then WMAP provides strong constraints
Even stronger if $g - 2$ is included (with one value of $R(e^+e^-)$ at low energy)

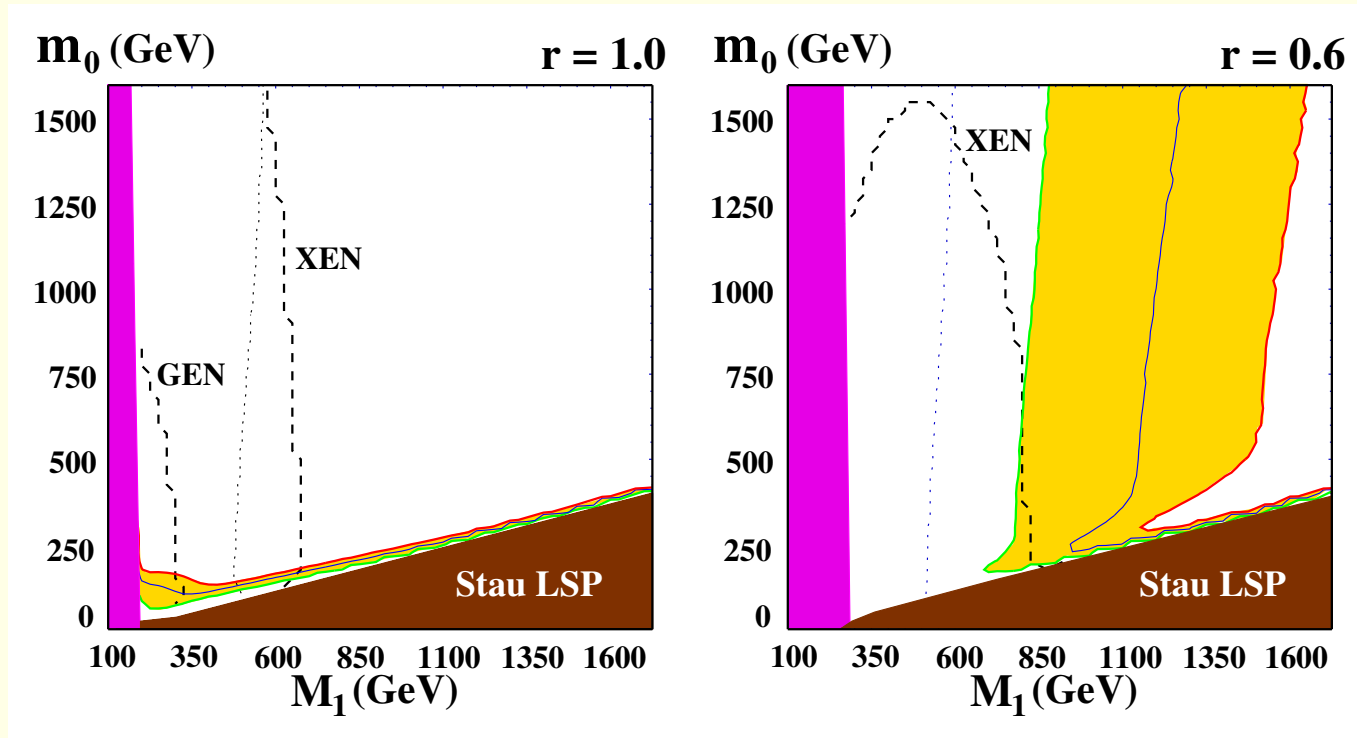


Plot from Ellis, Olive

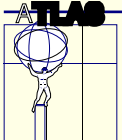


But constraints weaken outside minimal SUGRA

$R = M_2/M_3$ at GUT scale.



Plot from Birkdahl-Hansen et al



Extra Dimensions

Many theories (*e.g.* string) predict extra dimensions of size R

What is R ?. Old ideas $\Rightarrow 1/M_P$. Unobservable.

Larger value of R can allow scale of Gravity to be smaller

Arkani-Hamed...

$$G_N = 8\pi R^\delta M_D^{-(2+\delta)}$$

$$M_D \sim 1 \text{ TeV} \quad R \sim 10^{32/\delta-16} \text{ mm}$$

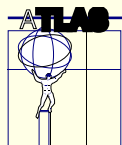
Attractive because no hierarchy between M_W and M_D

But hierarchy between $1/R$ and M_W still exists

Compactified dimension implies tower of states with $\Delta m \sim 1/R$

\Rightarrow Standard Model fields must be stuck in $d = 4$ But many graviton (G) excitations can exist.

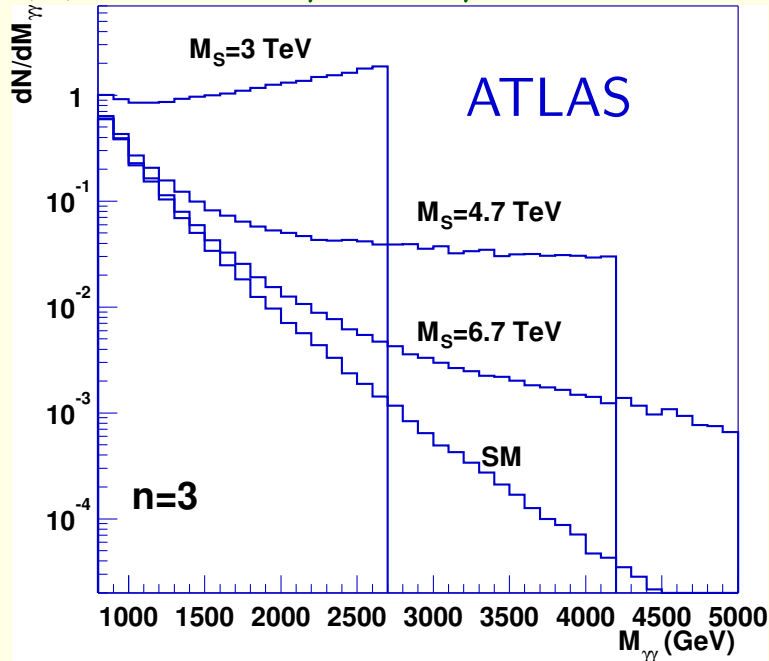
In simplest models processes such as $qg \rightarrow qG$ or $q\bar{q} \rightarrow \gamma G$ give missing energy signatures or distortions in rates due to exchanges



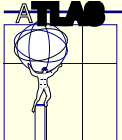
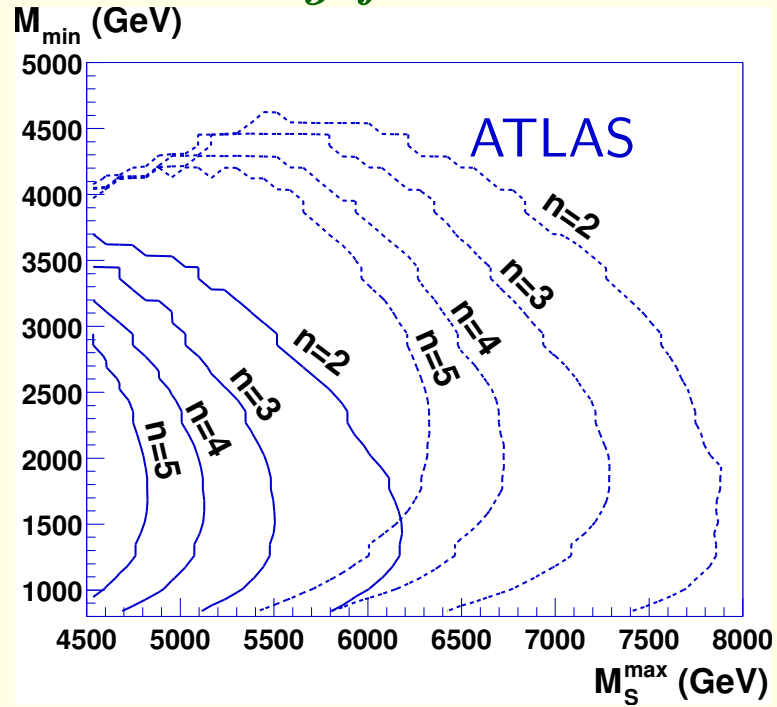
Studies have focused on jets + \cancel{E}_T , γ + \cancel{E}_T , $\gamma\gamma$, and ll final states.

Virtual effects from graviton exchange show up as excesses in the production rates

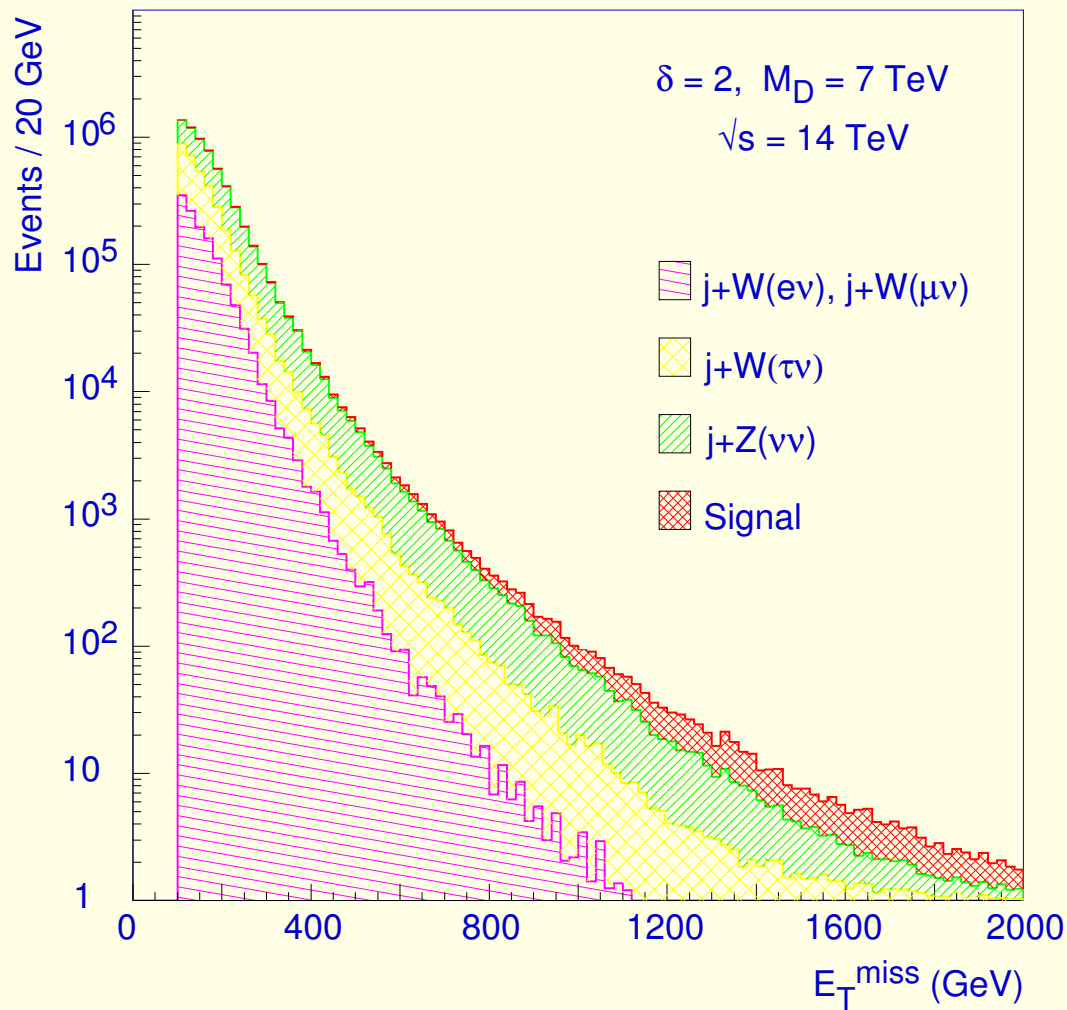
$\gamma\gamma$ Events/GeV/10 fb⁻¹



Sensitivity for 10 and 100 fb⁻¹

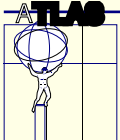


ADD extra dimensions produce jets + \cancel{E}_T , γ + \cancel{E}_T signals from graviton emission



red region is signal from jets for
 100 fb^{-1}
 Sensitivity

δ	M_D^{max} (TeV)
2	9
3	7
4	6



Warped Extra Dimensions – Randall Sundrum models

Model of 5-dim space with two branes of 4-dim. SM fields are stuck on one brane. Metric is “non-factorizable”

$$ds^2 = e^{-kR\phi} \eta_{\mu,\nu} dx^\mu dx^\nu + R^2 d\phi^2$$

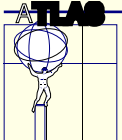
Scale $\Lambda = ke^{-kR\pi}$ in 4-D world

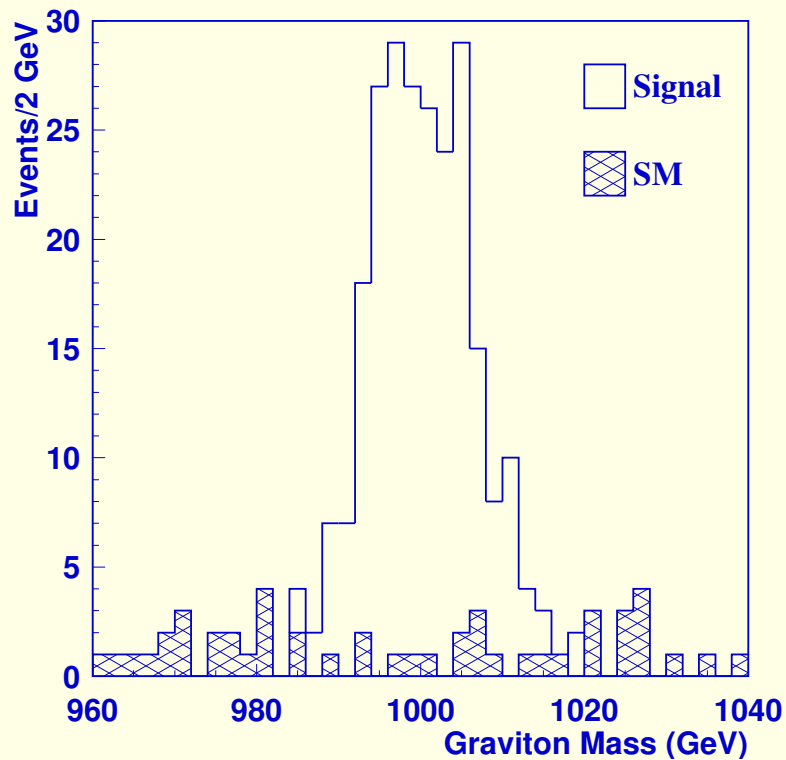
Can get $\Lambda \sim 1$ TeV with $Rk \sim 12$ and $k \sim M_P$

Graviton excited states have mass gaps of order Λ

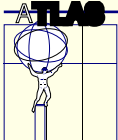
Properties are determined by k/M_P .

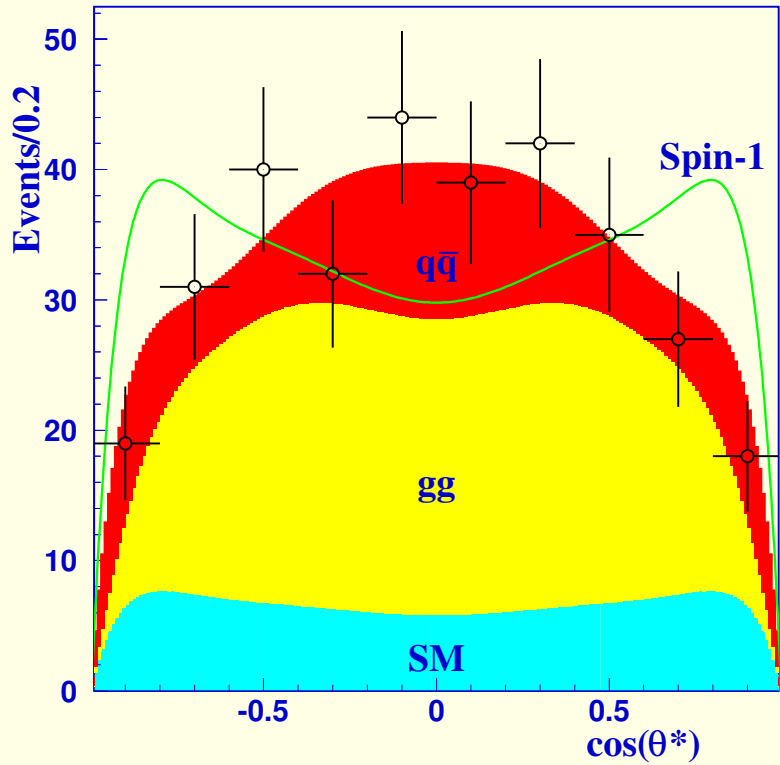
Simple models have $k/M_P \sim 0.01$; excited states are then narrow and weakly coupled



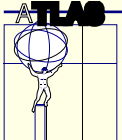


Look for a resonance in dilepton final states
e.g. $gg \rightarrow e^+e^-$ Discovery limit is \sim
 $1.8TeV$ for 100fb^{-1}



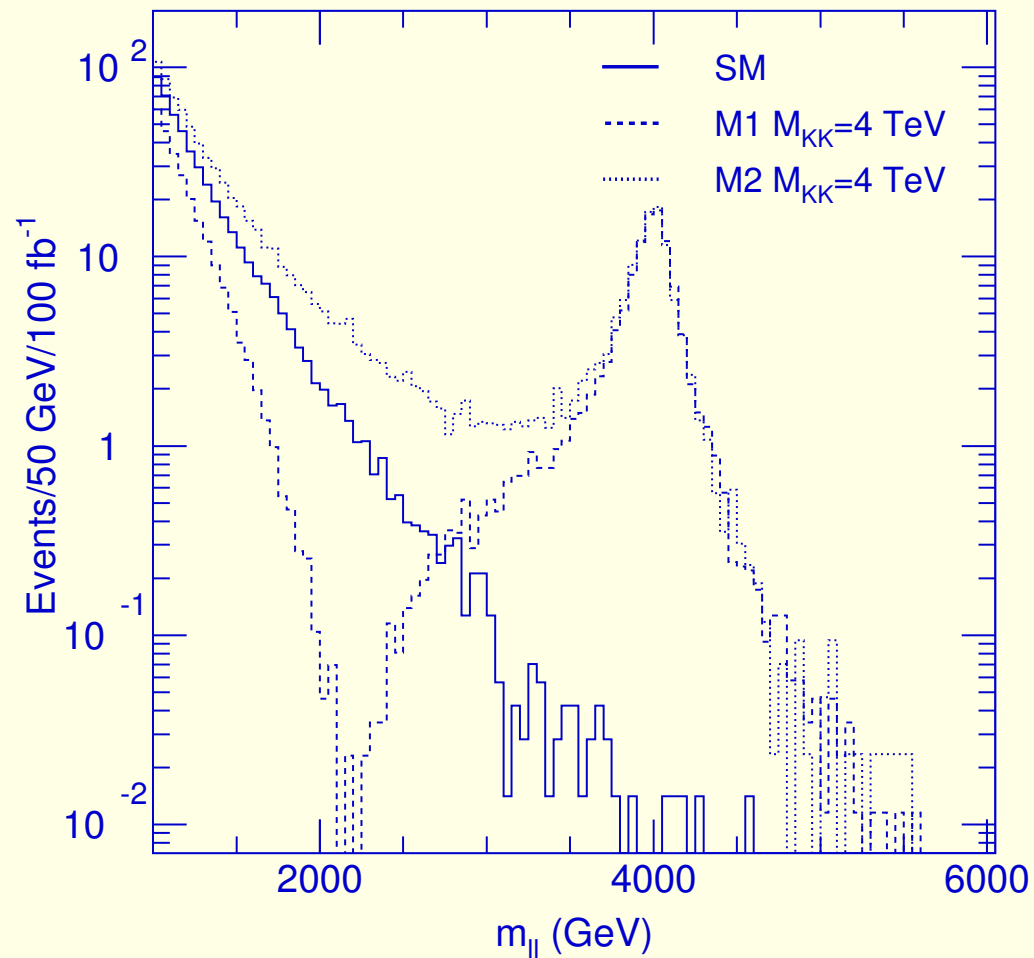


Resonance is Spin-2, confirm this by looking
 at lepton angular distribution
 Can determine spin properties for $M < 1.4 TeV$ for $100 fb^{-1}$

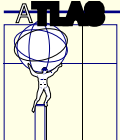


Can also have standard model fields in extra dim.
Excitations of SM particles

ATLAS e^+e^- preliminary



Insufficient reach to see second resonance



Conclusions

- We are 42 Months from first data
- Much work remains in completing and commissioning hardware and software
- Set of ongoing data challenges to test out software, Physics readiness document in 2006 – updates to Physics TDR (1999).
- Full capability of detector requires restoring staged components – vital for full luminosity operation and some physics (*b-physics*)
- Serious thinking has started about what might be done at 10^{35} and what machine and detector upgrades are needed.

