

Hadron Collider Physics

- Experimental Overview – Part I -

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Outline

- **Part I** **Introduction to Hadron Colliders**
Studies of the Strong Interaction (QCD)

- **Part II** **Electroweak Physics**
Top Physics

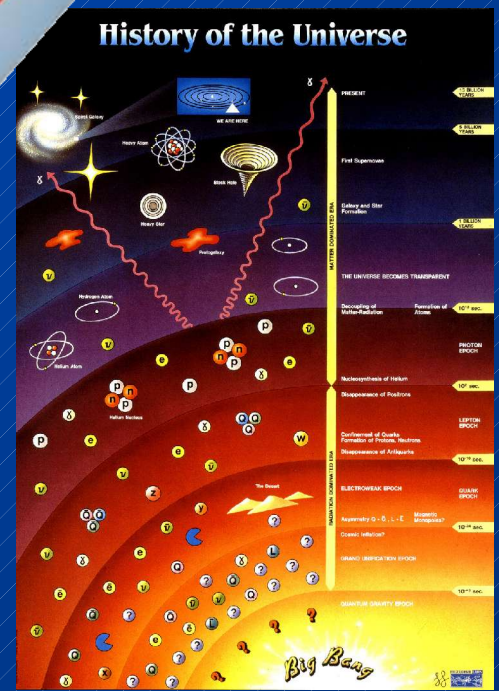
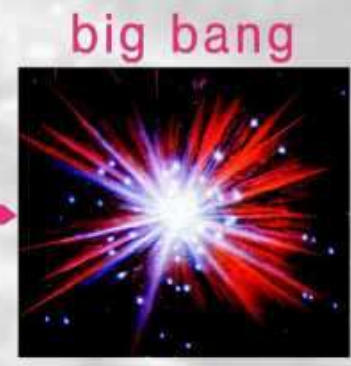
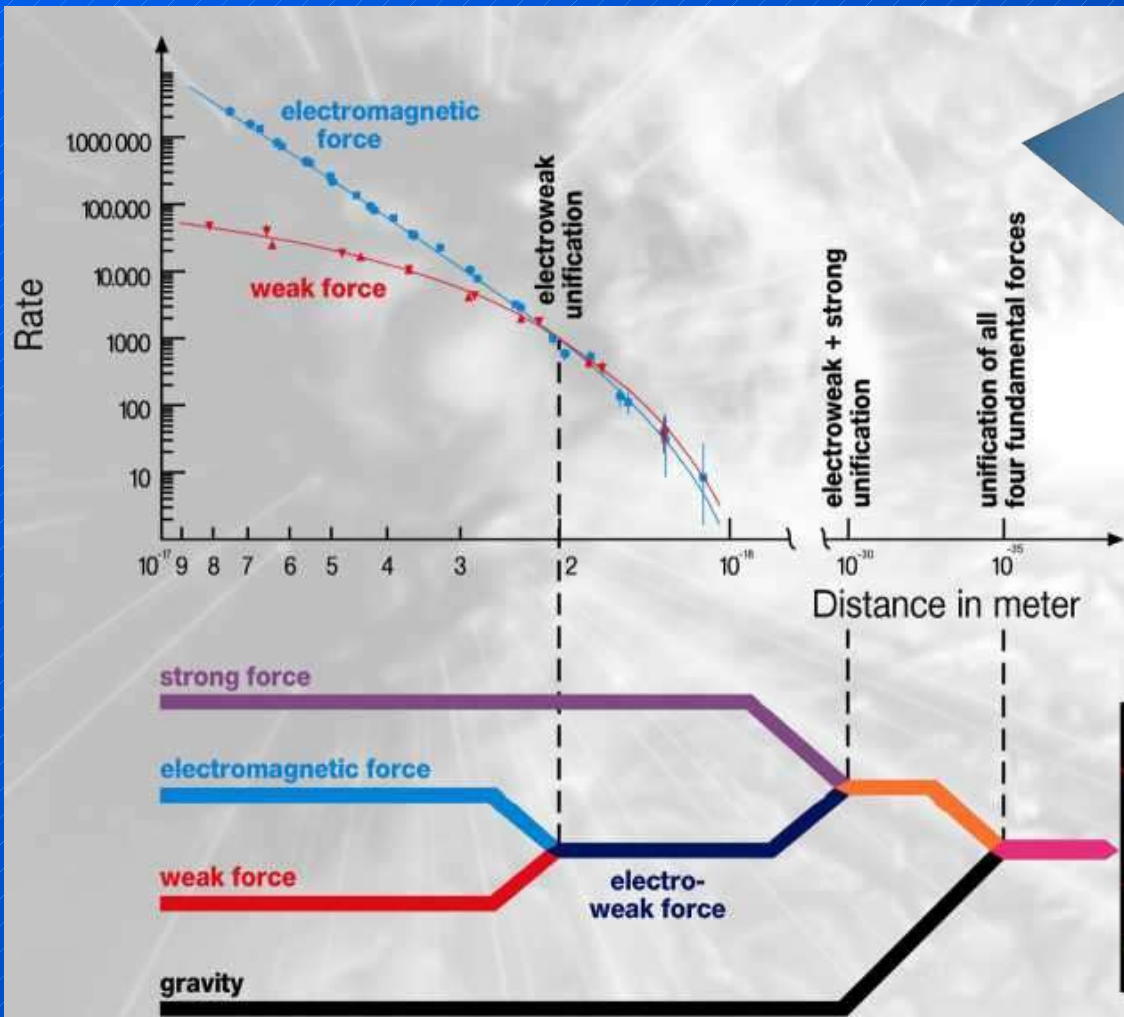
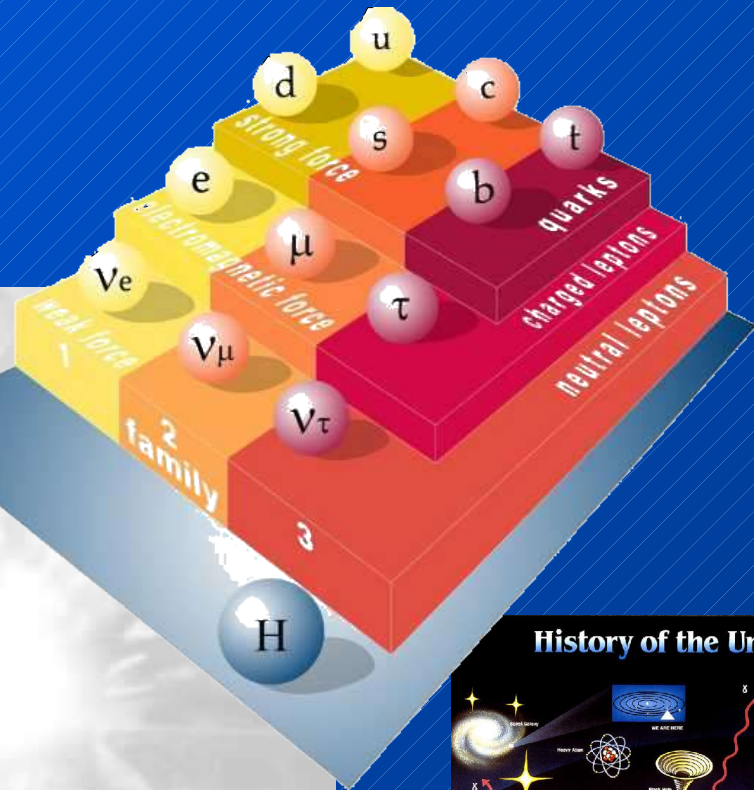
- **Part III** **Search for the Standard Model Higgs Boson**
Search for the Higgs beyond the Standard Model
Search for New Phenomena

Part I

- **Brief Summary of the Standard Model**
- **Structure of the Proton**
- **Hadron Colliders :**
PS – SPPS – Tevatron – LHC
- **The Experiments :**
CDF – DØ – ATLAS – CMS
- **Studies of the Strong Interaction (QCD)**

The Standard Model

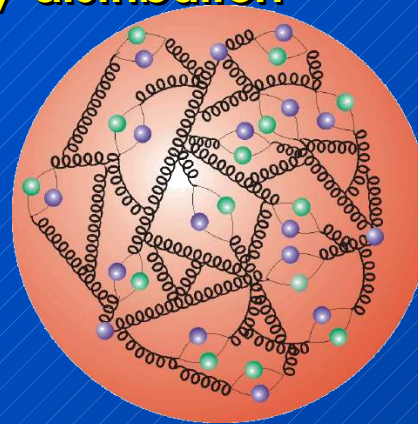
- $2 \times 3 \times 2$ fundamental particles (fermions)
- 4 fundamental forces + Higgs boson
- gauge theories with local gauge invariance



The Structure of the Proton (1)

Protons (hadrons) are not elementary particles

- they have substructure (partons = quarks and gluons)
- not well defined parton energy, but energy distribution
- various parton flavours can collide
- data analysis more complex than e^+e^- , QCD always involved



Why do we collide them ?

- much less synchrotron radiation losses
- not limited by cavity power
- but by ring size / magnet strength
- essentially no strict kinematic limit, only luminosity limited
- protons are easy to produce in vast quantities
- antiproton production ? (cooling techniques)

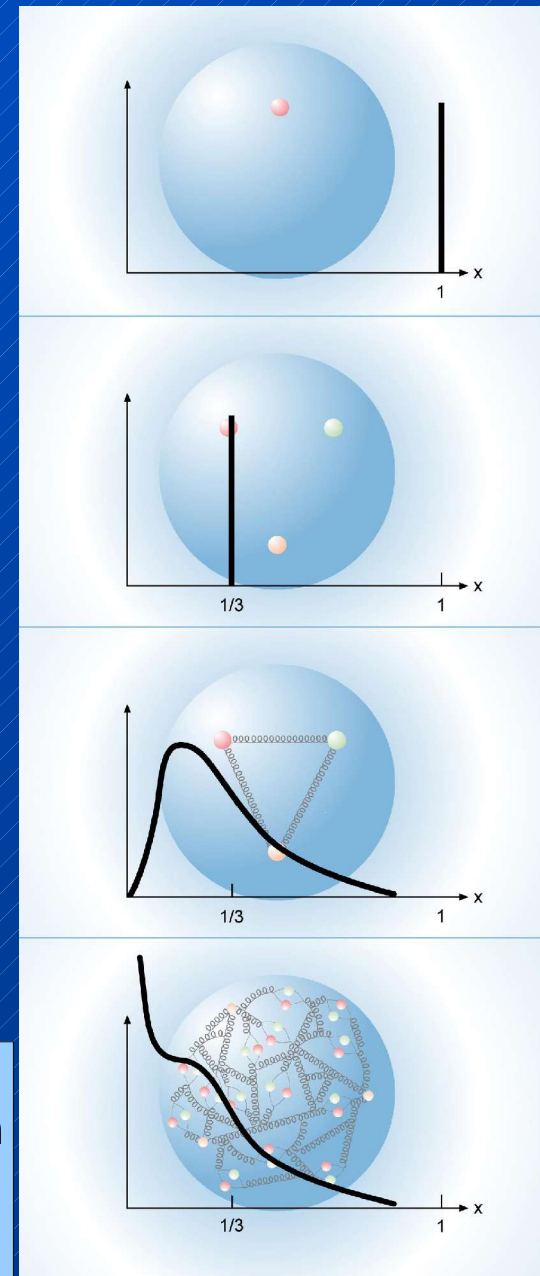
$$\Delta E = \frac{4 \pi \alpha \hbar c}{3 R} \cdot \left(\frac{E}{m} \right)^4$$

Kinematic variables:

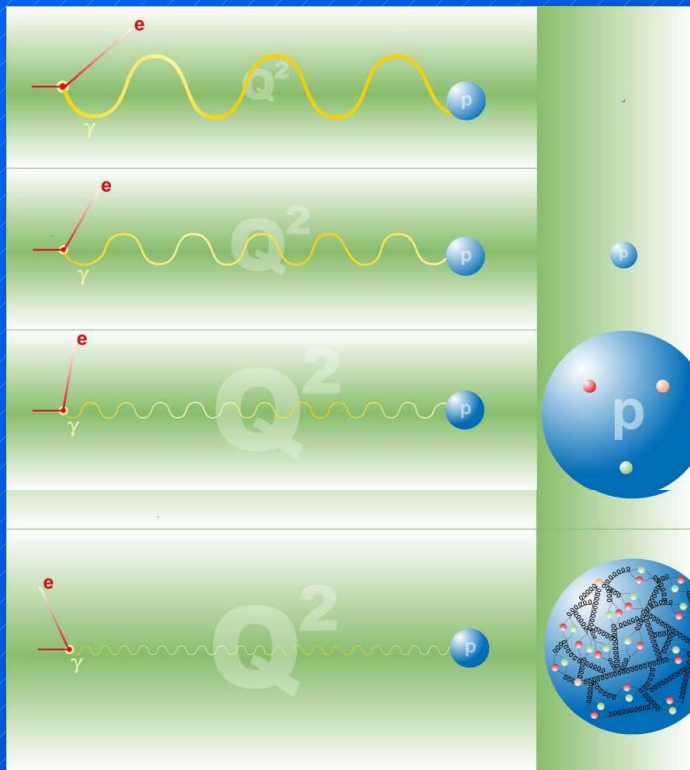
Bjorken- x = fraction of the proton momentum carried by the struck parton

Q^2 = 4-momentum transfer squared in hard interaction

$\sim 1/\text{transverse resolution power}$

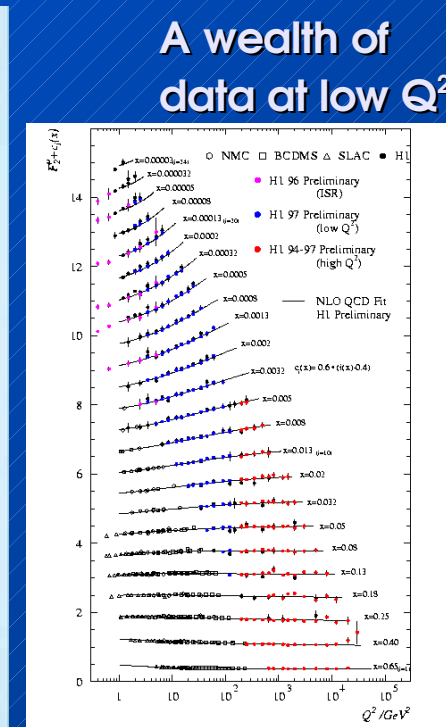
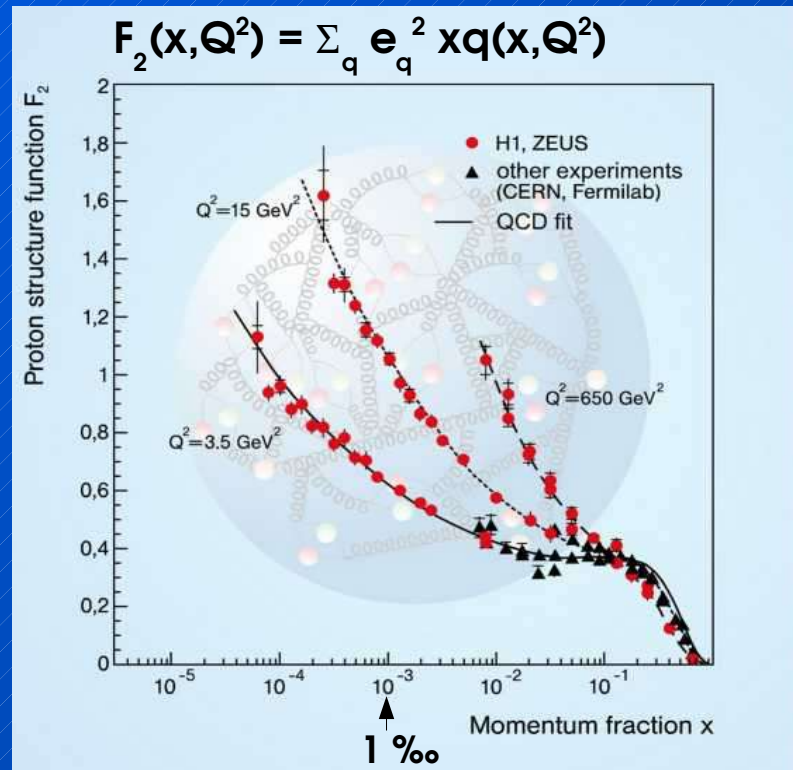


The Structure of the Proton (2)

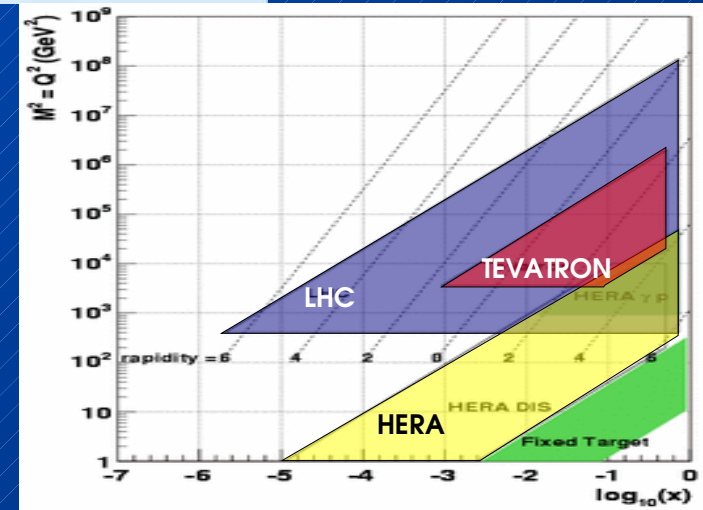


hadronic

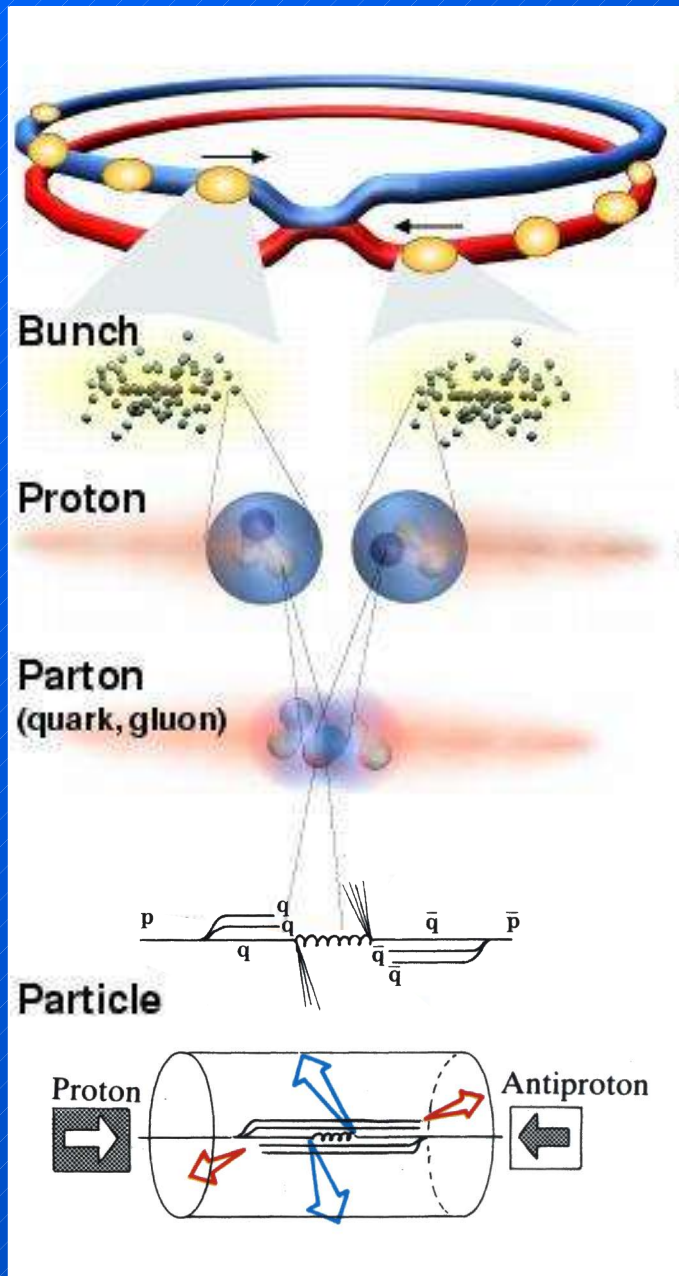
partonic



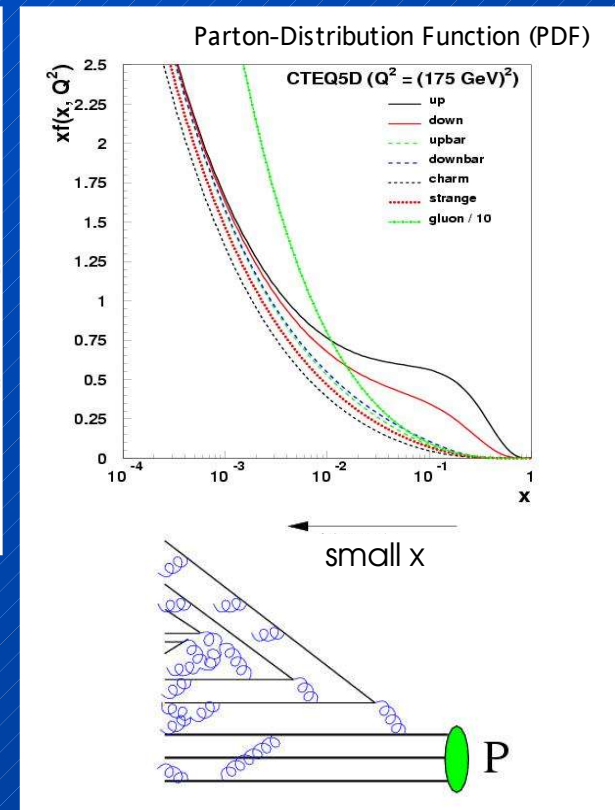
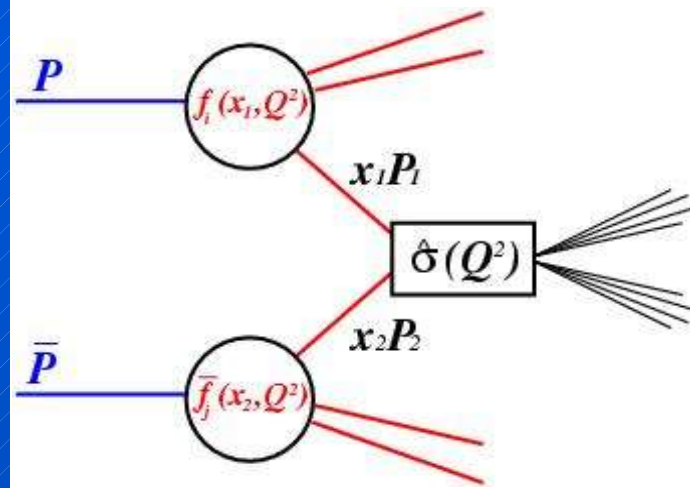
- Partonic structure at high Q^2
- Parton density (F_2) rises towards low x
- Extrapolation of HERA and fixed target measurements towards lower x and higher Q^2
- Need to measure (longitudinal) parton (momentum) distribution functions (PDFs) in every experiment parametrised by CTEQ, MRST ...



Parton-Parton Interactions



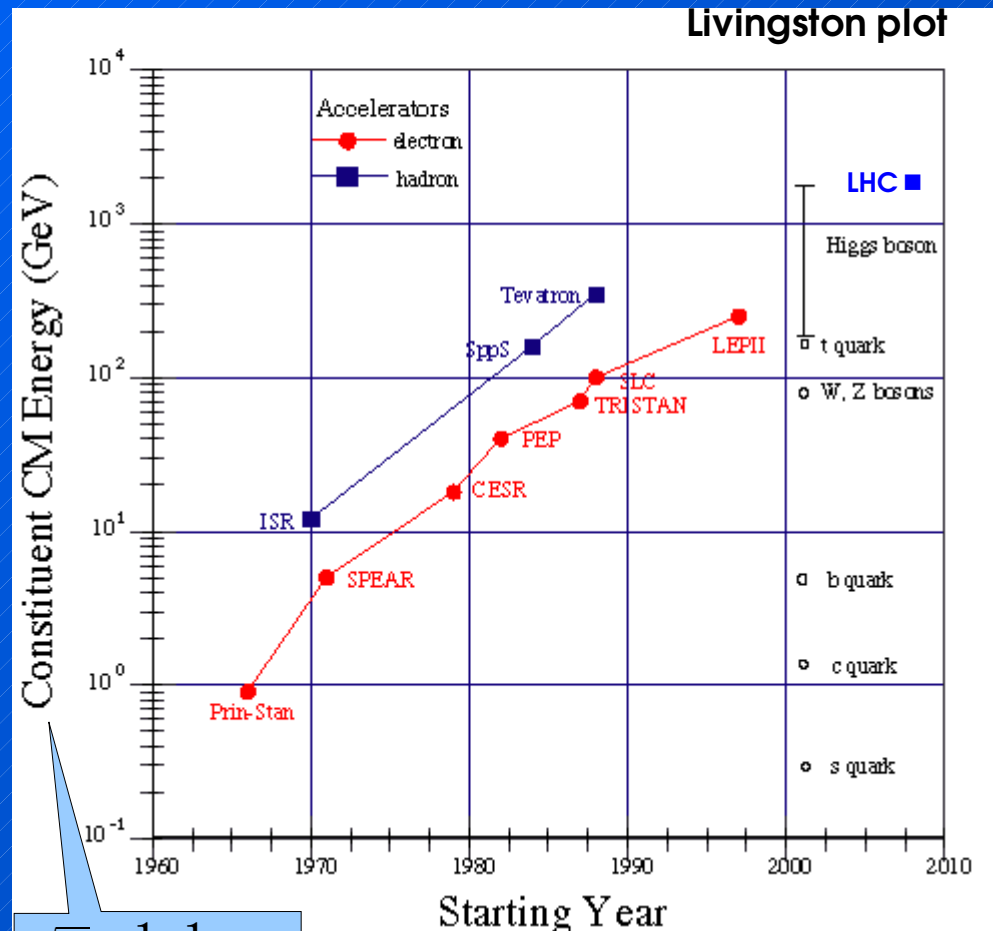
Factorization



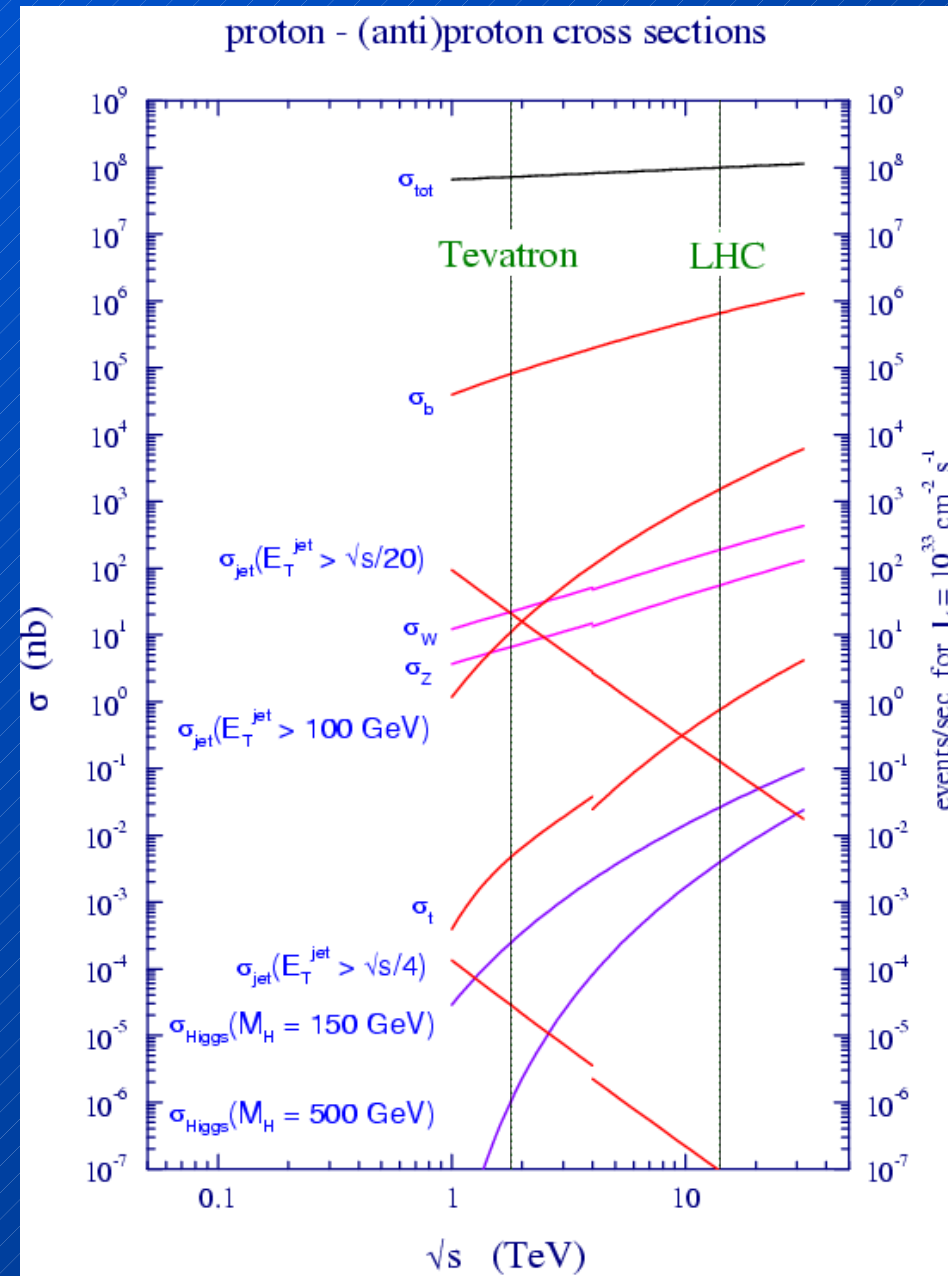
- if parton interaction requires **large CMS** (valence) **quark dominated**, e.g. top production at Tevatron
- if parton interaction possible with **low CMS** **gluon-gluon dominated**, e.g. top or light Higgs production at LHC

Parton-Parton Interactions

Progress in High Energy Physics depends on advancing the energy frontier :



$$\sqrt{s'} \approx \frac{1}{2} \cdot \frac{1}{3} \sqrt{s}$$



Kinematics

- **parton-parton center of mass energy**, estimate: $\sqrt{s'} \approx 1/6 * \sqrt{s}$
- partons have longitudinal momentum distribution !
⇒ **boost of CMS system** along beam line, a priori unknown

- focus on **transverse quantities**:

- transverse momentum P_T
- transverse Energy E_T
- missing transverse energy MET :

$$MET = \left| -\sum_{vis} \vec{P}_T \right|$$

- **angles** measured in :

- **rapidity**: $y = 1/2 \ln \{ (E+P_L) / (E-P_L) \}$

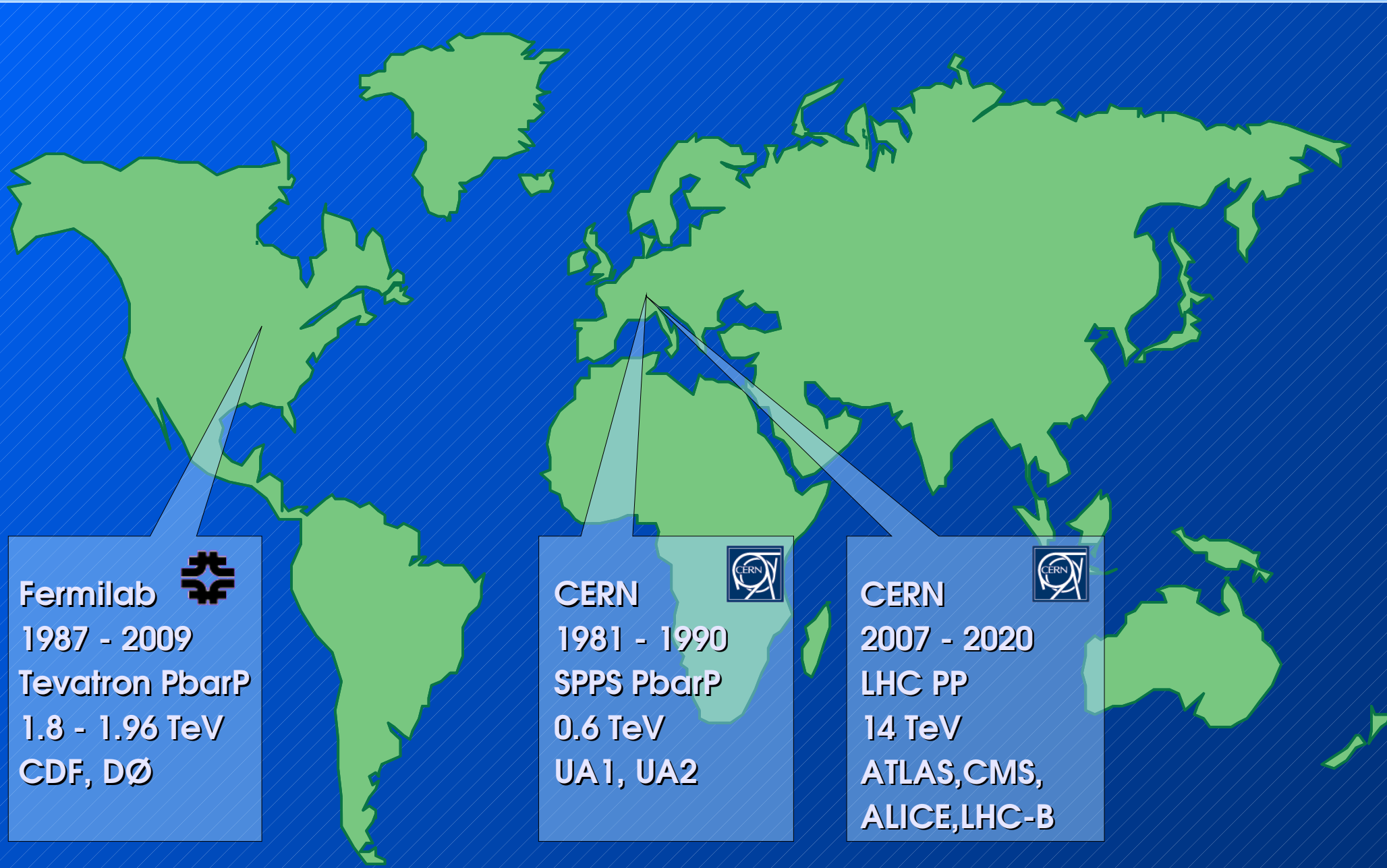
if $m \ll E, P_L$

- **pseudorapidity**: $\eta = -\ln \tan \theta/2$

rapidity intervals y are Lorentz-invariant

distance measure : $R^2 = (\Delta\phi)^2 + (\Delta\eta)^2$

Hadron Colliders



Fermilab
1987 - 2009
Tevatron PbarP
1.8 - 1.96 TeV
CDF, DØ

CERN
1981 - 1990
SPPS PbarP
0.6 TeV
UA1, UA2

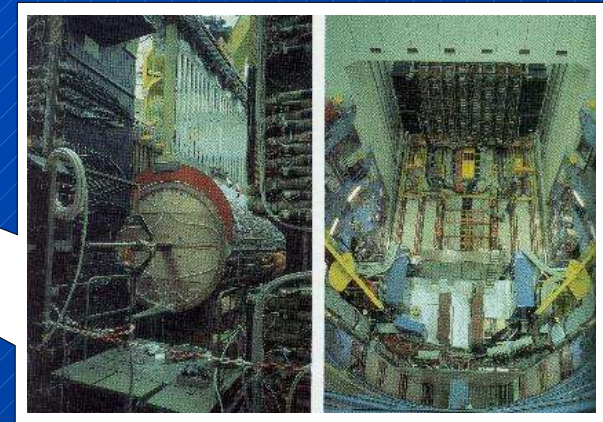
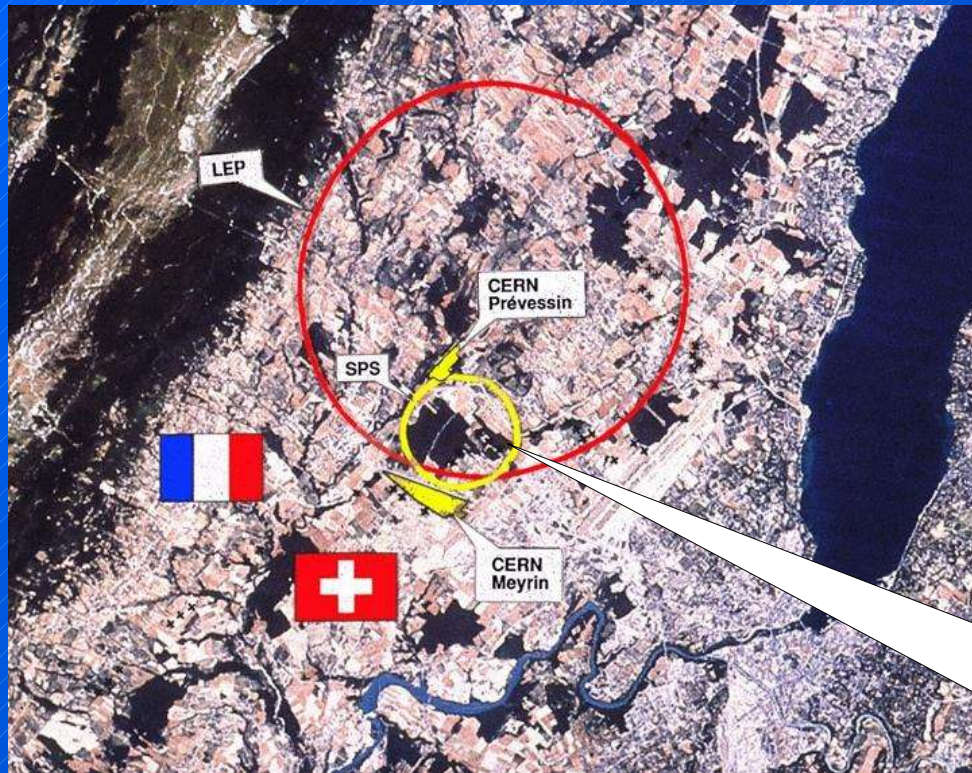
CERN
2007 - 2020
LHC PP
14 TeV
ATLAS, CMS,
ALICE, LHC-B

CERN's Proton Synchrotron



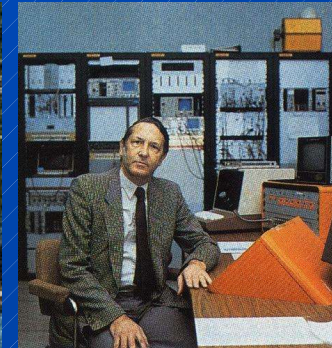
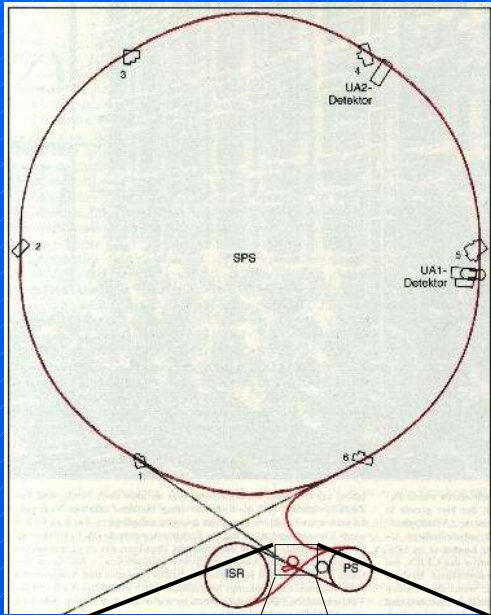
**proton-synchrotron (PS) accelerated first protons to 24 GeV in 1959
... still running today and in the future for the LHC ...**

CERN's Super-Proton Synchrotron

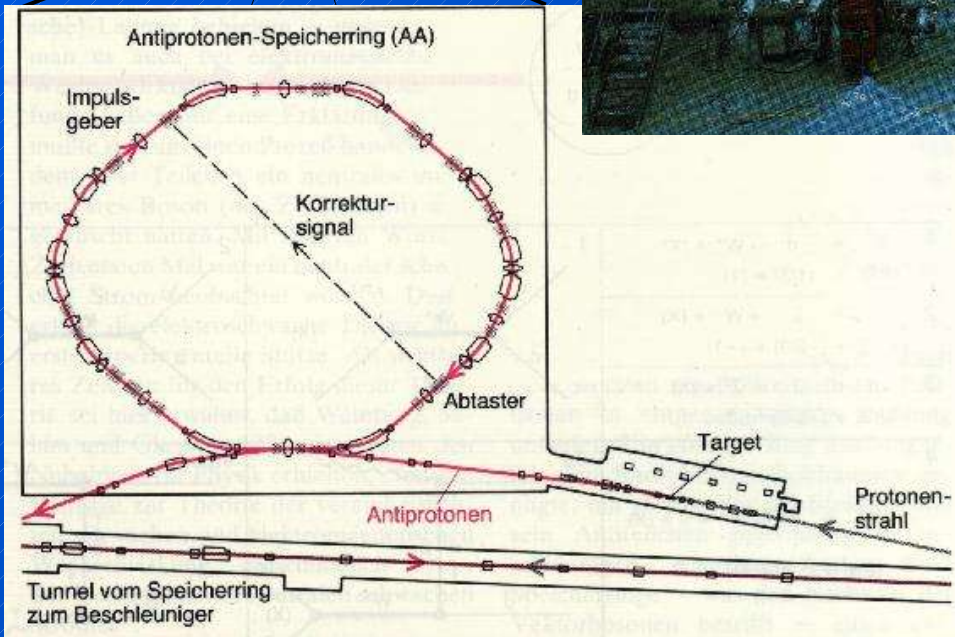


- first p-pbar reactions observed in Juli 1981 ... at 270 GeV per beam
- then turned into a proton-antiproton collider (one beampipe)
- collider experiments UA1 and UA2 (Nobel prize for machine and W/Z-discovery)

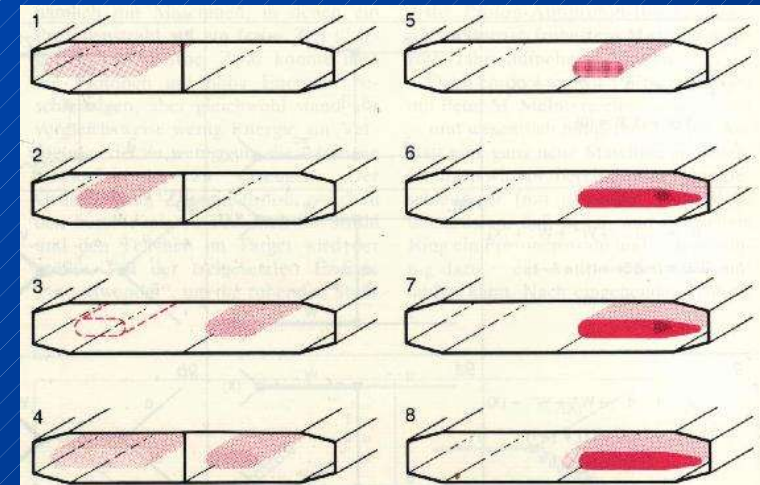
Stochastic Cooling



Simon van der Meer

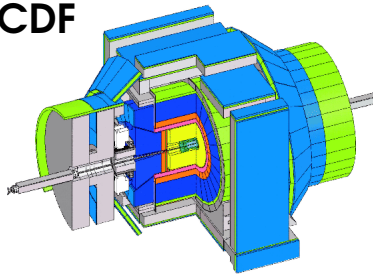


$5 \cdot 10^{11}$ Pbar/beam
1 Pbar per $3 \cdot 10^5$ P



The TEVATRON at Fermilab (1)

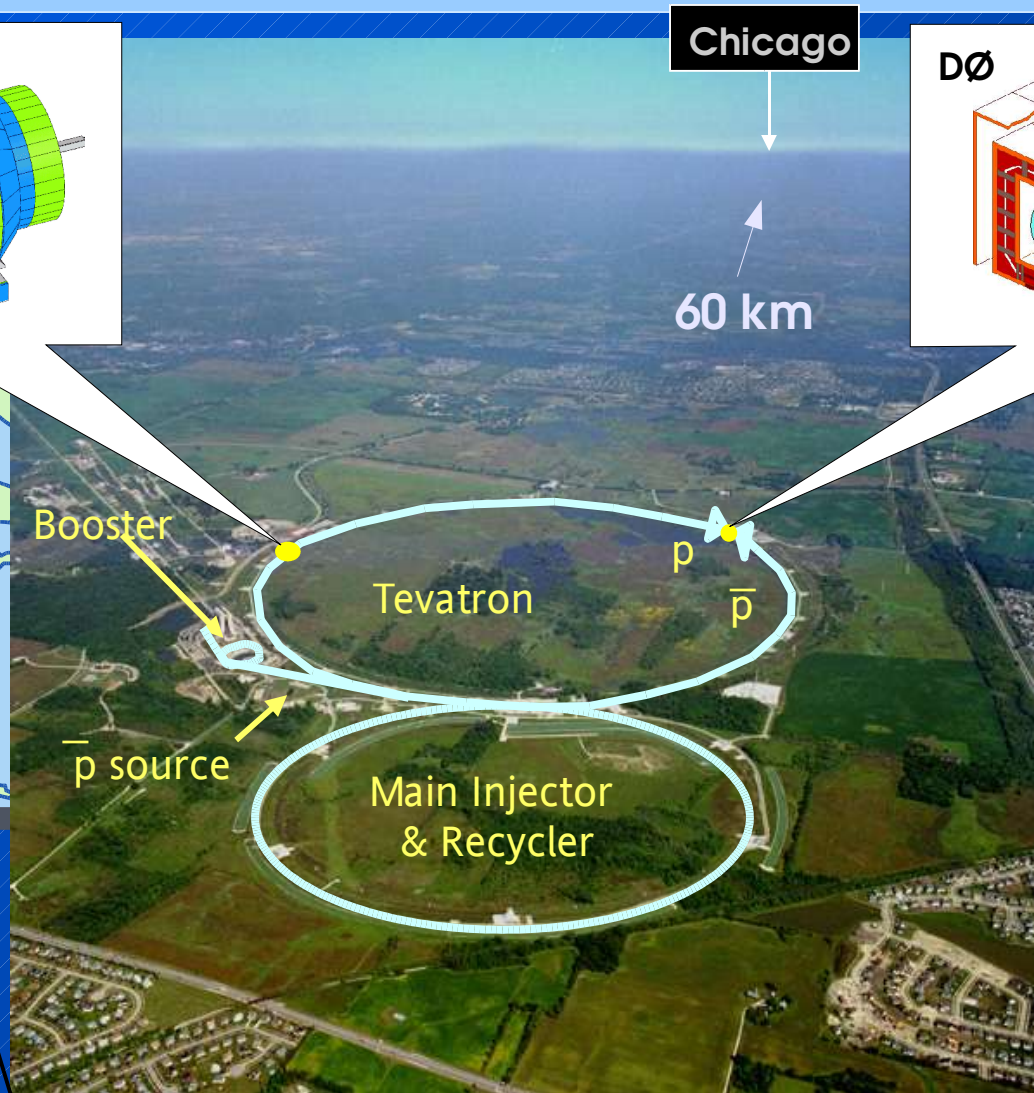
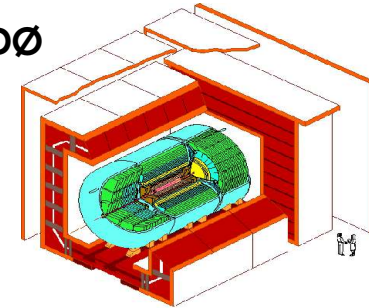
CDF



Chicago

60 km

DØ



$$\sqrt{s} = 1.96 \text{ TeV}$$

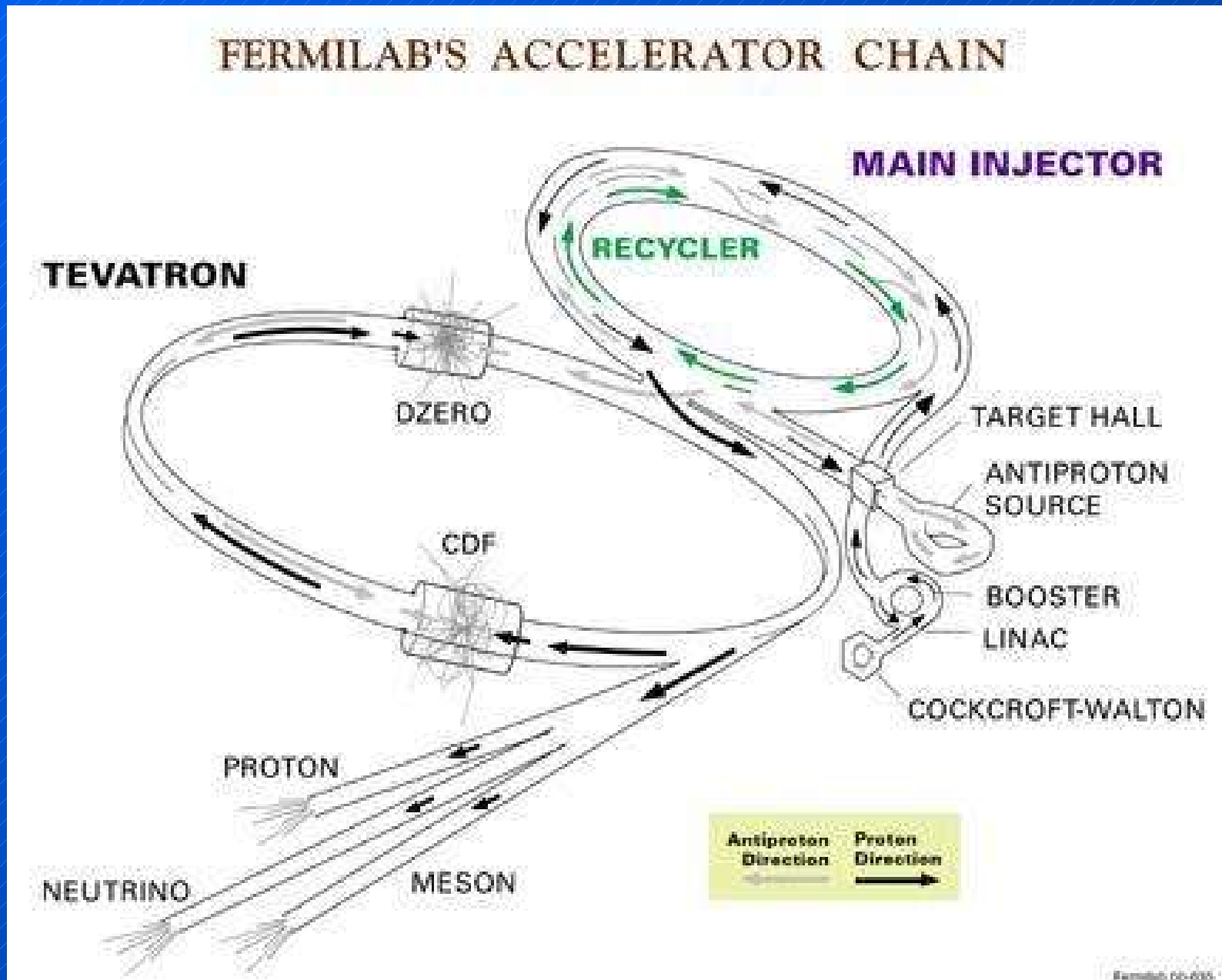
$$\Delta t = 396 \text{ ns}$$

Run I 1987 (92)-95 $L_{\text{int}} \sim 125 \text{ pb}^{-1}$

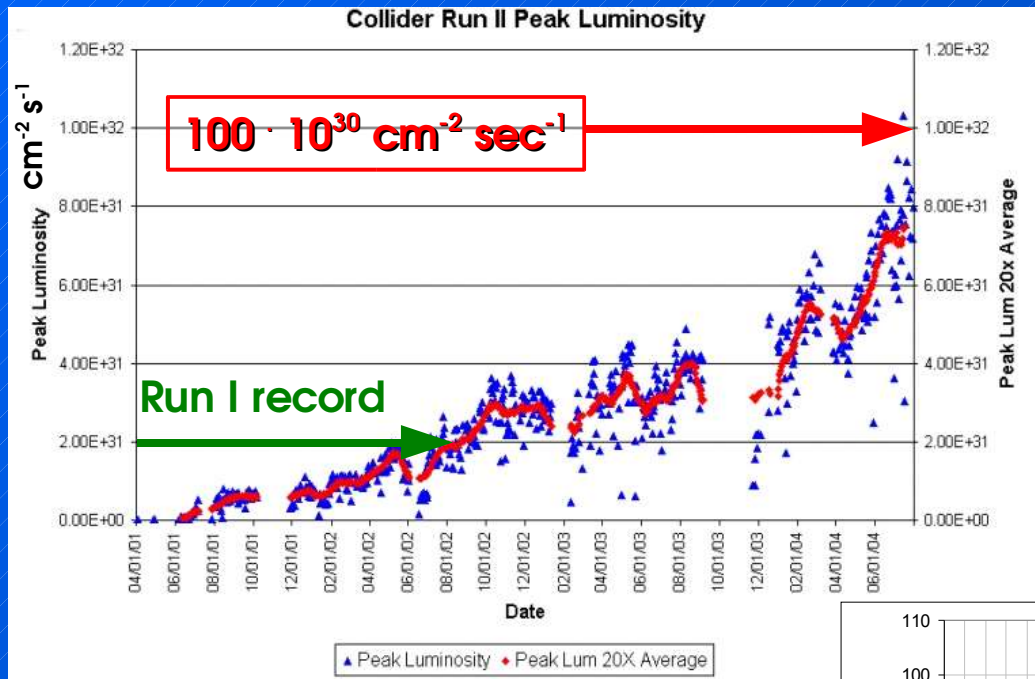
Run II 2001-09(?)

> 40-times larger dataset
at increased energy

The TEVATRON at Fermilab (2)

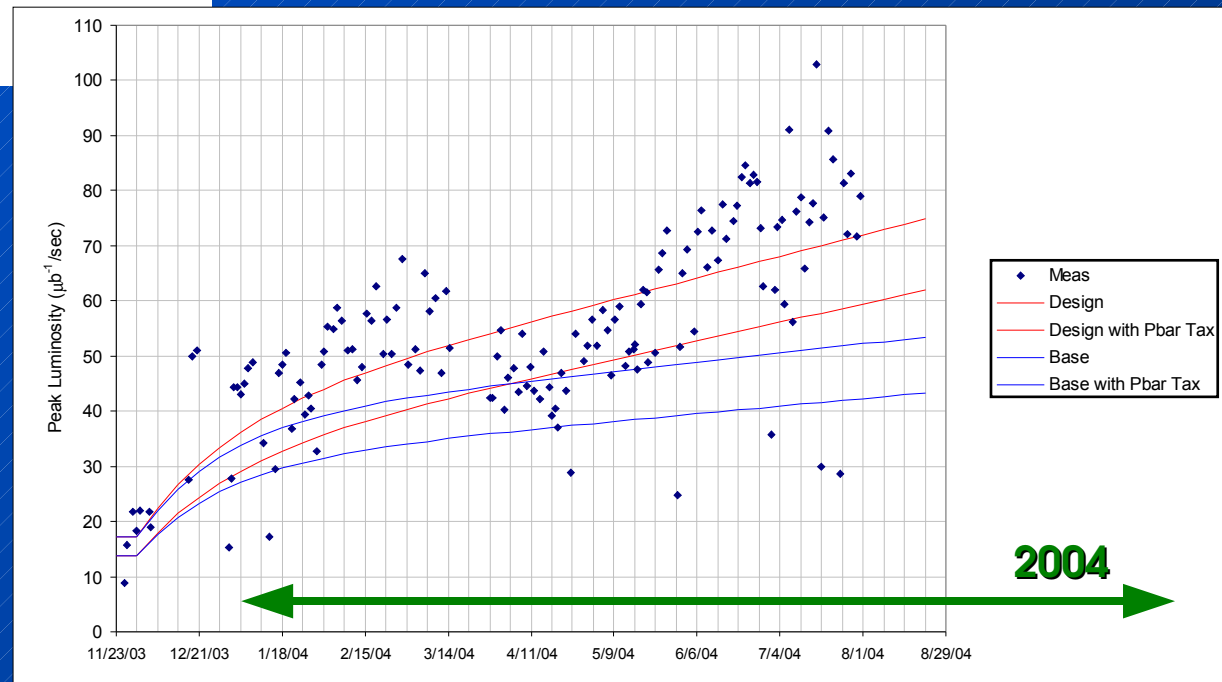


The TEVATRON Performance (1)

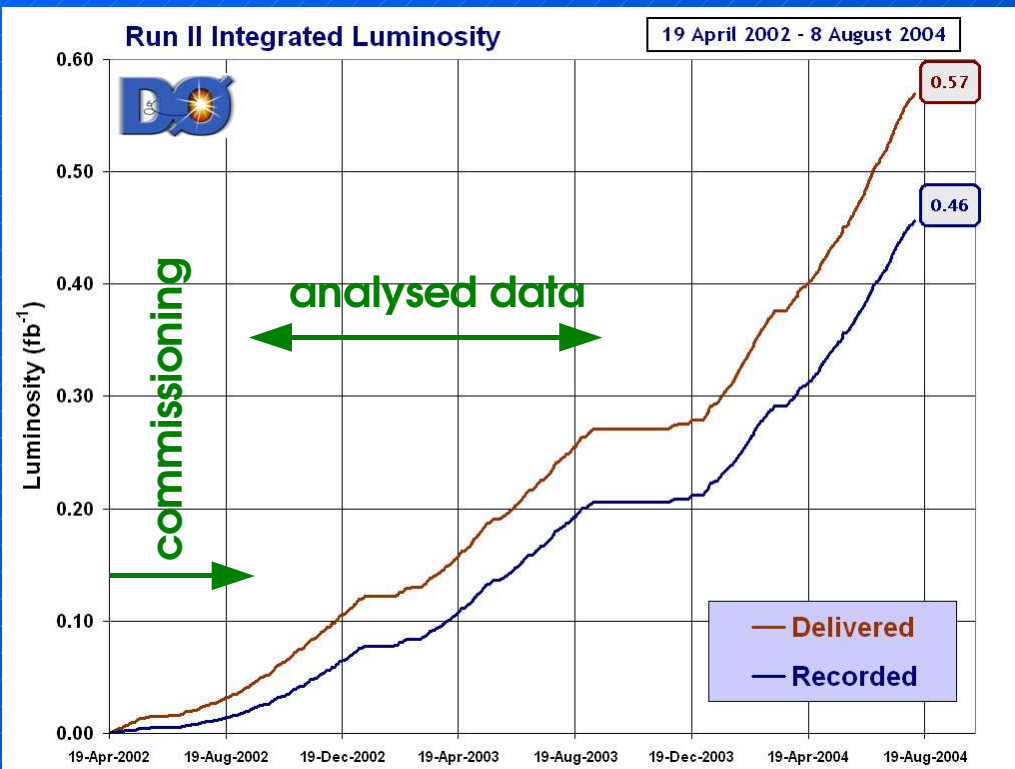


Peak luminosity:

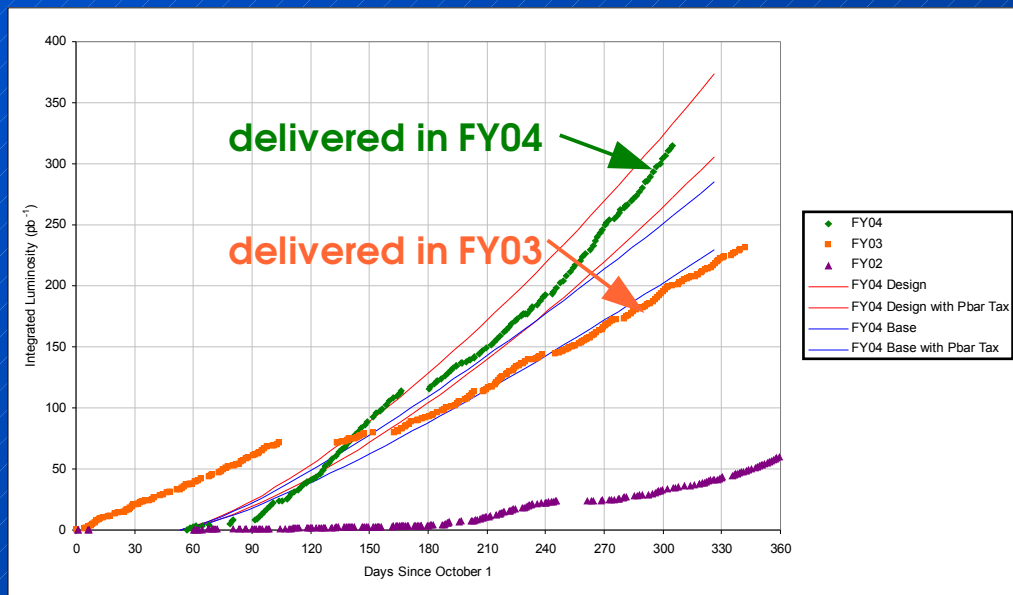
- 1 · 10¹² \bar{P} per beam, 'only' twice SPPS !
- in spring 2002 passed Run I record
- in 2004 often above optimistic design scenario



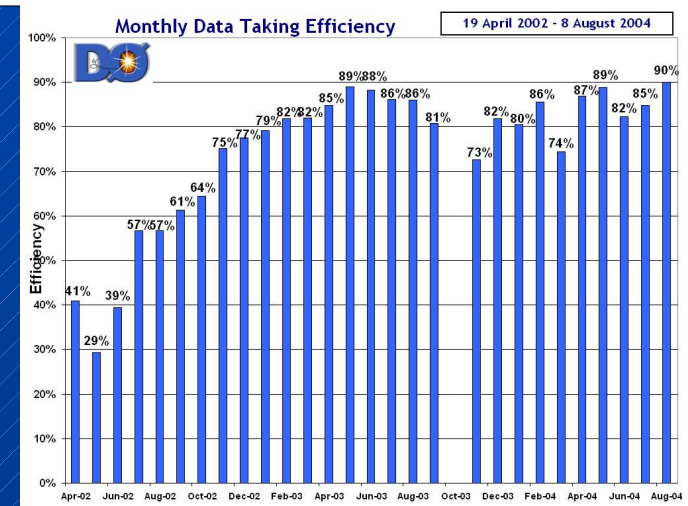
The TEVATRON Performance (2)



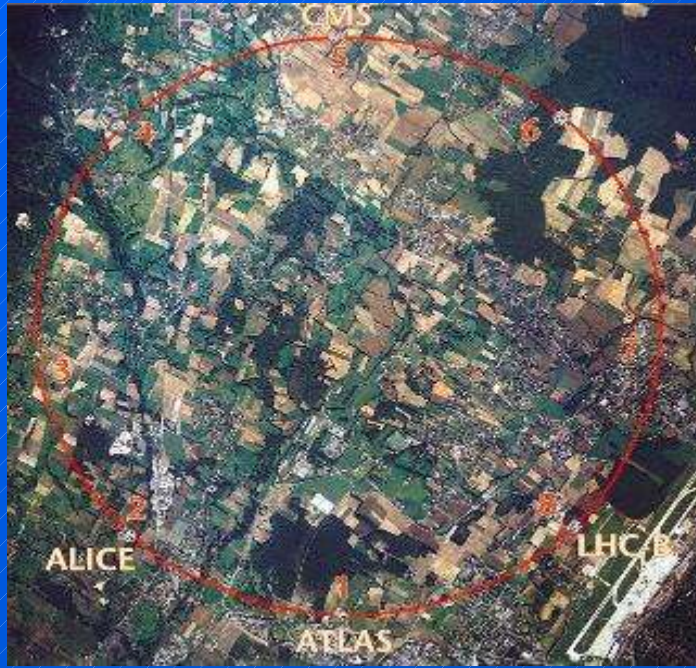
TEVATRON at upper limit of optimistic scenario



- CDF and DØ $\sim 500 \text{ pb}^{-1}$ recorded each
- typical data taking efficiency $\sim 90\%$
- commissioning is over



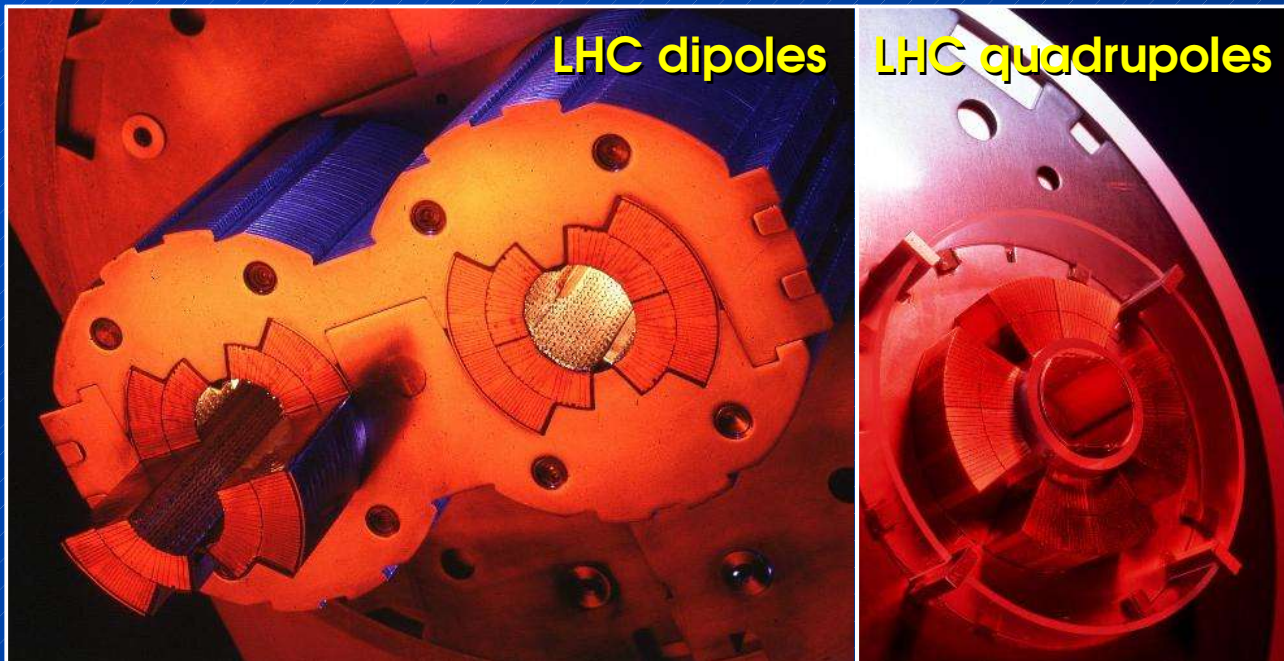
The Large Hadron Collider



The Large Hadron Collider:

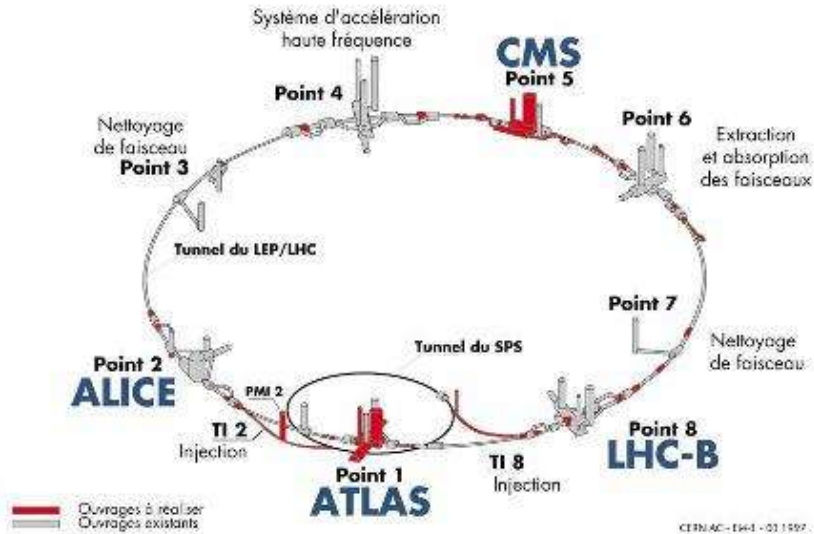
- proton-proton collisions
- no antiprotons, $3 \cdot 10^{14}$ P
- at 14 TeV center of mass energy
- 40 Mio. collisions per second
- first collisions in 2007
- 4 experiments:
ATLAS, CMS, ALICE, LHC-B

- two separate beampipes
- two accelerators



The Large Hadron Collider

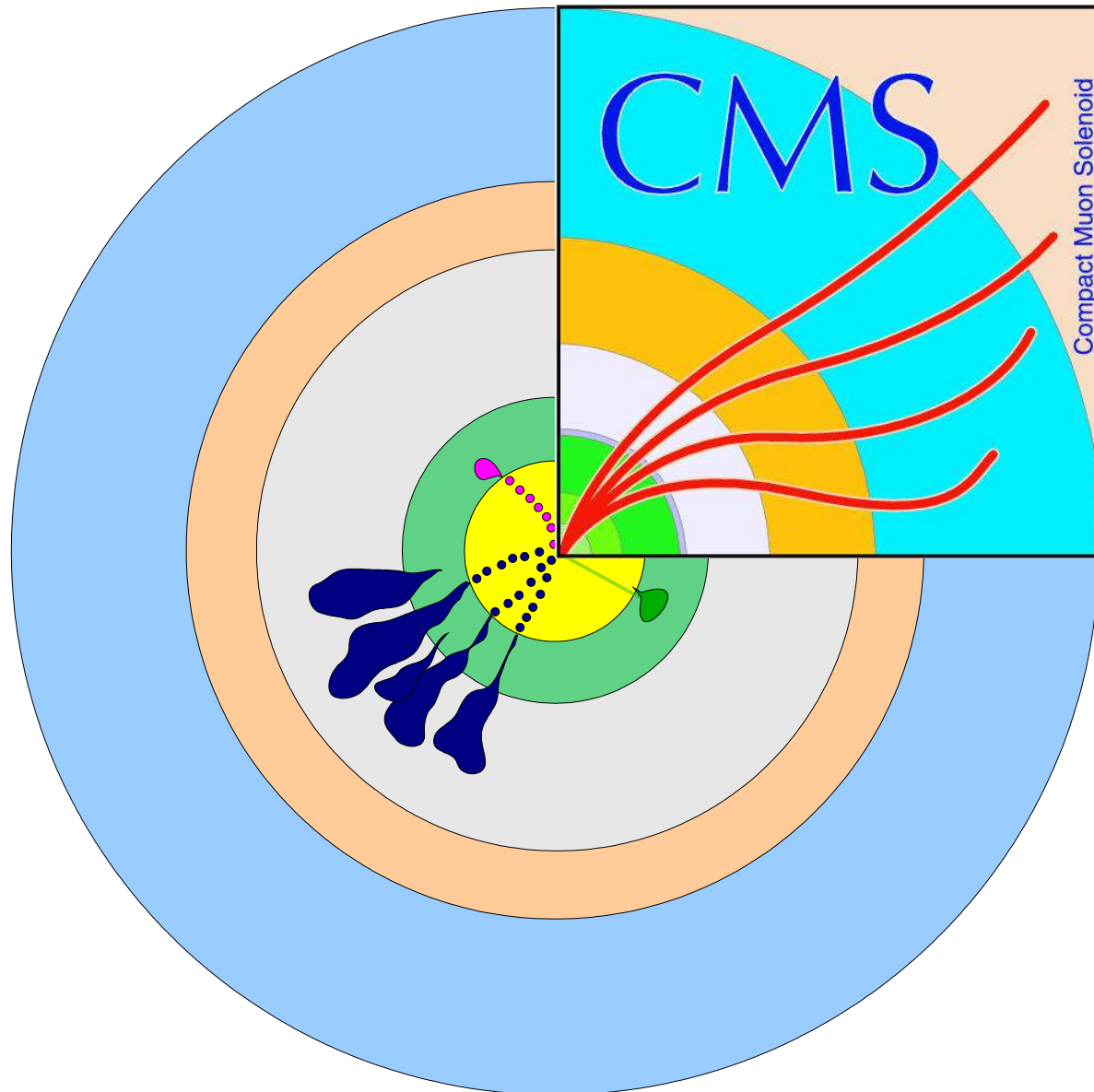
Vue d'ensemble des ouvrages souterrains du LHC



Collider Comparison

	<i>SPPS</i>	<i>Tevatron</i>	<i>LHC</i>
Physics start	1981	1987	2007
Particles	Pbar-p	Pbar-p	p-p
cm energy (TeV)	0.62	1.96	14
Lumi ($10^{30} \text{ cm}^{-2} \text{ s}^{-1}$)	6	50-100	$0.1 - 1.0 \times 10^4$
Lumi ($\text{fb}^{-1} \text{ year}^{-1}$)	0.05	0.5	100
Bunch spacing (ns)	3800	396	25
Particles per bunch (10^{10})	P: 15, Pbar: 8	P: 24, Pbar: 3	P: 11.5
Max.no Pbar in accumulator	1.2×10^{12}	2.6×10^{12}	-
Bunches	6 + 6	36 + 36	2835 + 2835
Circumference (km)	6.9	6.28	26.7
Nr. dipoles	232	774	1232 (main dipoles)
Magnet type	warm	cold, warm iron	cold, cold iron
Peak magnetic field (T)	1.4	4.4	8.3

Particle Identification








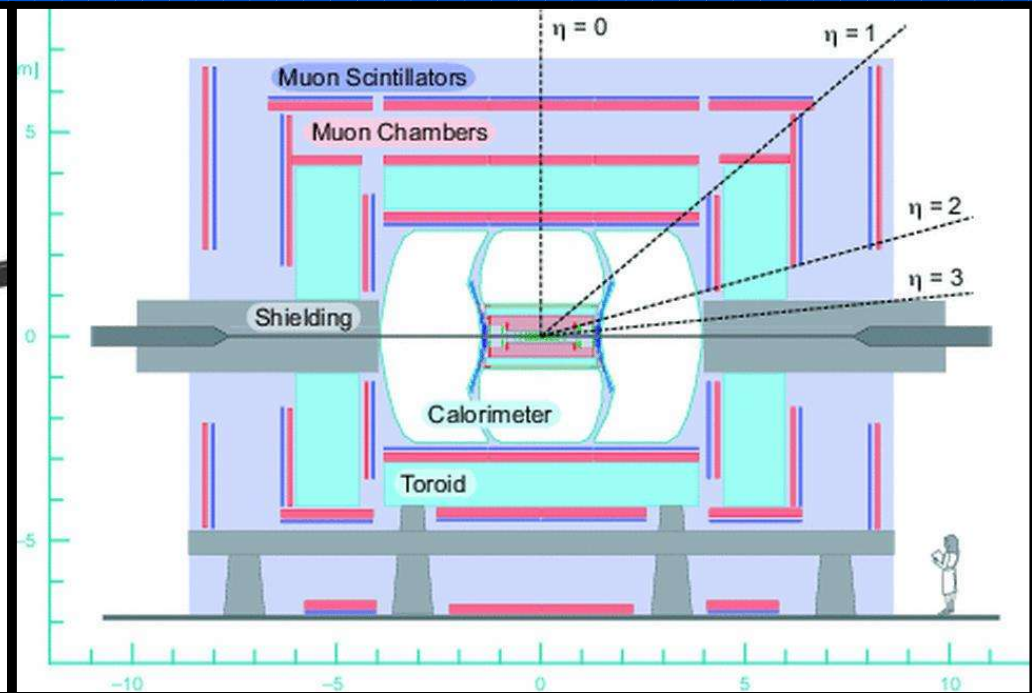
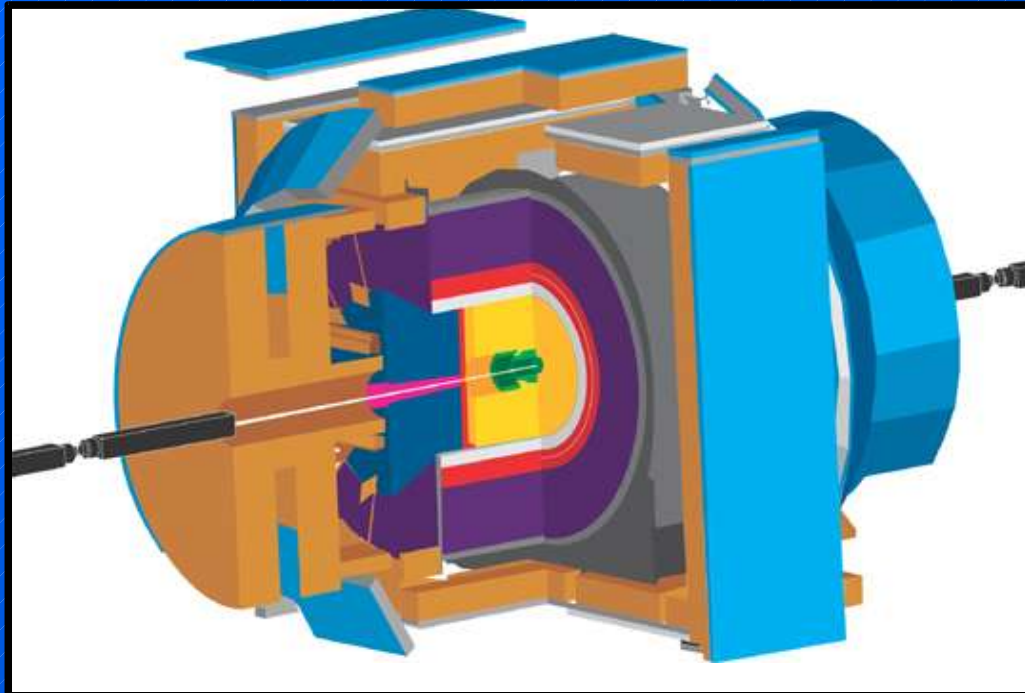
Muon (high energy)
(medium ene.)
(low energy)

Photon

Electron

Quark → Jet

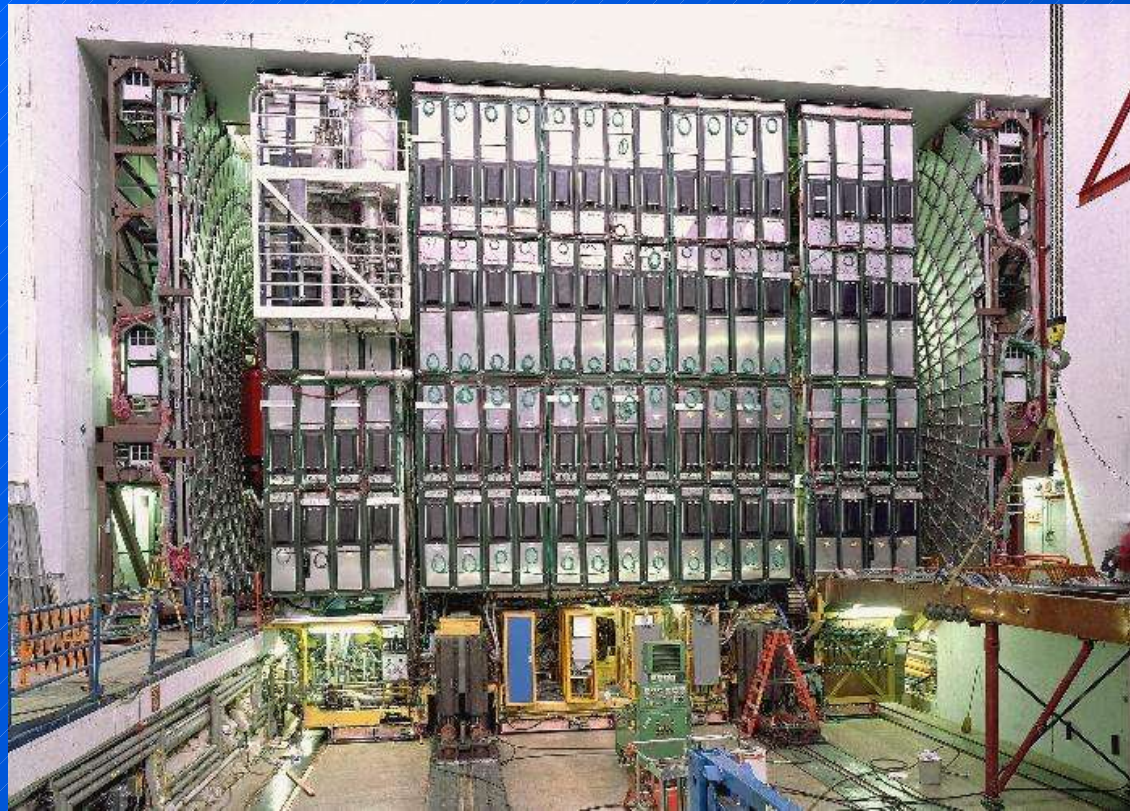
	Tracking detector
	Electron calorimeter
	Hadron calorimeter
	Magnet coil
	Muon chambers



- new bigger silicon, new drift chamber, TOF
- Upgraded calorimeter and muon system
- Upgraded DAQ/trigger
- Displaced track trigger
- ~750 physicists

- new silicon and fibre tracker
- new ~2 T solenoid
- upgraded muon system
- upgraded (track) trigger/DAQ
- 19 countries, 83 institutes, 664 physicists

The DØ Experiment



compensating
Uranium/Liquid Argon calorimeter



Calorimeter:

Transverse Segmentation

$$\Delta\eta \times \Delta\phi = 0.1 \times 0.1$$

(EM3)

$$\Delta\eta \times \Delta\phi = 0.1 \times 0.1$$

L1/L2 fast trigger readout in

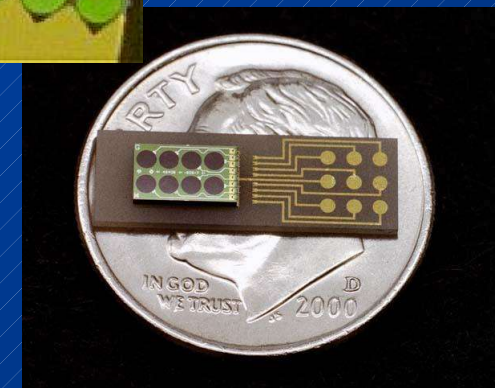
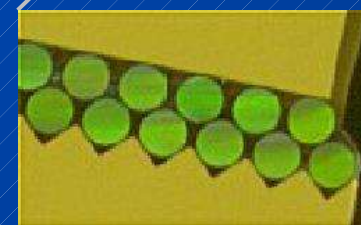
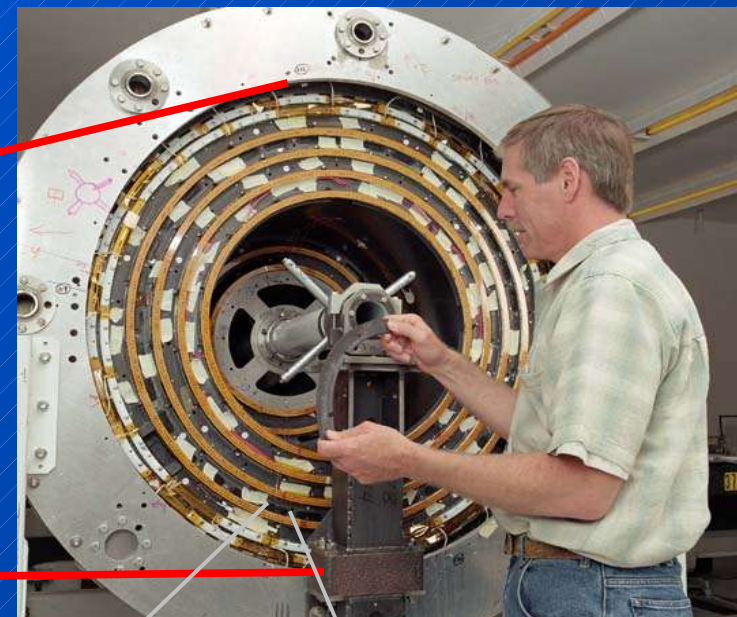
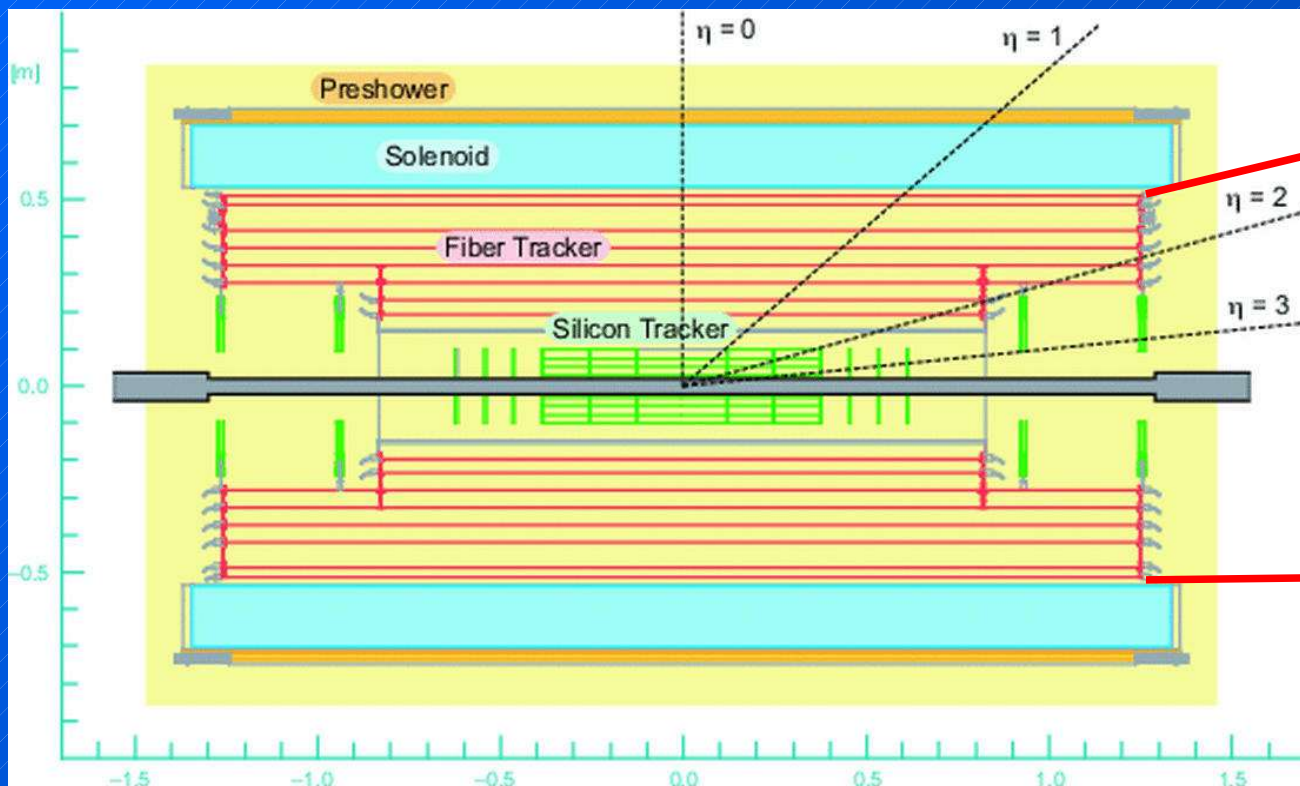
0.2 x 0.2 towers

resolutions:

EM: $\sigma_E/E = 15\% / \sqrt{E}$

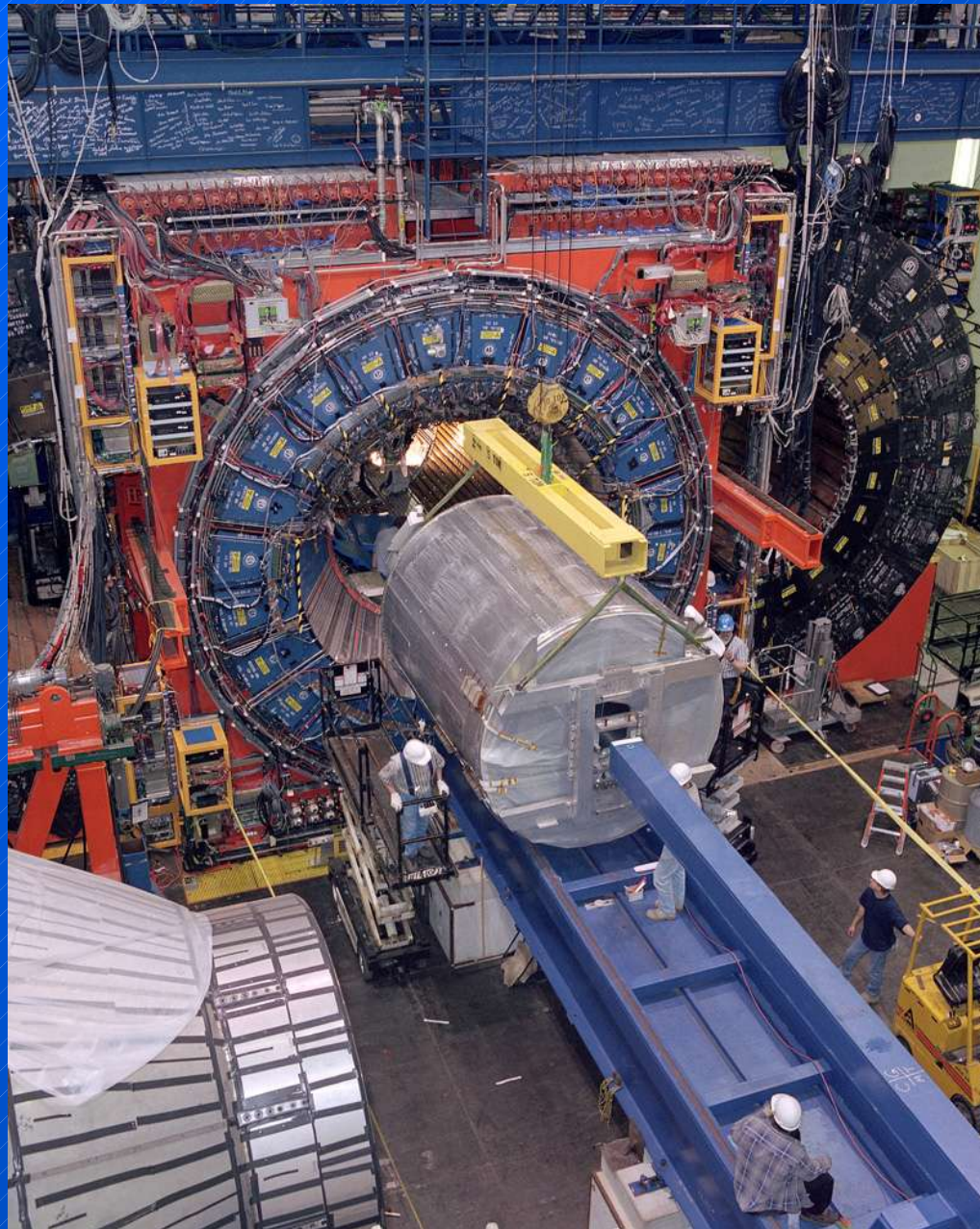
HAD: $\sigma_E/E = 50\% / \sqrt{E}$

The DØ Fiber Tracker



- **Central Fiber Tracker:**
 - 77k fibers in eight barrels, 800 μm diameter fibers
 - 3° stereo layers in each barrel
 - VLPC readout, ~ 7 photo-electrons/track at $\eta = 0$

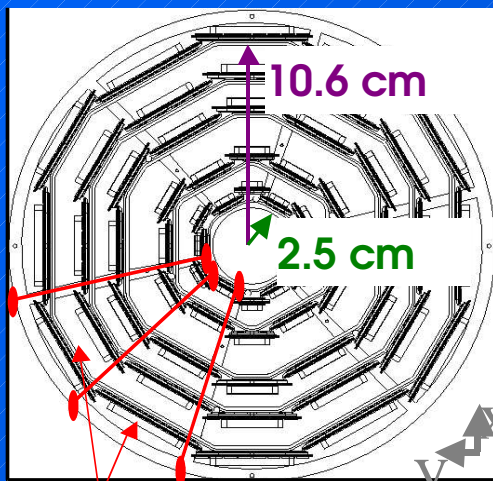
The CDF Experiment



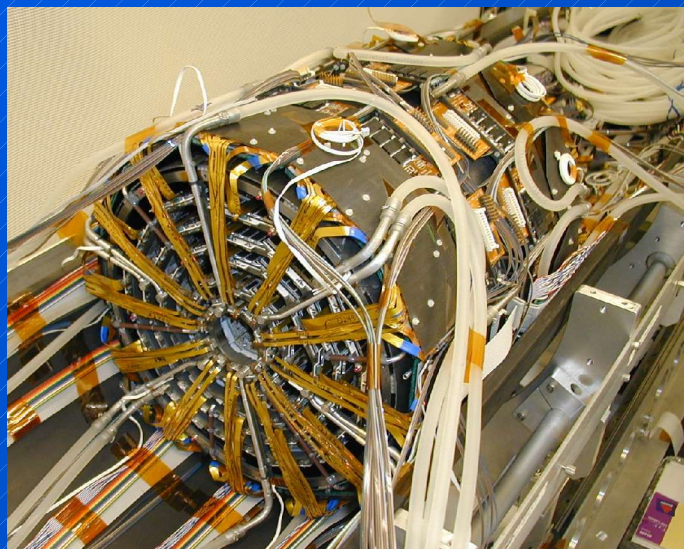
drift chamber (COT)
radius 40 - 137 cm, $|z| < 160$ cm
60 axial, 24 stereo hits (3°)

scintillator based calorimeter
EM: $\sigma_E/E = 13.5\% / \sqrt{E}$
HAD: $\sigma_E/E = 50\% / \sqrt{E}$

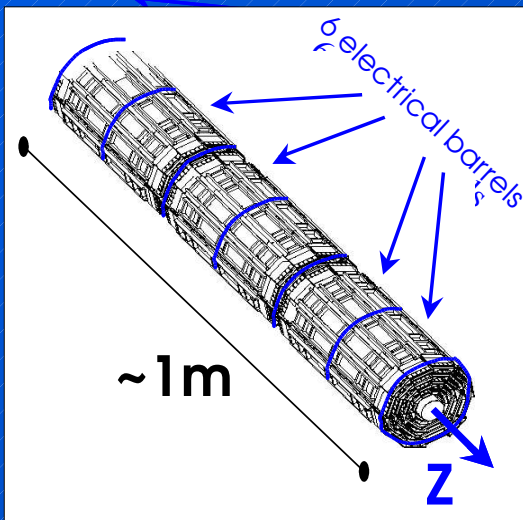
Baseline Upgrade: SVX II



Note wedge symmetry



- 7-8 silicon layers
- 722, 432 channels
- r_ϕ, r_z views
- $z^{\max} = 45 \text{ cm}, \eta^{\max} = 3$
- $1.3 < r < 30 \text{ cm}$



- 5 double-sided layers
 - 5 axial, 3 x 90°, 2 x 1.2°
- Tight alignment tolerances
 - for displaced track trigger
- Highly symmetric
 - 12-fold in ϕ
 - 6-fold in z

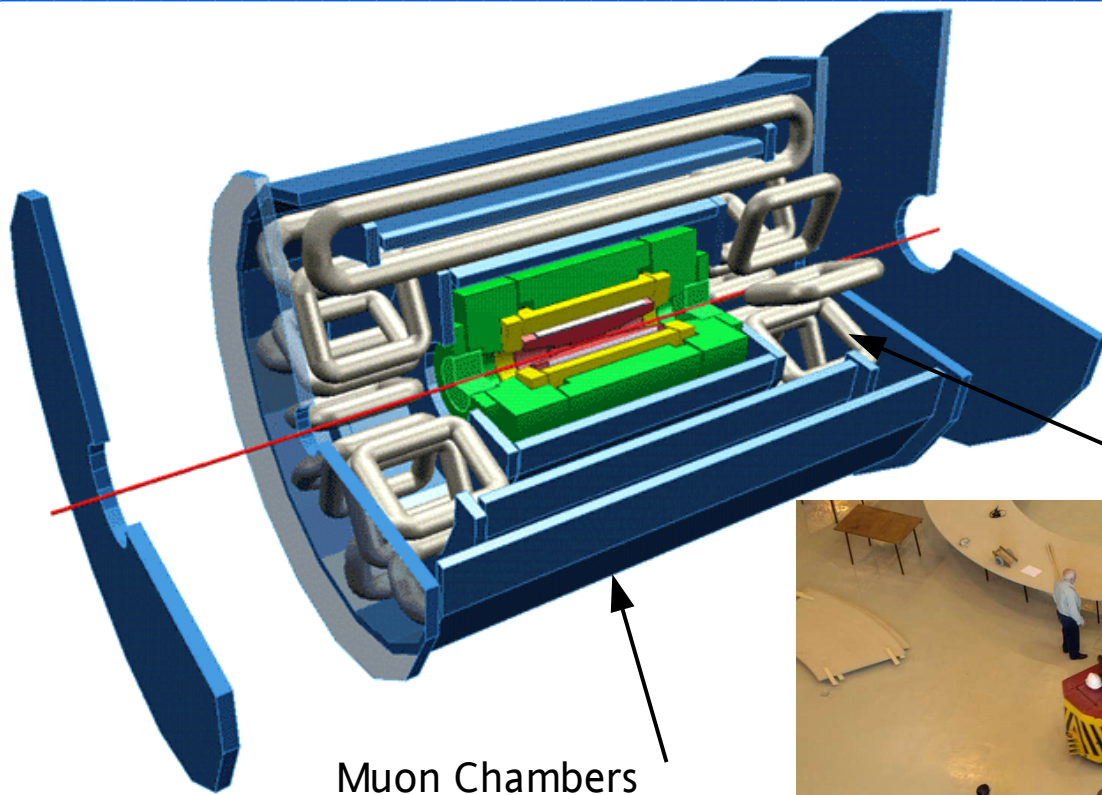
Run II physics goals:

- properties of top quark
- precision Electroweak
- CKM, B_s mixing
- search for new phenomena
- tests of QCD

heavy flavour tagging:

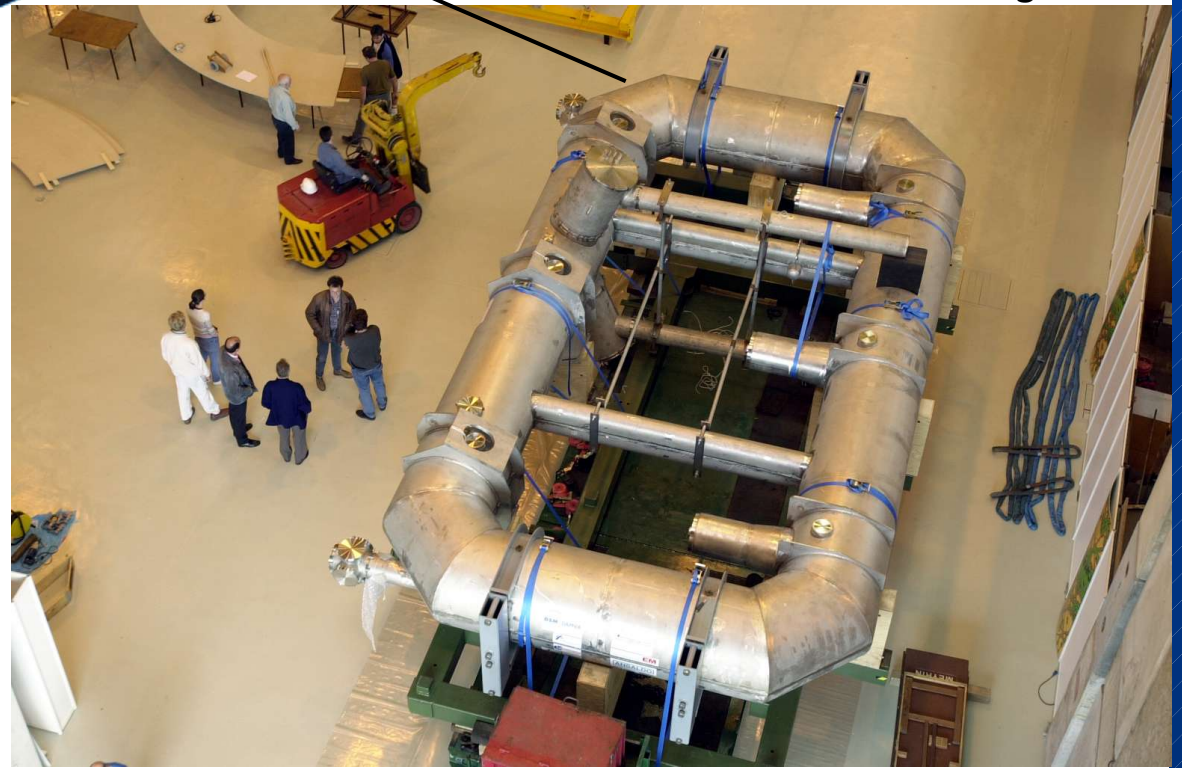
- B reconstruction efficiency
- increased forward acceptance
- improved σ_{d0} ($\sim 20 \mu\text{m}$)

The ATLAS Experiment



Toroid Magnets

Muon Chambers

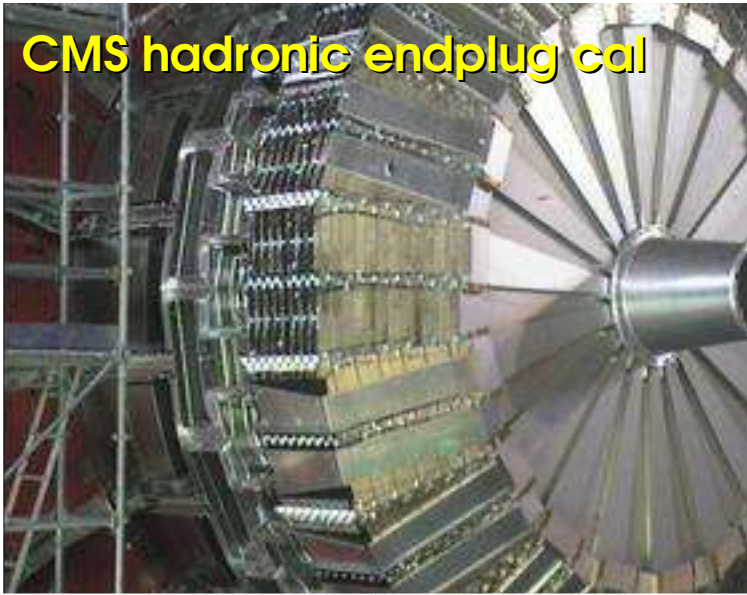


The CMS Experiment

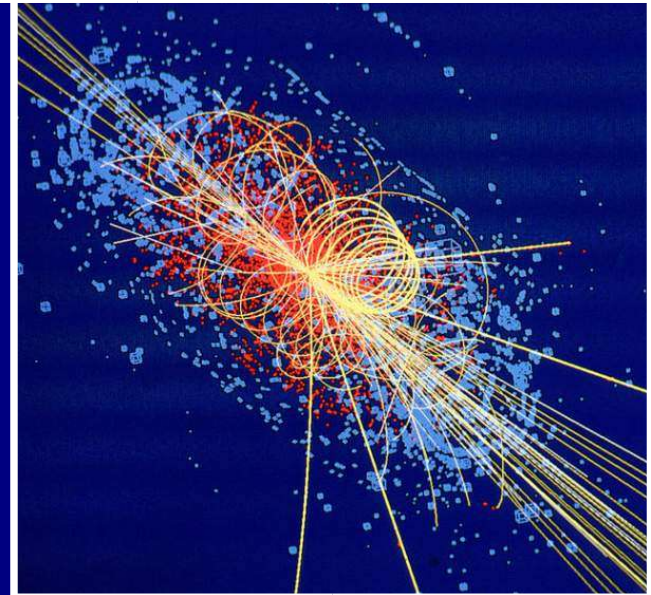
CMS return yoke and solenoid cryostat



CMS hadronic endplug cal



- assembly advanced
- lowering into cavern in 2005
- software and simulation being prepared as well ...

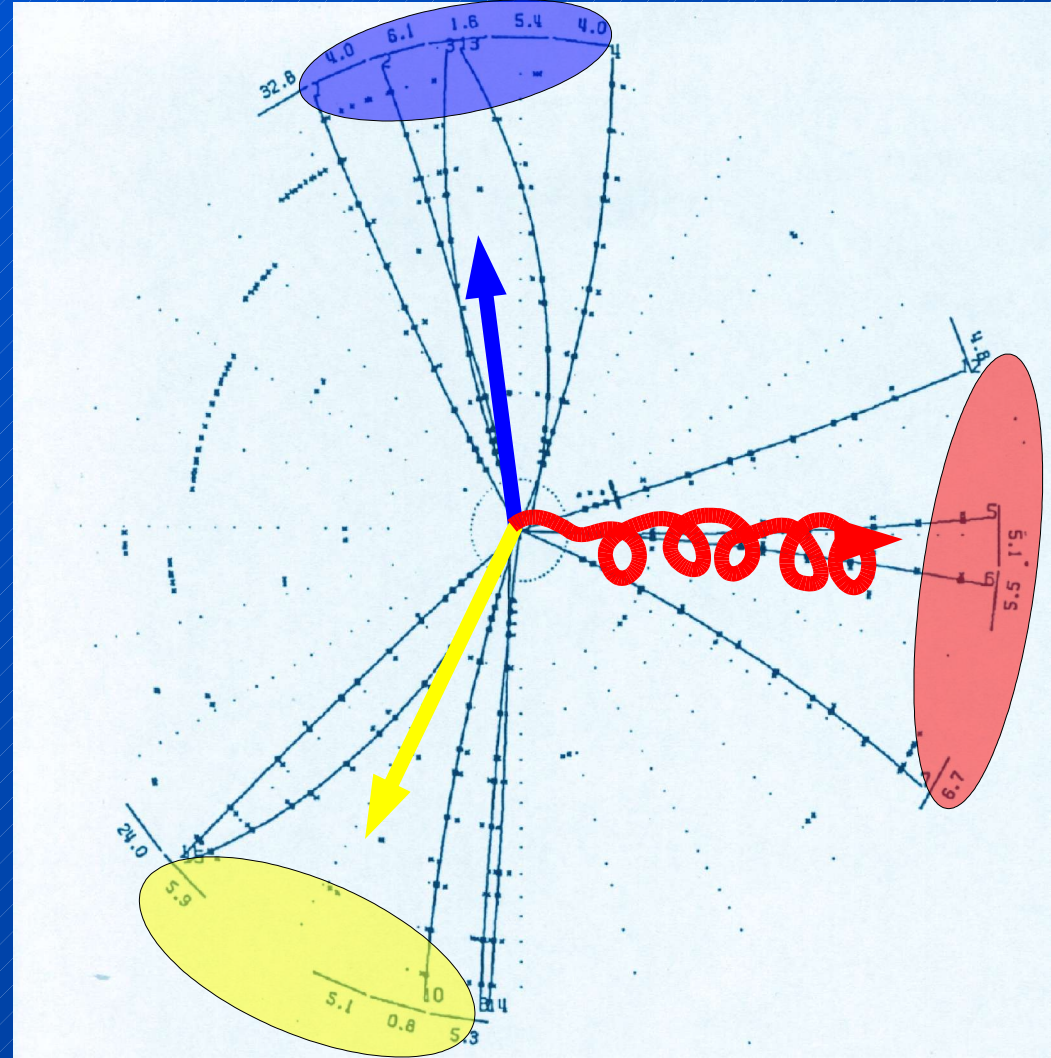


QCD Studies

Quantum Chromo Dynamics (QCD)

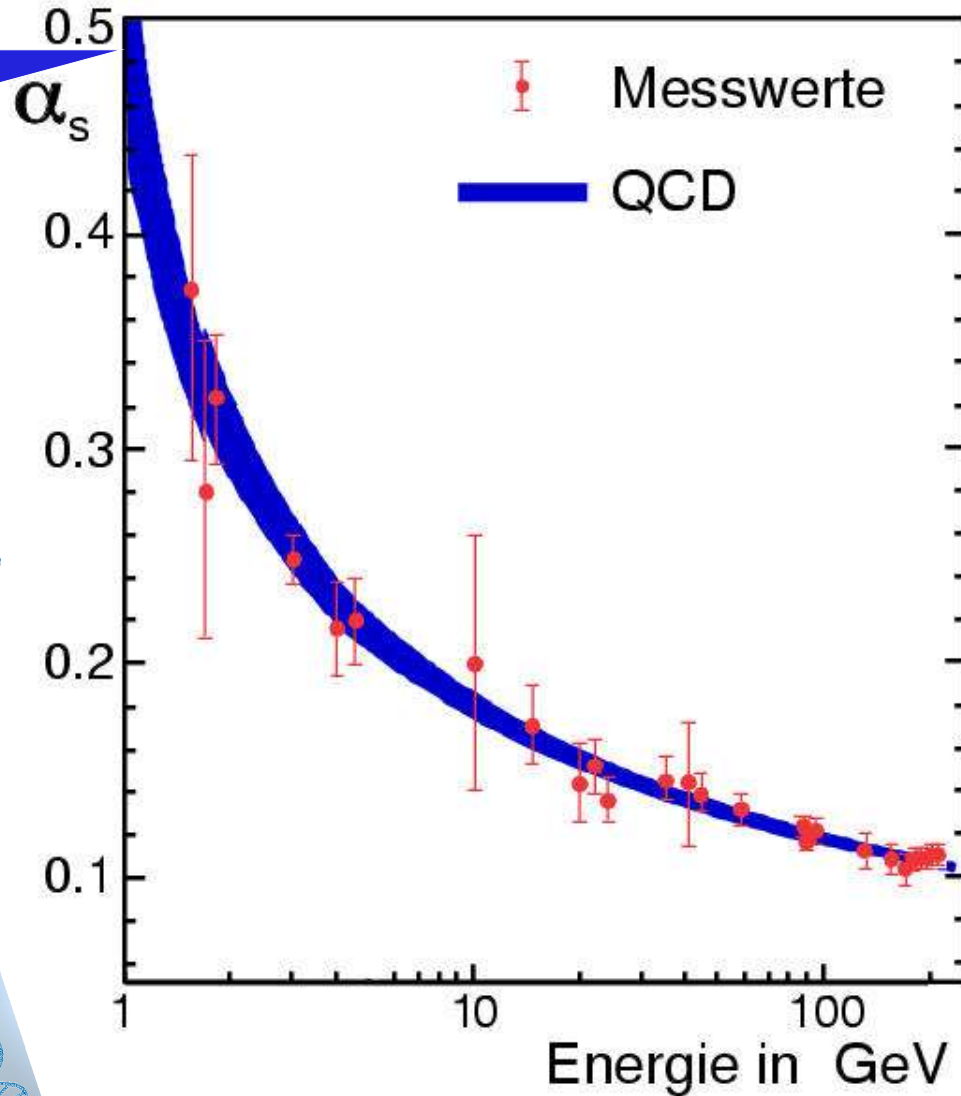
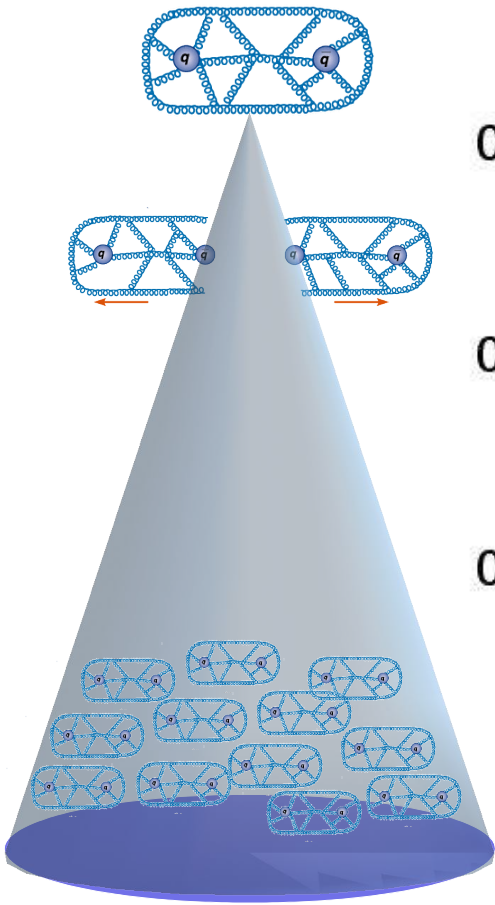
- Theory of **Strong Interactions**
- Acts on all quarks
- Mediated via **Gluons**
... discovered in 70ties at PETRA / DESY ...

Jets



Quantum Chromo Dynamics (QCD)

quark
confinement



asymptotic
freedom

quasi-free quarks

large

distances

small

distances

QCD Overview

- ♦ Rapidity Gaps/Diffractive/Elastic Physics
- ♦ PDF's: double parton interactions, W charge asymmetry
- ♦ non-perturbative QCD: jet shapes, W/Z boson p_T spectra
- ♦ perturbative QCD, particle cross section, W/Z bosons, prompt photons, jets, top, b/c quarks
- ♦ Perturbative QCD with W/Z bosons
- ♦ Perturbative QCD with jets: α_s , jet topologies
- ♦ ...

- ♦ **Inclusive jet cross section** ... here only a small selection ...
- ♦ **Dijet mass spectra**
- ♦ **uncorrected dijet $\Delta\varphi$**
- ♦ **jet multiplicities in W+jet events**

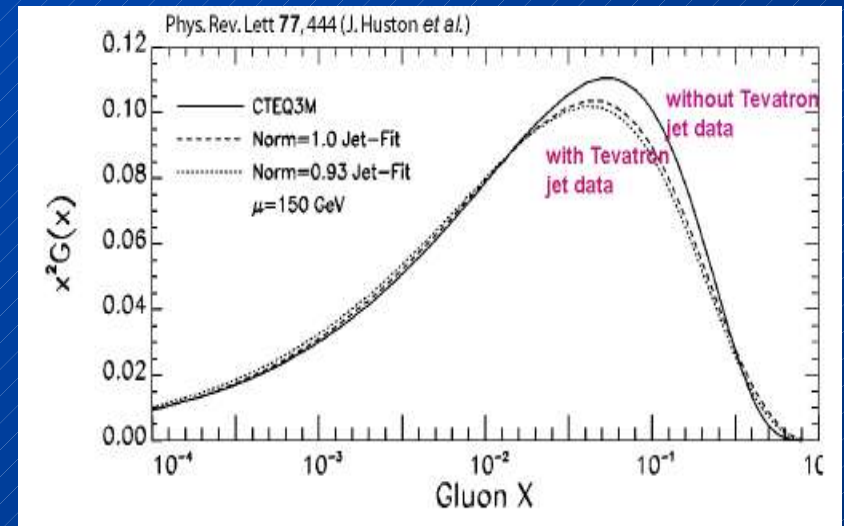
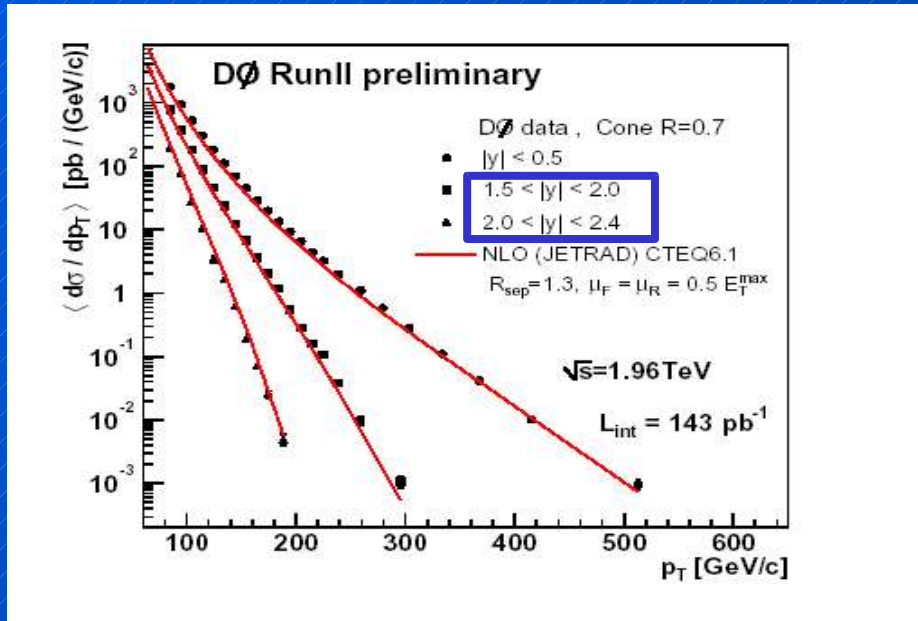
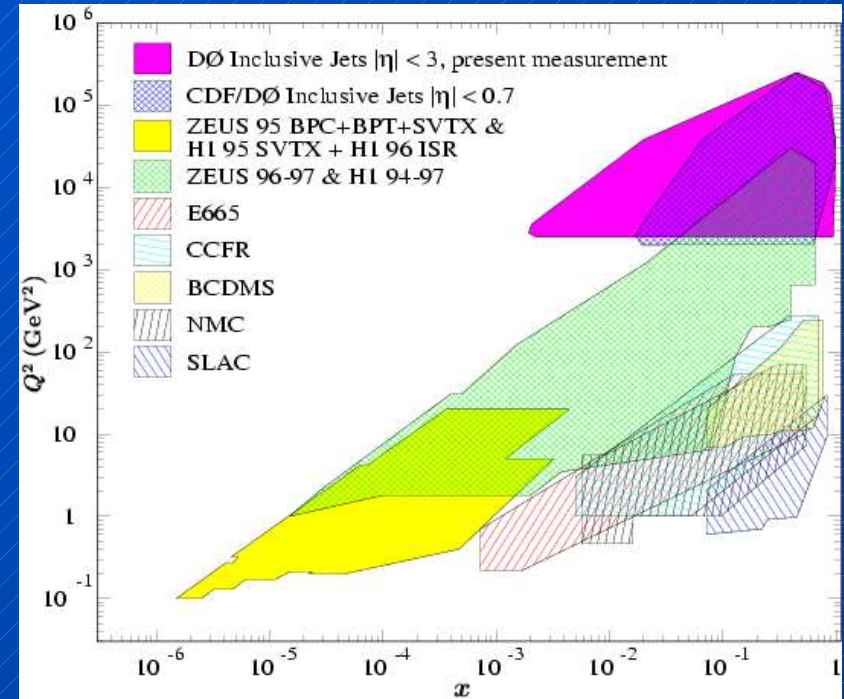
Inclusive Jet Cross Section (1)

High E_T jets probe large x

PDF's, especially gluon

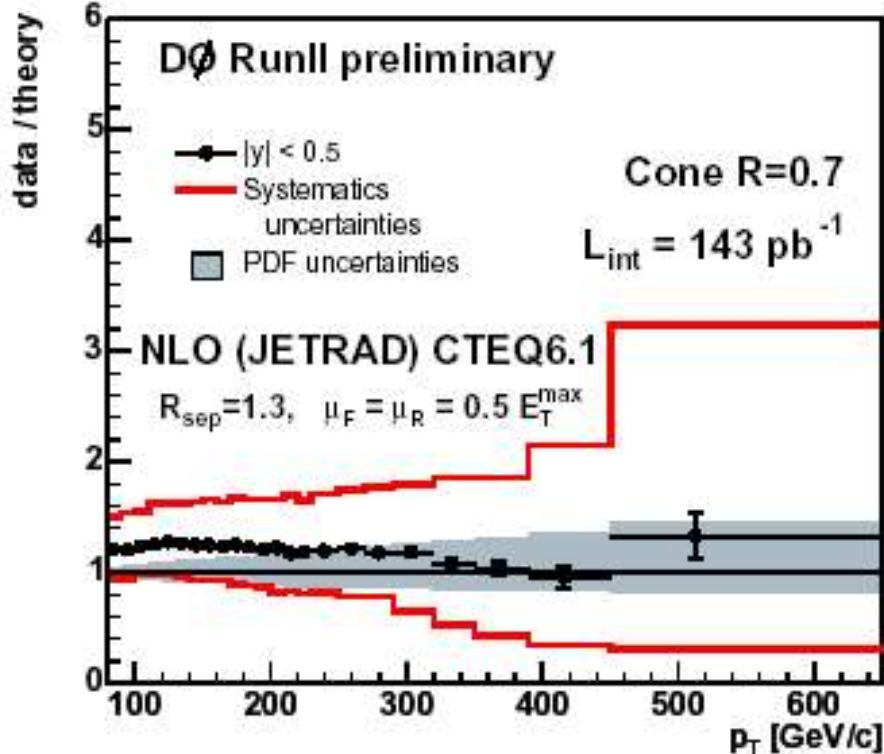
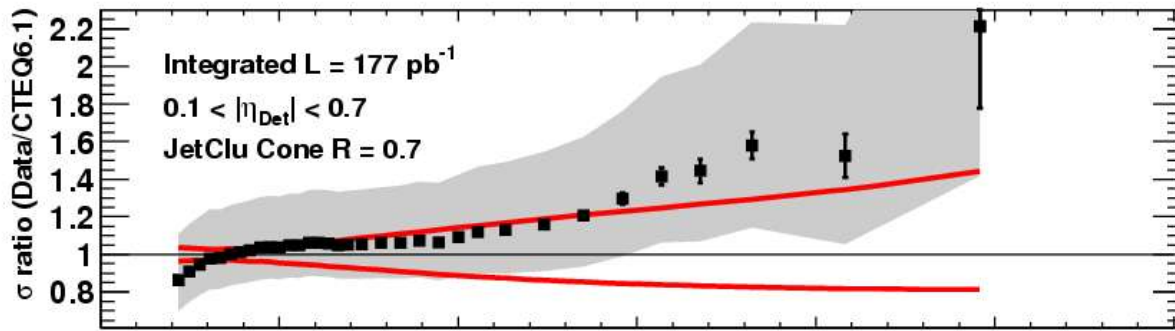
PDF: Run II has extended reach in jet E_T

important information is in
the cross section versus rapidity



Inclusive Jet Cross Section (2)

CDF Run II Preliminary



Uncertainties dominated by jet energy scale. Jet energy is systematics dominated in central. Needs data and time to study and understand more precisely.

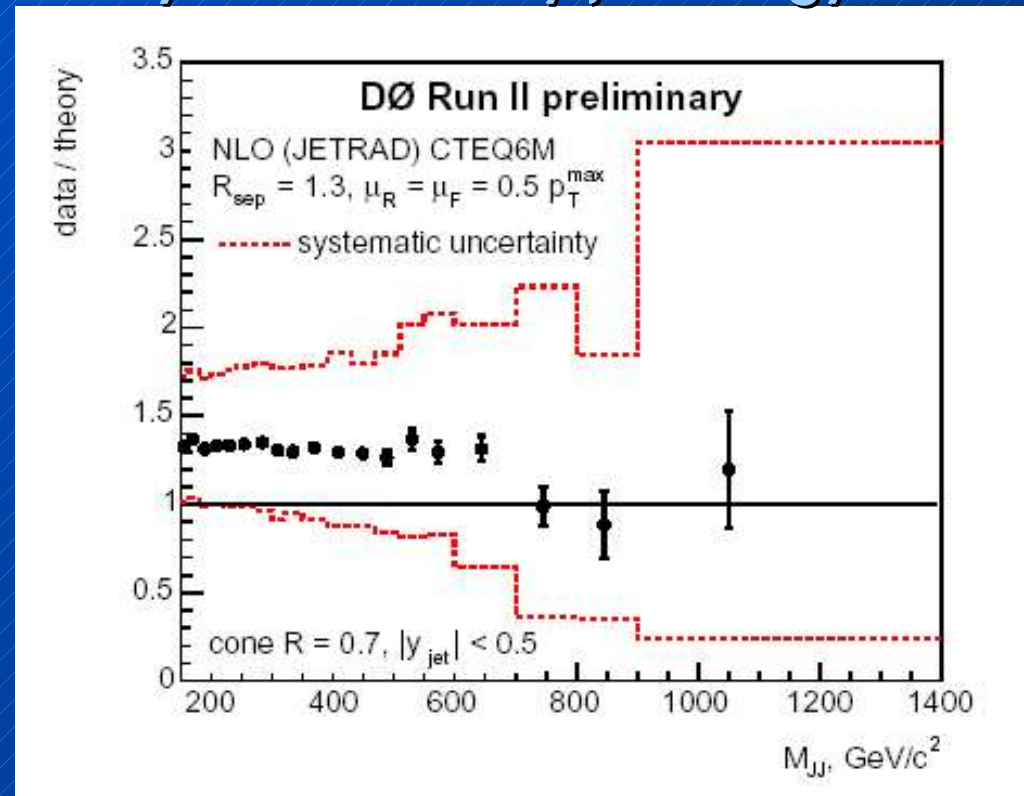
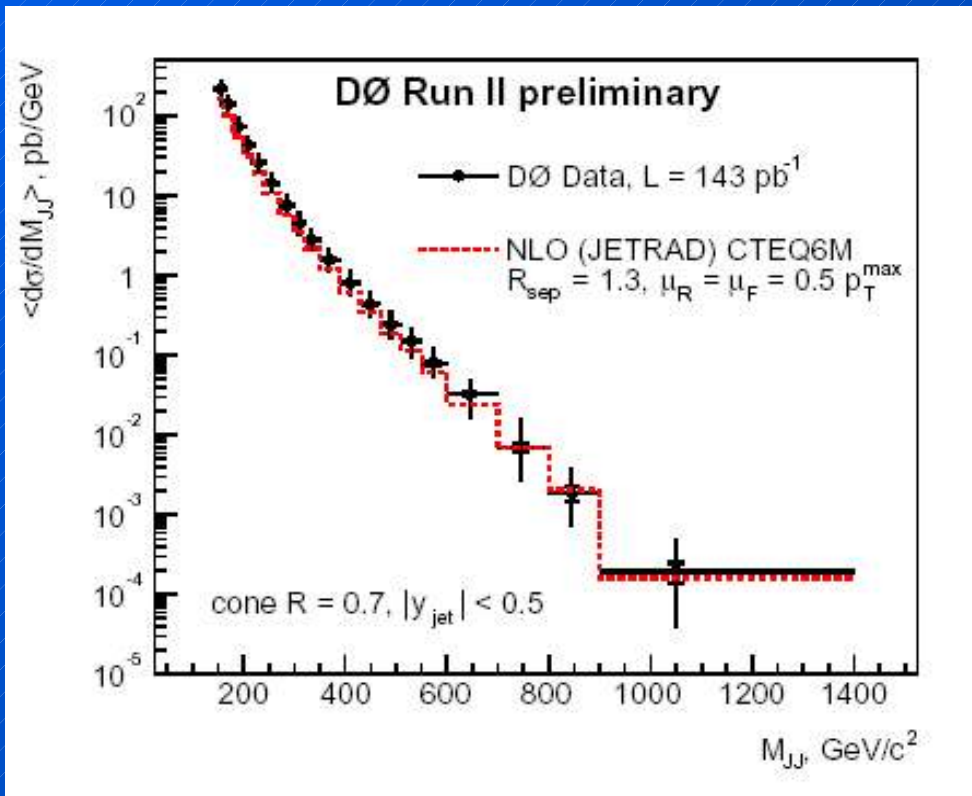
Uncertainty for central jets in central region (GeV)

	Run I	Run II
CDF	17	30
DØ	14	14

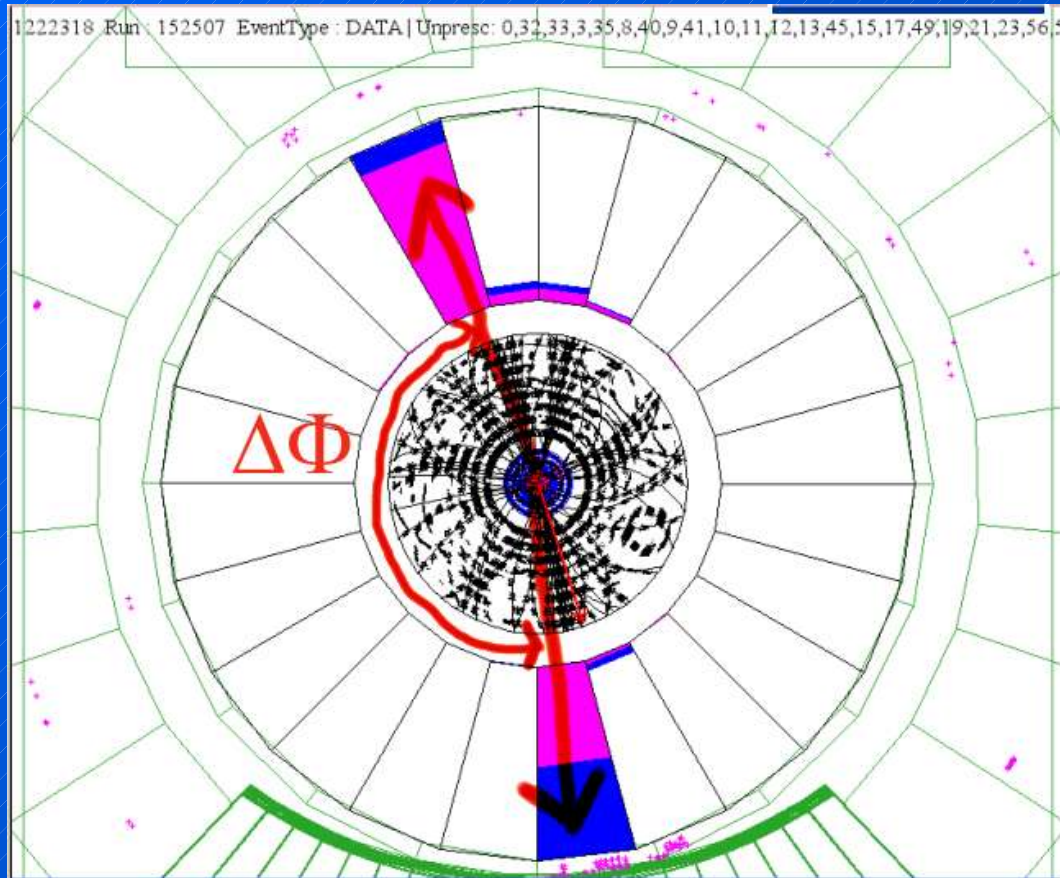
Dijet Cross Section

Often used to search for new resonances ...

Uncertainty dominated by jet energy scale



Φ Decorrelation



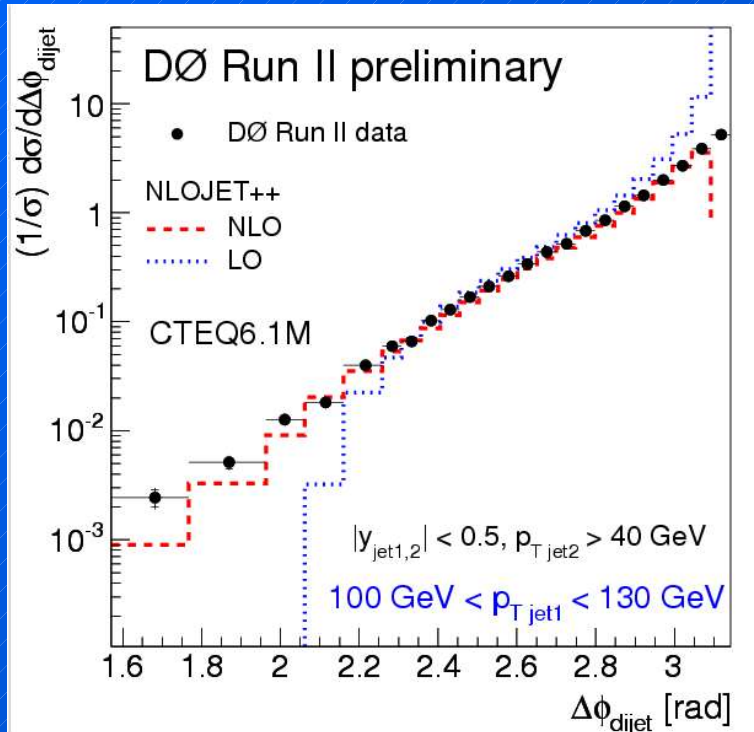
high mass event:

$M_{JJ} = 1364 \text{ GeV}$, $ET = 633, 666 \text{ GeV}$

Sensitive to radiation of
beyond 2 jets without actually
measuring them. Test of pQCD

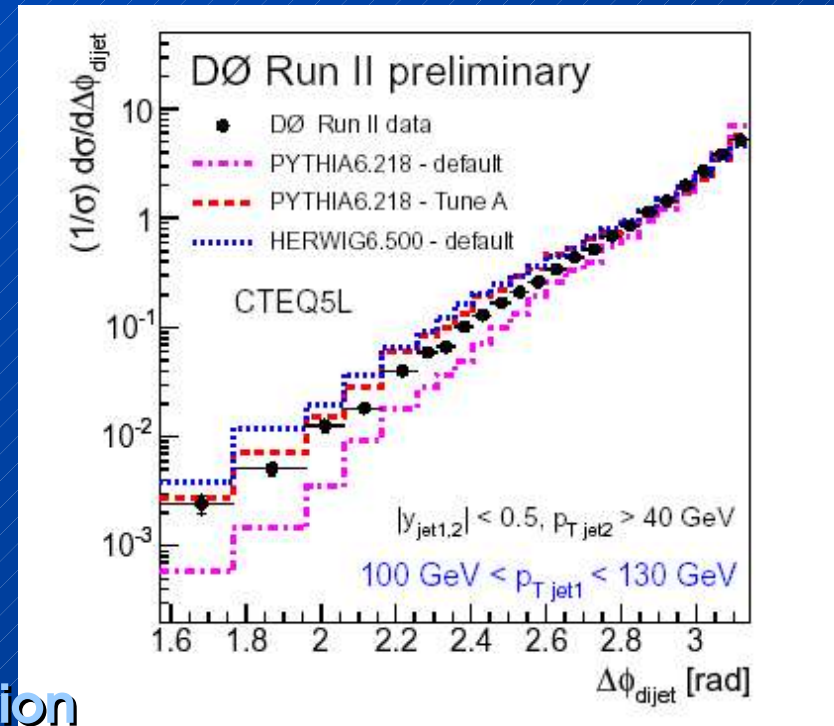
$\Delta\phi = 180^\circ$ at LO

Φ Decorrelation



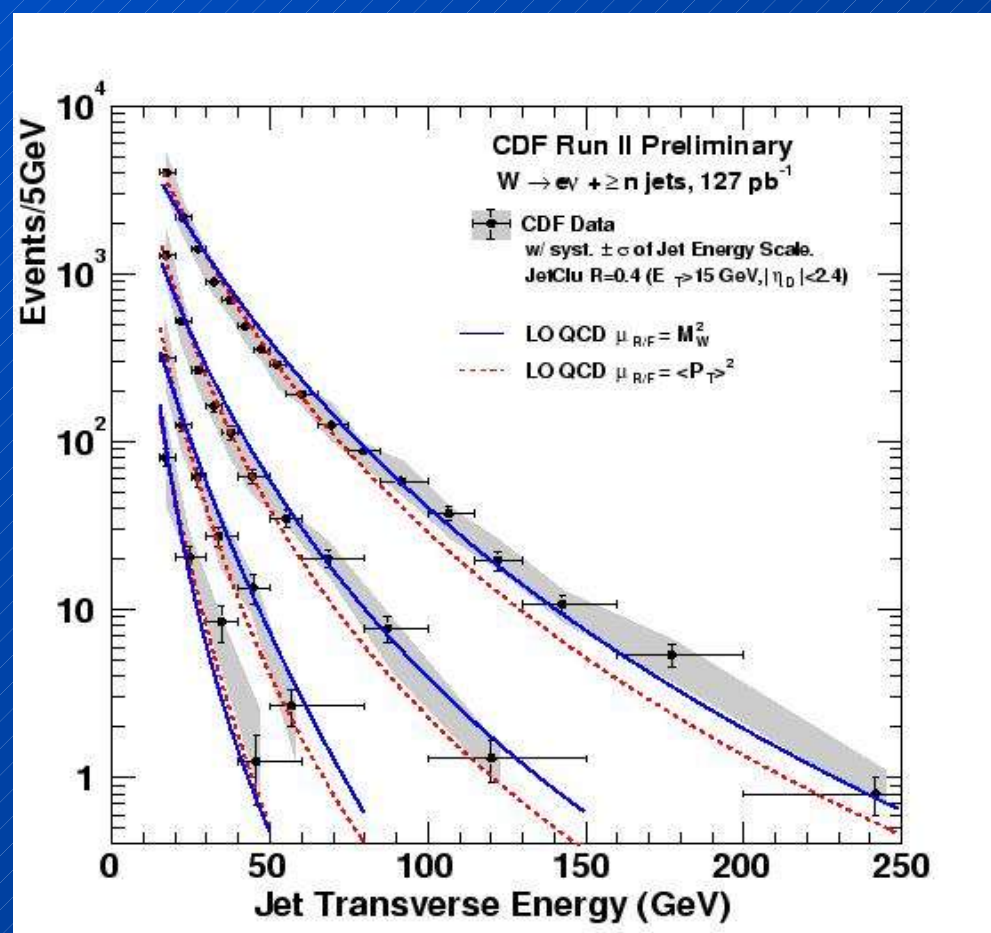
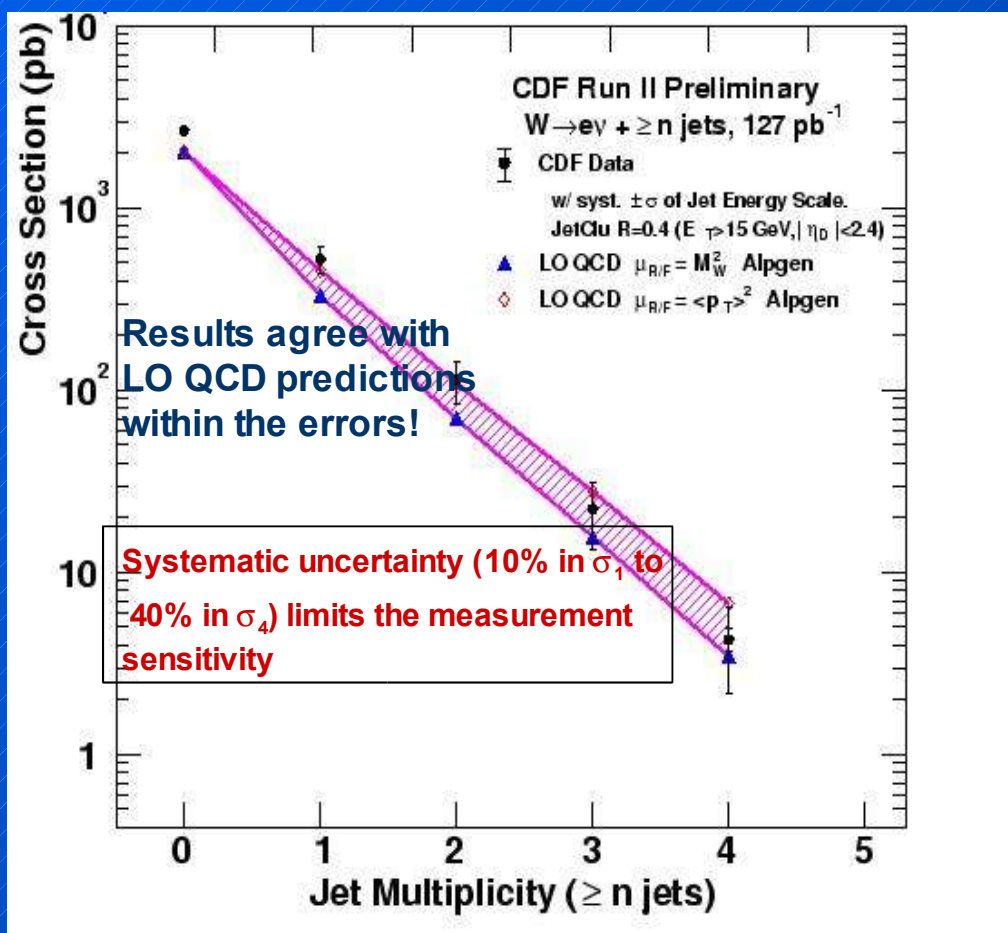
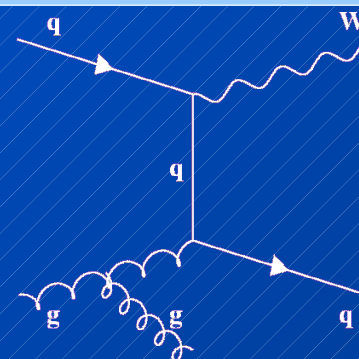
- ♦ Run II differential measurement at small $\Delta\eta$
- ♦ LO (in 3rd jet) perturbative calculation (JETRAD) does not agree
- ♦ lots of 4 jet events at smaller $\Delta\phi$
- ♦ NLO calculation pretty good

- ♦ good agreement with HERWIG and PYTHIA
- ♦ tuned PYTHIA (soft underlying event)
- ♦ gives best agreement
- ♦ NLO calculation is better in intermediate region



W+Jet Events

- Crucial to be able to calculate/understand this process for top & Higgs physics
- ALPGEN LO matrix element interfaced to HERWIG for parton shower
- not more than one parton associated with a reconstructed jet



Summary - Part I

- **Brief Summary of the Standard Model**
- **Structure of the Proton**
- **Hadron Colliders :**
PS – SPPS – Tevatron – LHC
- **The Experiments :**
CDF – DØ – ATLAS – CMS
- **Studies of the Strong Interaction (QCD)**

Backup Slides

The Fermilab Site



Protons or Antiprotons – The Machine

- ♦ History of complex !
 - What machines do you already have ?
- ♦ Pbarp, single magnetic ring
 - SpbarS at CERN
 - Tevatron at FNAL
- ♦ BUT need pbar production and cooling, a limitation
- ♦ pp two magnetic rings, but helped by 2 in 1 design
 - (SSC)
 - Large Hardon Collider

Tevatron Luminosity

Frequency

Protons in
Bunch

Total
Antiprotons

$$L = \frac{f N_p (B \cdot N_{\bar{p}})}{2\pi (\sigma_p^2 + \sigma_{\bar{p}}^2)} F(\sigma_z / \beta^*)$$

Beam Sizes

Beam Shape Form factor
at Intersection

Antiprotons



Cooling

- multiple stochastic cooling systems
- different bandwidth systems react to different characteristics of the beam

Acceleration

- *Main Injector* 8 to 150 GeV
- *Tevatron* 150 – 1000 GeV

Recycler Ring

- 8 GeV
- permanent storage ring; magnetic field controlled by mechanical construction of magnets (reliable, less dependent on power glitches !)
- further storage for accumulation of antiprotons
- filled from *Accumulator*
- *Electron Cooling* cooling antiprotons after production and accumulator
- will support several hundred times 10^{10} antiprotons

Production

- 120 GeV Protons impact on target
- 8 GeV antiprotons produced, large angles
- focussed using Lithium Lens

Accumulation

- pbars injected into large aperture accelerators
- *Debuncher*
- *Accumulator*

Both Accumulator and Recycler used to feed single store in TEVATRON !

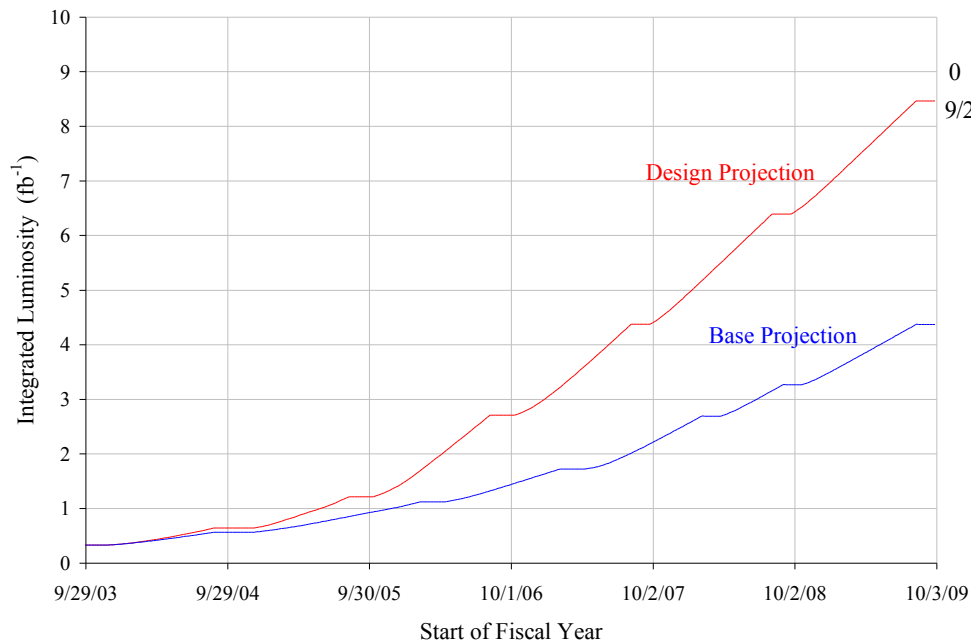
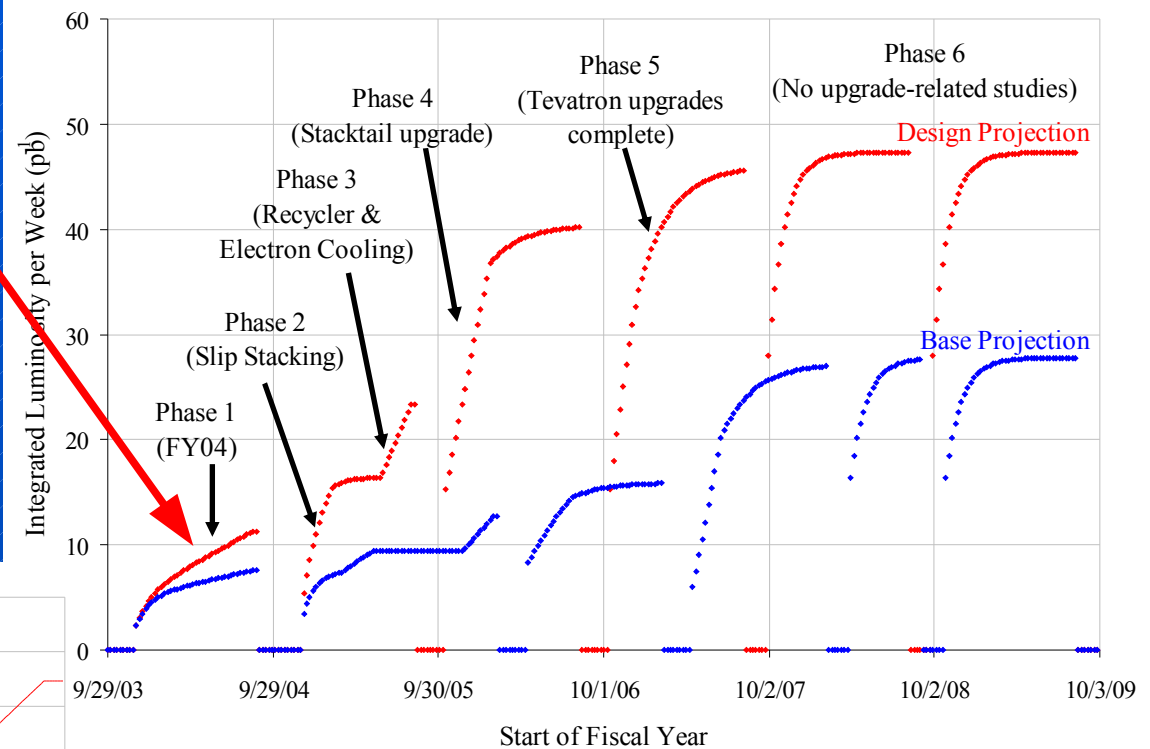
Record Luminosity ($> 1 \cdot 10^{32} \text{ cm}^{-2} \text{ sec}^{-1}$)

The TEVATRON until 2009

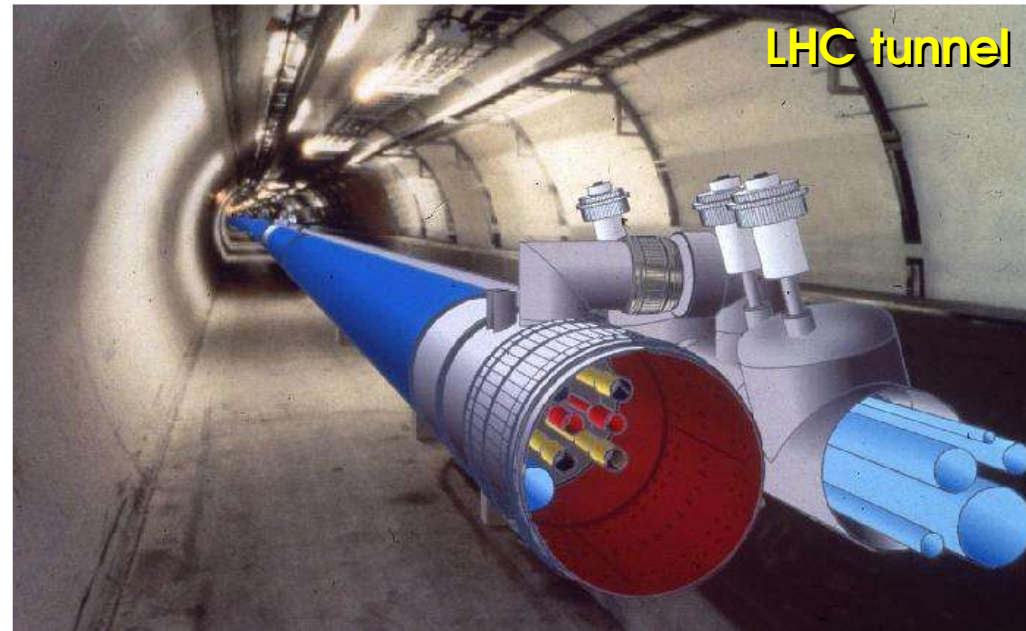
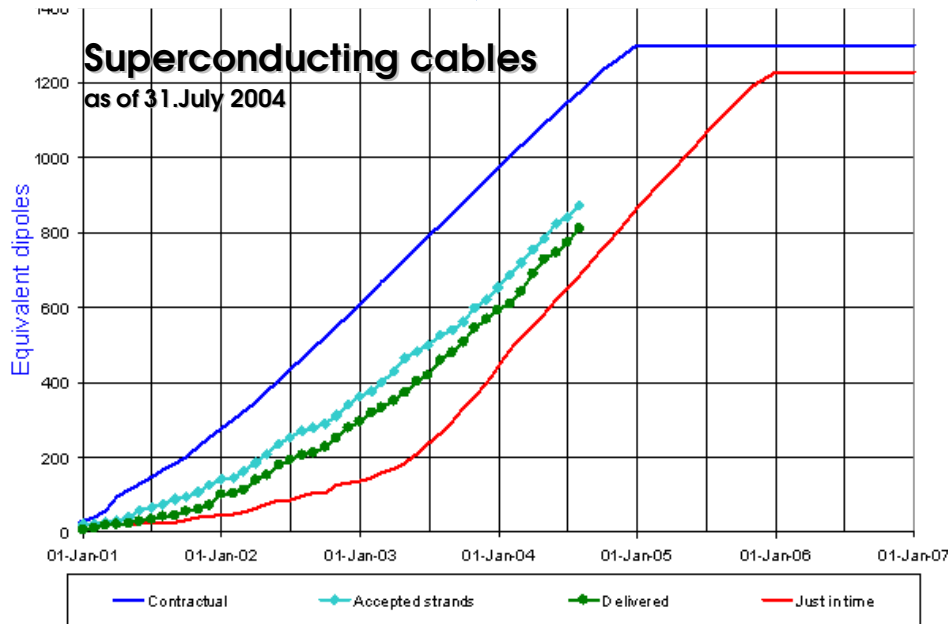
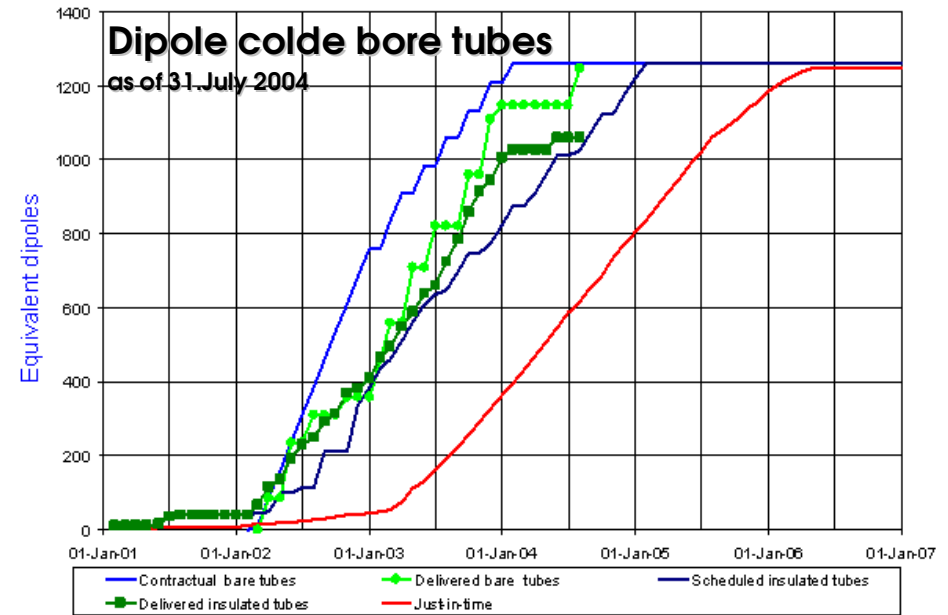
luminosity projection (fb^{-1})

year	baseline	design
2003	0.28	0.3
2004	0.59	0.68
2005	0.98	1.36
2006	1.48	2.24
2007	2.11	3.78
2008	3.25	6.15
2009	4.41	8.57

we are here



The Large Hadron Collider



The DØ Fiber Tracker

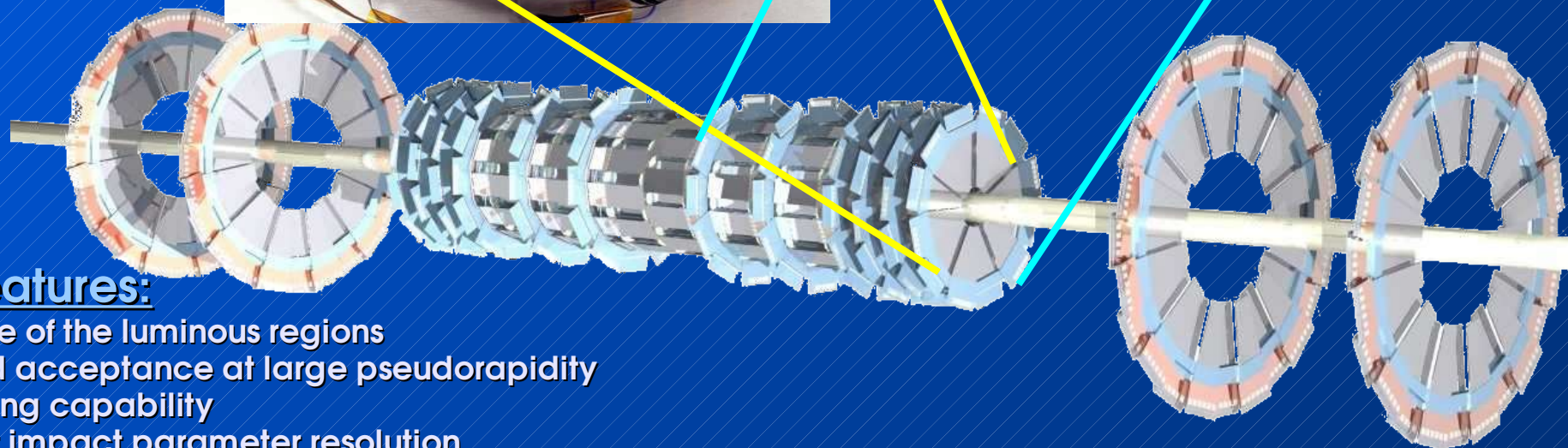
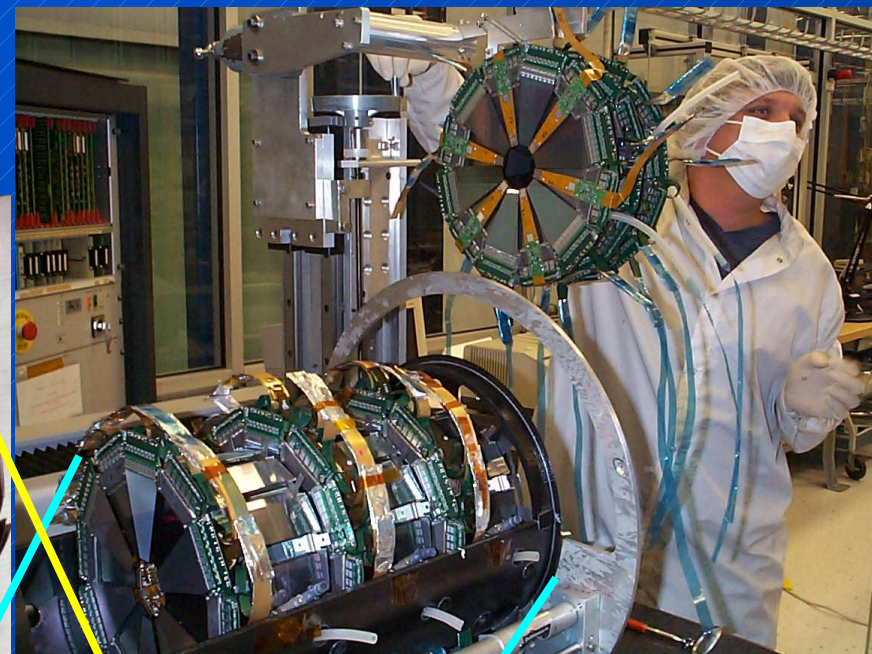
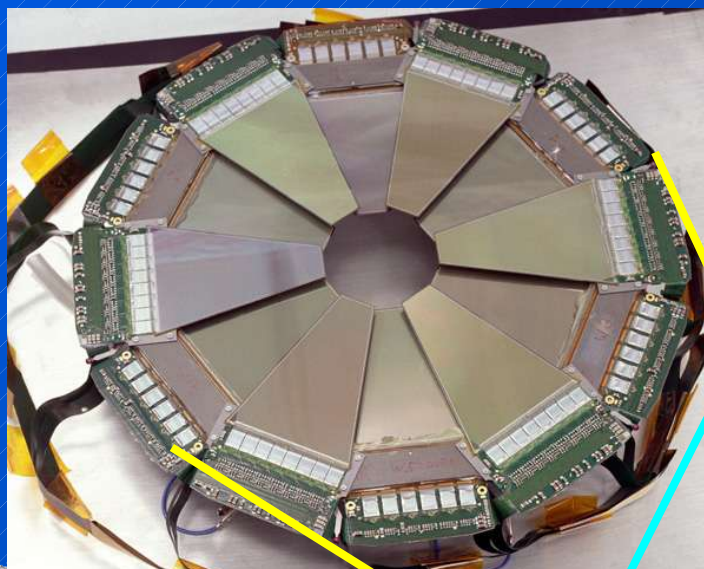


- **Silicon Microstrip Tracker (SMT):**

- 6 barrels, 14 disks
- Tracking out to $\sim \eta = 3$

- axial, double-sided
small-angle stereo and
double-sided 90°
detectors

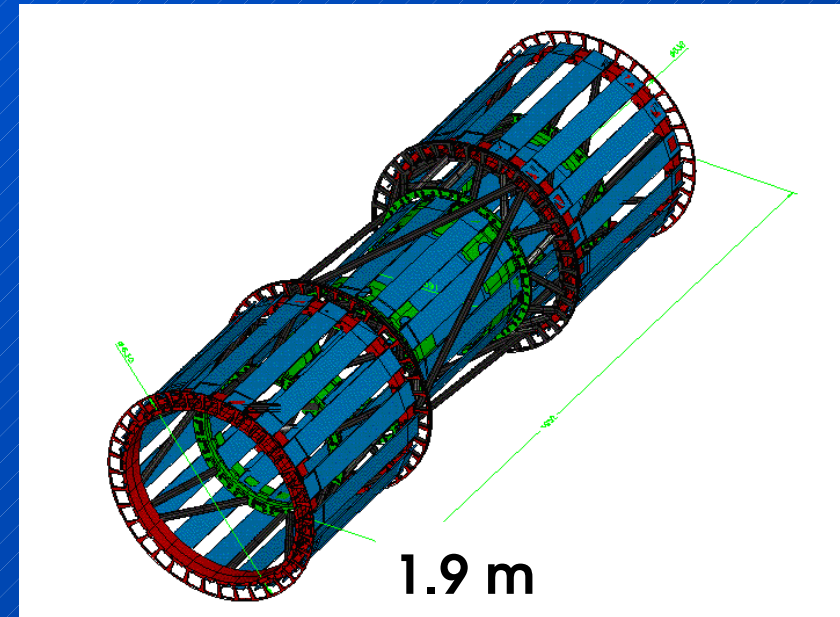
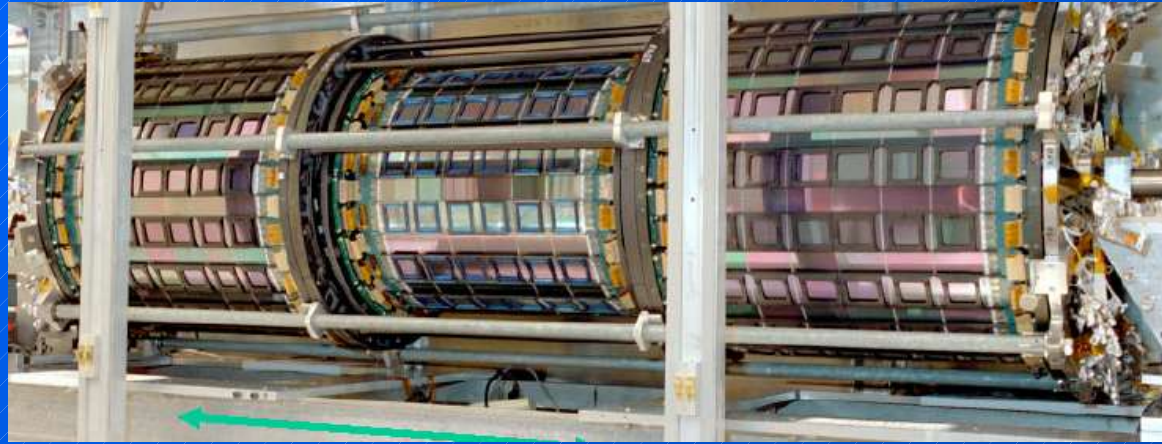
- 800k channels,
SVX2 readout



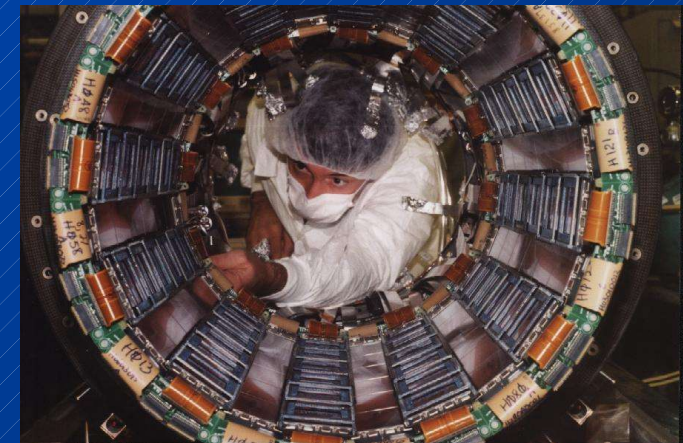
Main features:

- Coverage of the luminous regions
- Extended acceptance at large pseudorapidity
- 3D Tracking capability
- Excellent impact parameter resolution

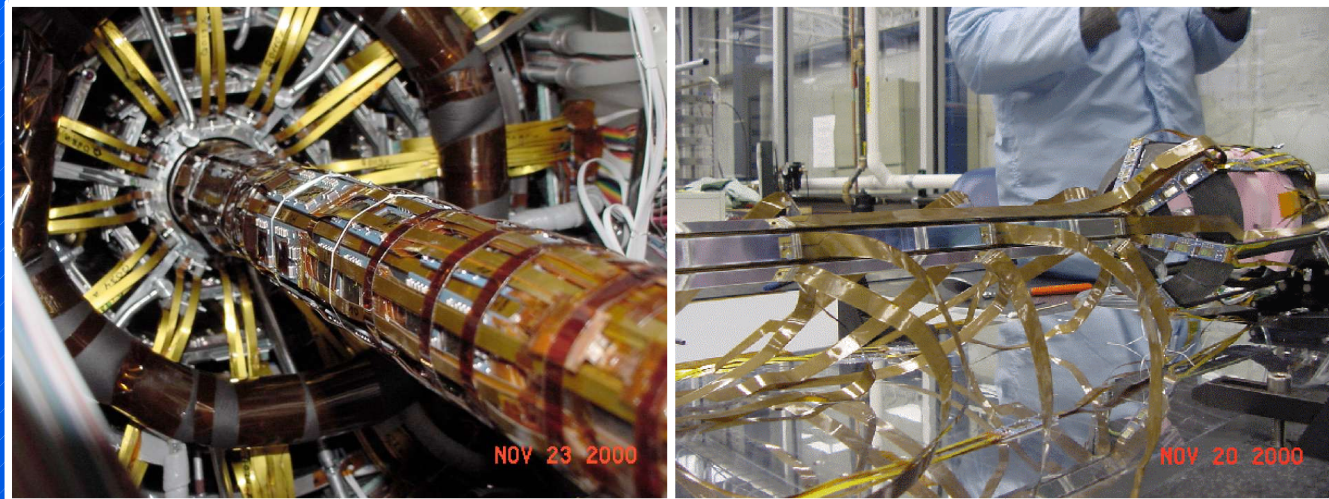
Intermediate Silicon Layer (ISL)



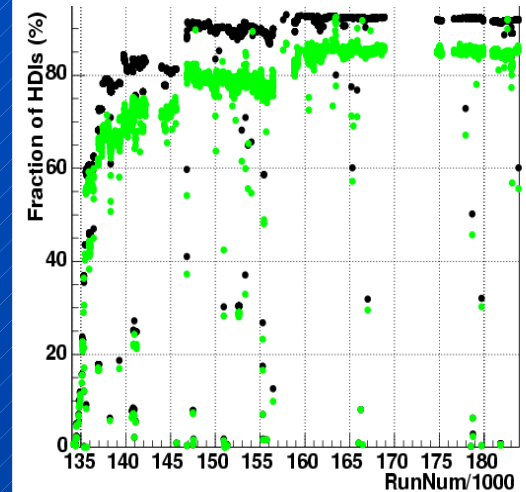
- One central layer ($|\eta| < 1$)
 - link tracks from drift chamber to SVXII
- two forward layers ($1 < |\eta| < 2$)
 - silicon tracking in forward regions
- simple design
 - not used in trigger (less stringent alignment)
 - hybrids mounted off silicon
 - one kind of double-sided sensor flavour



Upgrade more ... L00



total Si status since 2002



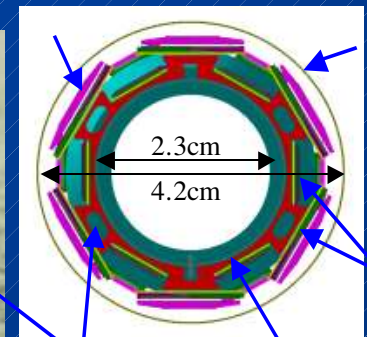
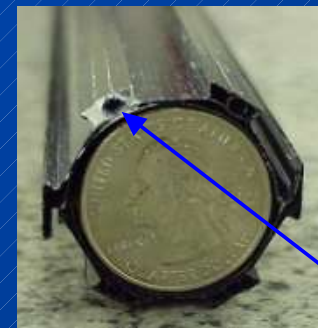
Powered
good (<1% errors)

Precision position measurement before scattering
single-sided layer mounted on beam pipe

- 24 μm pitch; every other strip read out
- 0.6 % X_0 (no cooling) –
1.0% X_0 (cooling, 15% of phi)
- actively cooled
- electronics at either end, large radii
- rad hard silicon; capable of 500V bias
(will outlast SVXII inner layer)

lightweight signal &
bias cables (Kapton)

SVX II inner
bore



sensors

cooling channel Be beampipe

Comparison between ATLAS and CMS

	ATLAS	CMS
Magnet(s)	<p>`dr-core' toroid + inner solenoid</p> <p>2 independent magnet systems</p> <p>Calorimeter outside of solenoid</p>	<p>solenoid</p> <p>1 magnet</p> <p>Calorimeter inside solenoid</p>
Tracking detector	<p>Si pixel + strips (3 + 4 layers)</p> <p>TRD for particle ID, B = 2 T</p> <p>$\sigma/p_T \sim 5 \cdot 10^{-4} p_T \oplus 0.01$</p>	<p>Si pixel + strips (2 + 10 layers)</p> <p>B = 4 T</p> <p>$\sigma/p_T \sim 1.5 \cdot 10^{-4} p_T \oplus 0.005$</p>
EM calorimeter	<p>Pb liquid argon</p> <p>$\sigma/E \sim 10\% / \sqrt{E}$</p> <p>Presampler + longitudinal segmentation</p>	<p>PbWO4 crystals</p> <p>$\sigma/E \sim 3-5\% / \sqrt{E}$</p> <p>No longitudinal segmentation</p>
Hadron calorimeter	<p>Fe-scintillator (barrel) +</p> <p>Cu liquid Argon (endcaps) (10 λ)</p> <p>$\sigma/E \sim 50\% / \sqrt{E} \oplus 0.03$</p>	<p>Cu scintillator (5.8 λ + catcher)</p> <p>$\sigma/E \sim 65\% / \sqrt{E} \oplus 0.05$</p>
Muons	<p>Air $\Rightarrow \sigma/p_T \sim 7\%$ at 1 TeV</p> <p>Standalone</p>	<p>Fe $\Rightarrow \sigma/p_T \sim 5\%$ at 1 TeV</p> <p>Combined with central tracker</p>

Run II Jet Algorithm

- ◆ Use Run II cone algorithm
- ◆ Combine particles in a $R=0.7$ cone
- ◆ Use the four vector of every tower as a seed
- ◆ Rerun using the midpoints between pairs of jets as seed
- ◆ Overlapping jets merged if the overlap area contains more than 50% of lower p_T jet, otherwise particles assigned to nearest jet.

Both experiments now using same algorithm

Reduced sensitivity to soft radiation

E-scheme recombination

$$P^J = (E^J, p^J) = \sum_{i=\text{all clusters} \in \text{jet}} (E^i, p_x^i, p_y^i, p_z^i)$$

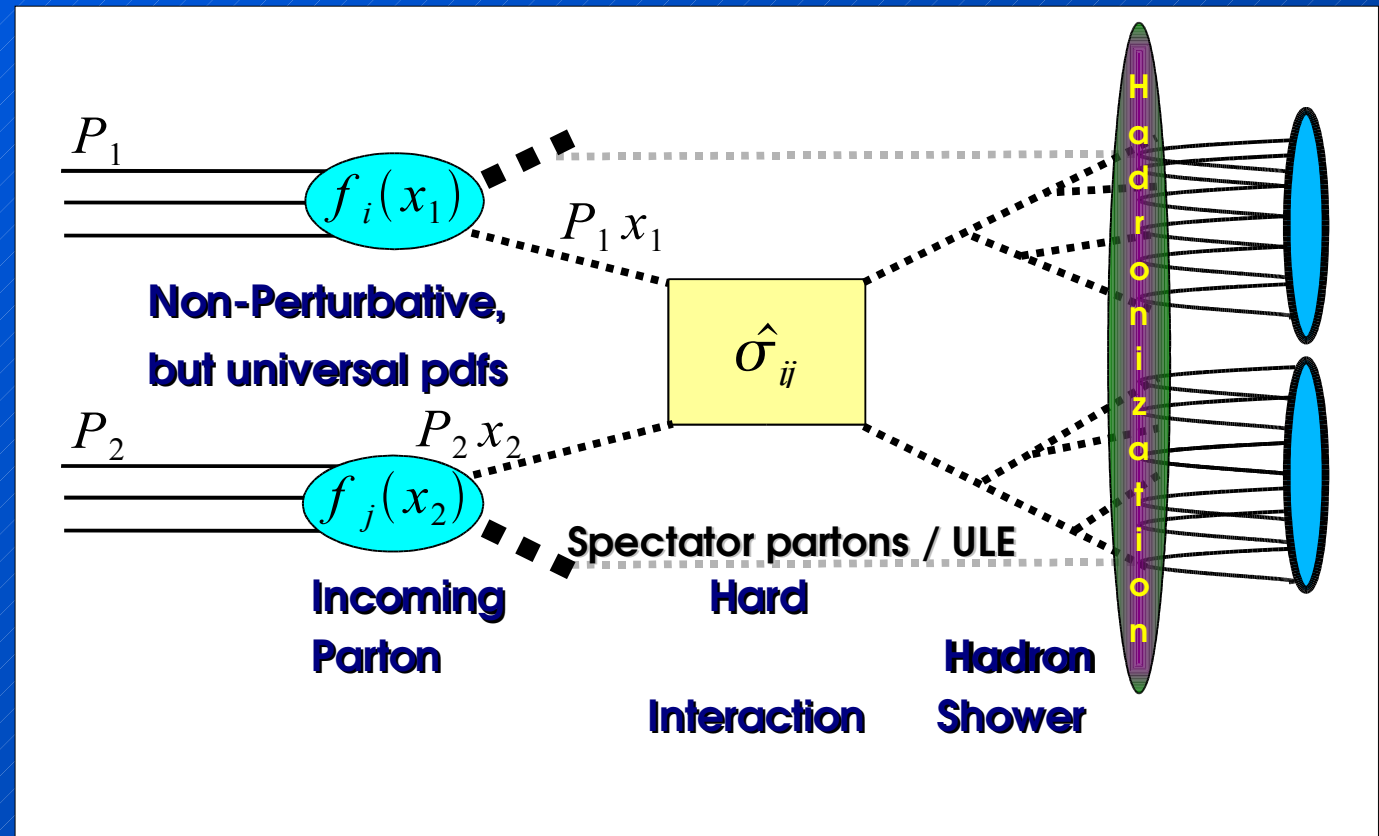
$$p_T^J = \sqrt{(p_x^J)^2 + (p_y^J)^2}$$

$$y^J = \frac{1}{2} \ln \frac{E^J + p_z^J}{E^J - p_z^J}$$

$$\phi^J = \tan^{-1} \frac{p_y^J}{p_x^J}$$

Soft Underlying Event

The “**underlying event**” consists of hard initial & final-state radiation plus the “beam-beam remnants” and possible multiple parton interactions.



Learn through studies of min bias events, Jet events. Look at distributions/correlations of charged particles with $\eta < 1$, $p_T > 500$ MeV

Also, studies of mini-jets in min bias events

Luminosity Measurement

event rate of inelastic events (luminosity detectors)

$$L = \frac{1}{\sigma_{eff}} \cdot \frac{dN}{dt}$$

need to know effective cross section for this process

$$\sigma_{eff} = \epsilon \cdot A \cdot \sigma_{inelastic}$$

Efficiency and geometrical acceptance
(studied and determined using Monte Carlos)

$$\sigma_{inelastic} \equiv \sigma_{total} - \sigma_{elastic}$$

Measured separately by several experiments
(CDF, E811)

	Run II	Run I
Inelastic	60.7 +- 2.4 mb	59.23 +- 2.3 mb
Single diffraction	9.6 +- 0.5 mb	9.6 +- 0.5 mb
Double diffraction	7.0 +- 2.0 mb	7.0 +- 2.0 mb

last two are part of the inelastic cross section
but with difference acceptance ...

Total Proton-Antiproton Cross-Section (1)

$$1) \quad \sigma_T = \frac{1}{L} (R_{el} + R_{inel})$$

$L = \text{integrated luminosity}$

$$2) \quad \sigma_T = \frac{4\pi}{k} \Im F(0) \quad (\text{optical theorem})$$

$$\sigma_T^2 = \frac{16\pi^2}{k^2} \cdot \frac{\Im F(0)^2}{\Im F(0)^2 + \Re F(0)^2} \cdot \underbrace{(\Im F(0)^2 + \Re F(0)^2)}_{|F(0)|^2}$$

$$|F(0)|^2 = \left(\frac{d\sigma}{d\Omega} \right)_{el}^{\Theta=0} = \frac{1}{L} \frac{1}{2\pi} \frac{4k^2}{2} \left(\frac{dR_{el}}{dt} \right)_{t=0}$$

$$t = \frac{-s}{2} (1 - \cos \Theta) \Rightarrow d \cos \Theta = \frac{2}{s} dt$$

$$s = 4k^2$$

$$3) \quad \sigma_T^2 = \frac{16\pi (Ac)^2}{(1+\rho^2)} \cdot \frac{1}{L} \left(\frac{dR_{el}}{dt} \right)_{t=0} \quad \text{with} \quad \rho = \frac{\Re F(0)}{\Im F(0)}$$

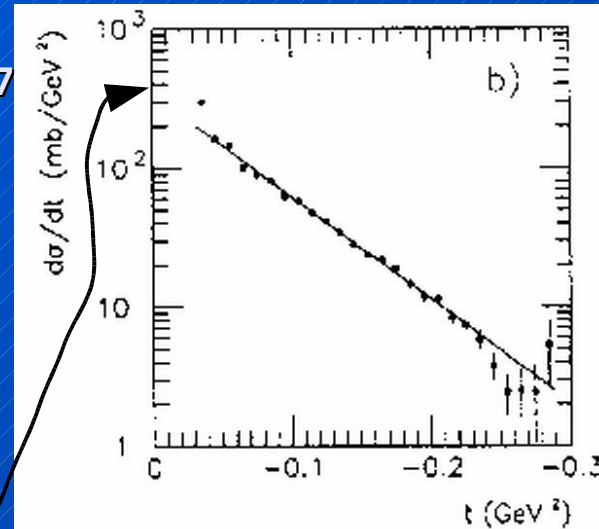
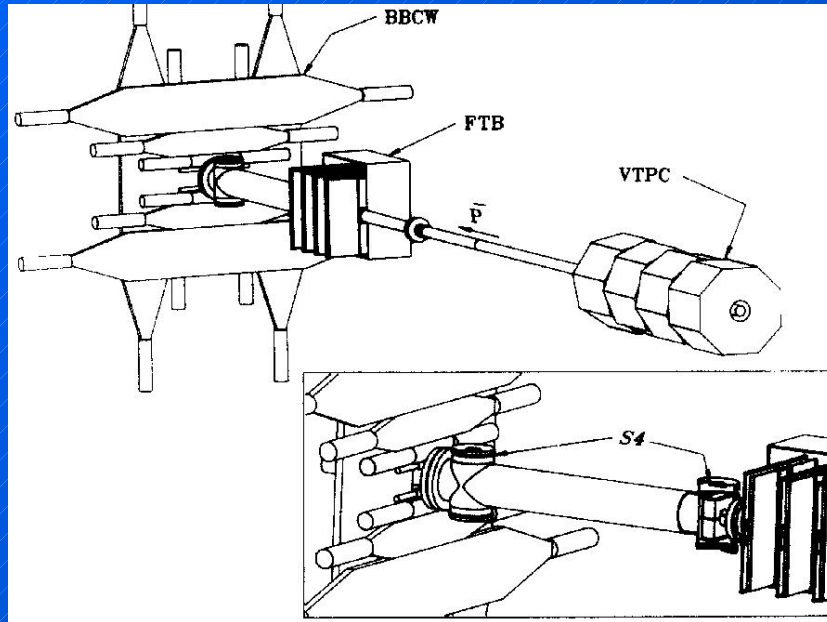
4) dividing 3) by 1) gives :

$$\sigma_T = \frac{16\pi (Ac)^2}{(1+\rho^2)} \cdot \frac{\left(\frac{dR_{el}}{dt} \right)_{t=0}}{R_{el} + R_{inel}}$$

... total cross-section
can be determined from event
rates, i.e. without knowing the
luminosity !!!

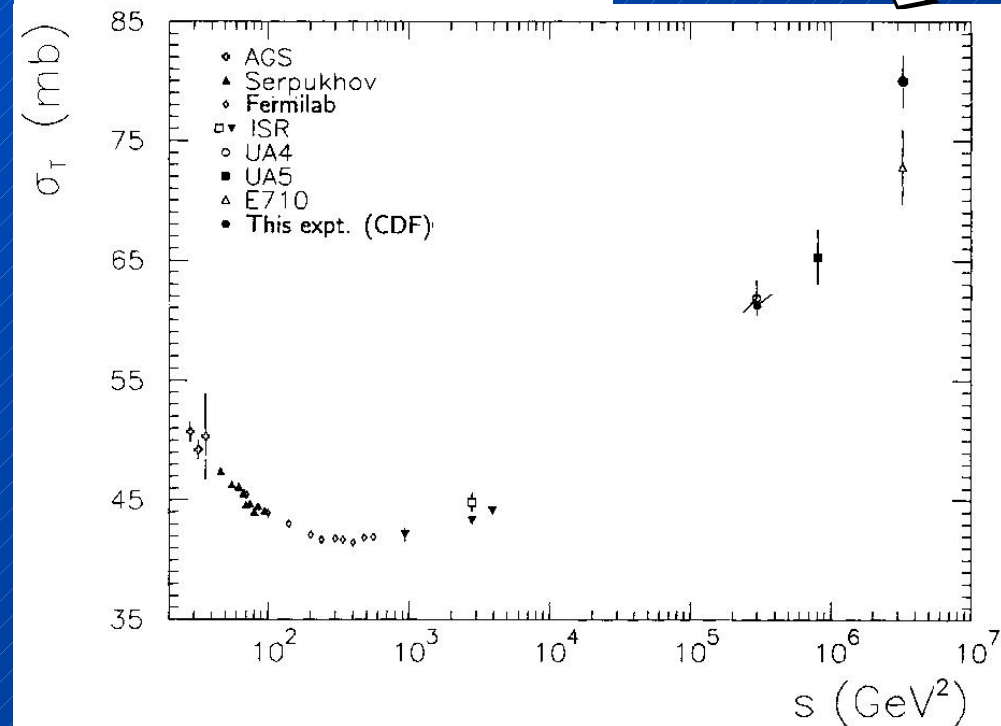
Total Proton-Antiproton Cross-Section (2)

inelastic proton-antiproton events identified by TPC, telescopes, drift chambers and scintillators up to $|\eta| < 6.7$



here shown is the differential t -distribution; needed is only the event rate at the $t=0$ intercept; t is measured using roman pot detectors.

CDF result on total proton-antiproton cross-section



CDF total proton-antiproton cross-section at $\sqrt{s} = 1800$ GeV:

Phys.Rev D 50, 5550 (1994) & Phys.Rev D 50, 5518 (1994)

=0.15 from UA4/2 Phys.Rev.Lett. 68, 2433 (1992)

total inelastic	240 982	$\pm 2 967$
total elastic	78 691	$\pm 1 463$
total	319 673	$\pm 3 308$
$(dE_{el} / dt)_{t=0}$	1 336 532	$\pm 40 943$

$$\Rightarrow \sigma_T^{p\bar{p}}(\sqrt{s} = 1800 \text{ GeV}) = 80.03 \pm 2.24 \text{ mb}$$

Optical Theorem (1)

- Consider beam of **particles incident** in z-direction (plane wave) on a spinless particle (neglect time-dependence $e^{i\omega t}$ here)
 $k = 1/\tilde{\lambda}$ with $2\pi\tilde{\lambda} = \text{deBroglie wavelength}$
- At sufficiently large radial distances $kr \gg 1$:

$$\psi_i = e^{ikz} = \frac{i}{2kr} \sum_l (2l+1) [(-1)^l e^{-ikr} - e^{ikr}] P_l(\cos \Theta)$$
- scattering center can alter phase and amplitude of the **outgoing /th partial waves** by $2\delta_l$ and $2\eta_l$ ($1 > \eta_l > 0$), respectively:

$$\psi_{total} = \frac{i}{2kr} \sum_l (2l+1) [(-1)^l e^{-ikr} - \eta_l e^{2i\delta_l} e^{ikr}] P_l(\cos \Theta)$$
- **scattered wave = total - incoming** :

$$\begin{aligned} \psi_{scatt} &= \psi_{total} - \psi_i = \frac{e^{ikr}}{kr} \sum_l (2l+1) \frac{(\eta_l e^{2i\delta_l} - 1)}{2i} P_l(\cos \Theta) \\ &= \frac{e^{ikr}}{r} F(\Theta) \end{aligned}$$

with the scatt. amplitude $F(\Theta) = \frac{1}{k} \sum_l (2l+1) \left(\frac{\eta_l e^{2i\delta_l} - 1}{2i} \right) P_l(\cos \Theta)$

... elastically scattered wave, since k taken to be same for incoming and outgoing ...
- **scattered outgoing flux in solid angle $d\Omega$, at radius r**

$$v_o \psi_{scatt} \psi_{scatt}^* r^2 d\Omega = v_o |F(\Theta)|^2 d\Omega$$

where v_o is velocity of outgoing particle relative to scattering center

THIS IS the product of the scattering cross-section and the incident flux ($= v_i \psi_i \psi_i^* = v_i$); **elastic scattering** $\Rightarrow v_i = v_o$

$$v_o d\sigma = v_o |F(\Theta)|^2 d\Omega$$

$$\Leftrightarrow$$

$$\left(\frac{d\sigma}{d\Omega} \right)_{el} = |F(\Theta)|^2$$

Optical Theorem (2)

- Legendre polynomials P_l obey orthogonality condition :

$$\int P_l P_{l'} d\Omega = \frac{4\pi\delta_{l,l'}}{2l+1}$$

total elastic scattering cross-section, integrated over angle

$$\sigma_{el.} = 4\pi\lambda \sum_l (2l+1) \left| \frac{\eta_l e^{2i\delta_l} - 1}{2i} \right|^2$$

- if no absorption of incoming wave ... ($\eta_l = 1$):

$$\sigma_{el.} = 4\pi\lambda \sum_l (2l+1) \sin^2 \delta_l$$

- if $\eta_l < 1$ get reaction cross-section from probability conservation:

$$\sigma_r = \int (|\psi_{in}|^2 - |\psi_{out}|^2) r^2 d\Omega$$

ψ_{in} is first term of ..., ψ_{out} is second term of ...

$$\Rightarrow \sigma_r = \pi\lambda^2 \sum_l (2l+1) (1 - \eta_l^2)$$

- total cross-section :

$$\sigma_{tot} = \sigma_{el} + \sigma_r = \pi\lambda^2 \sum_l (2l+1) 2 (1 - \eta_l \cos 2\delta_l)$$

- since $P_l(1) = 1$ for all l and $\cos(\Theta = 0) = 1$

$$\text{Im } F(0) = \frac{1}{2k} \sum_l (2l+1) (1 - \eta_l \cos 2\delta_l)$$

- $\Rightarrow \text{Im } F(0) = \frac{k}{4\pi} \sigma_{tot}$

... total cross-section is related to the imaginary part of the forward elastic scattering amplitude ...